

Ronan Gautier

High-Fidelity Control and Stabilization of Cat Qubits

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Reviewers / Liang Jiang
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Examinators / Christiane Koch
Jean-Michel Raimond
Mario Sigalotti

Invited / Mazyar Mirrahimi
Alexandre Blais

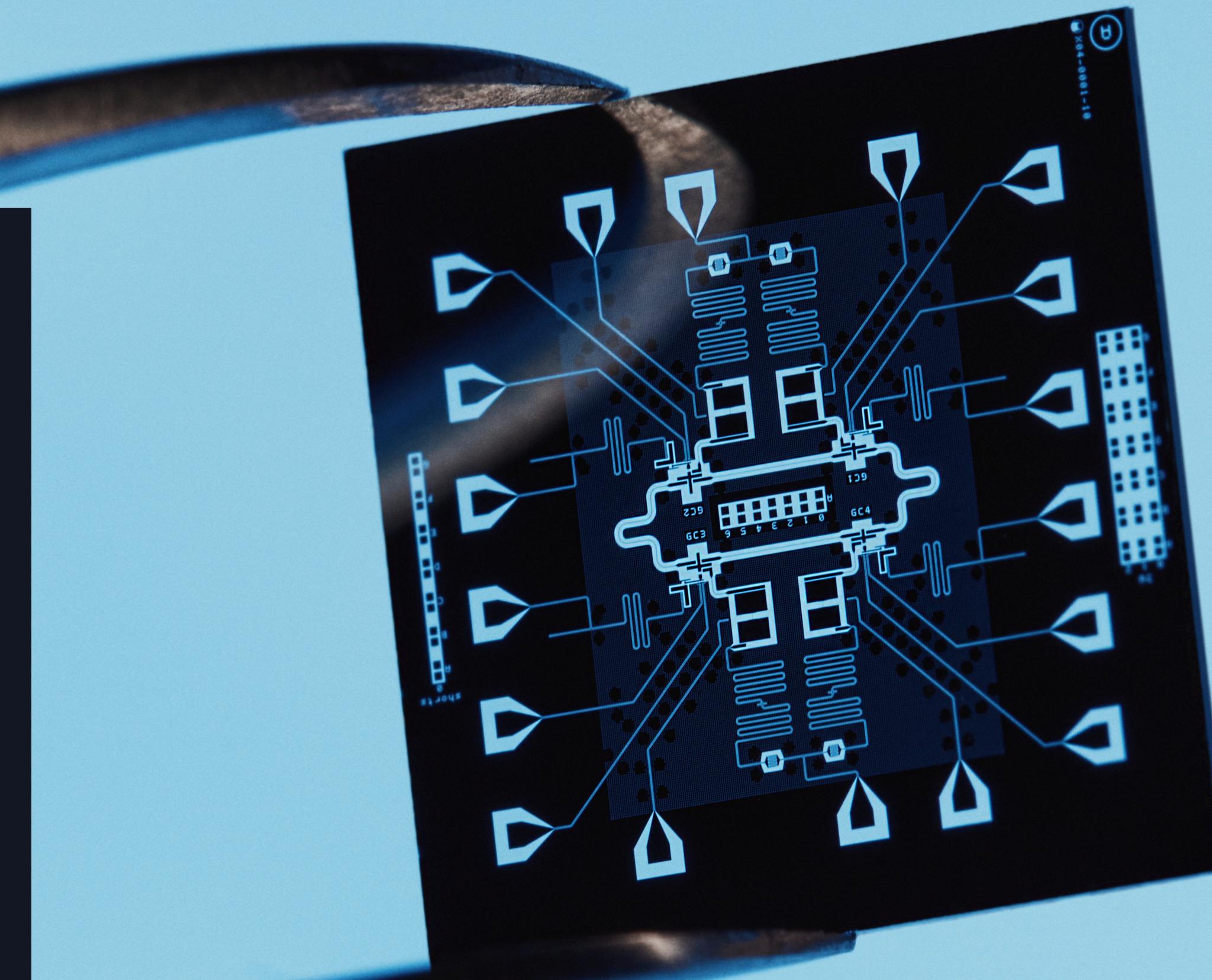
PhD defence / 4th December 2023



Université de
Sherbrooke



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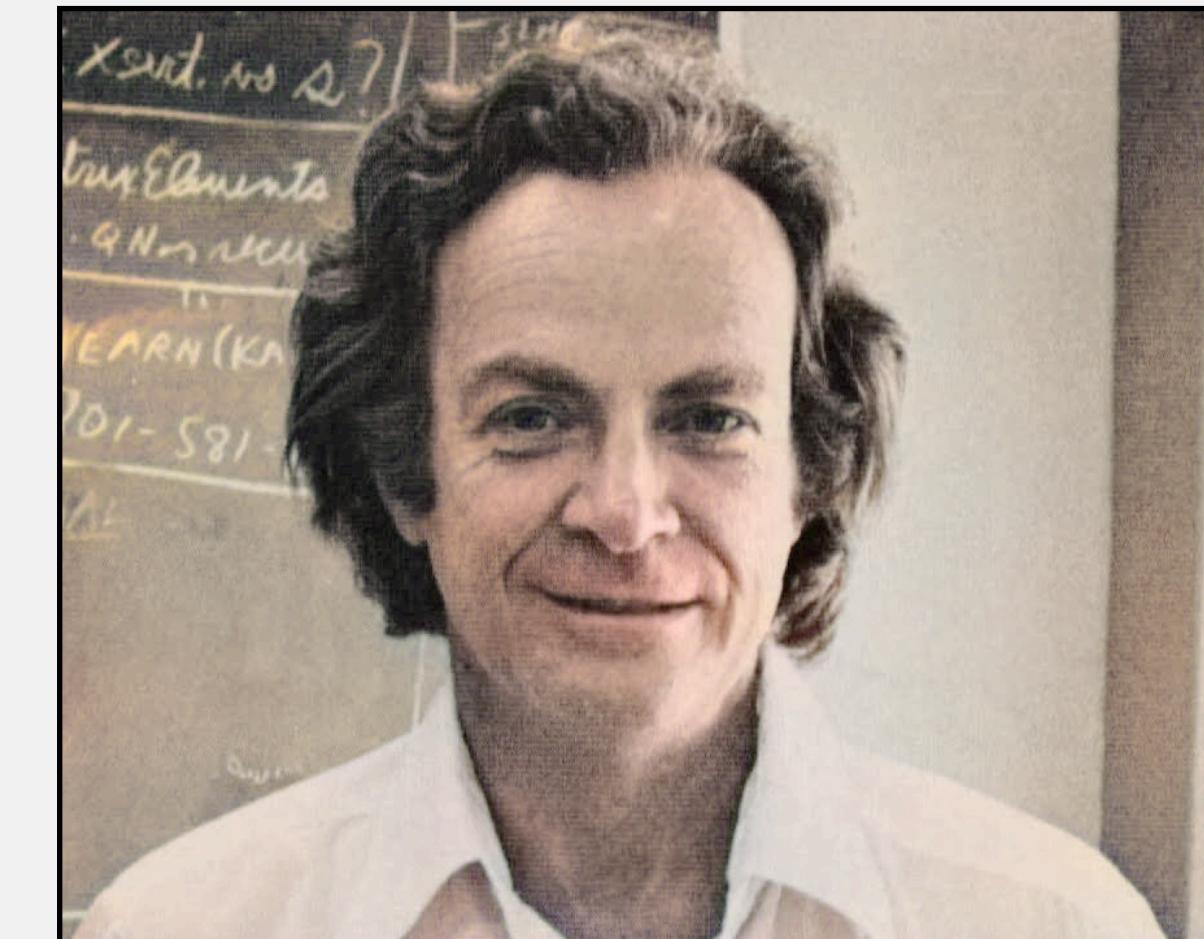


A quantum computer for simulating Nature

Feynman's 1981 talk

“ The full description of quantum mechanics [...] cannot be simulated with a normal computer.

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Richard Feynman at Caltech, circa 1980

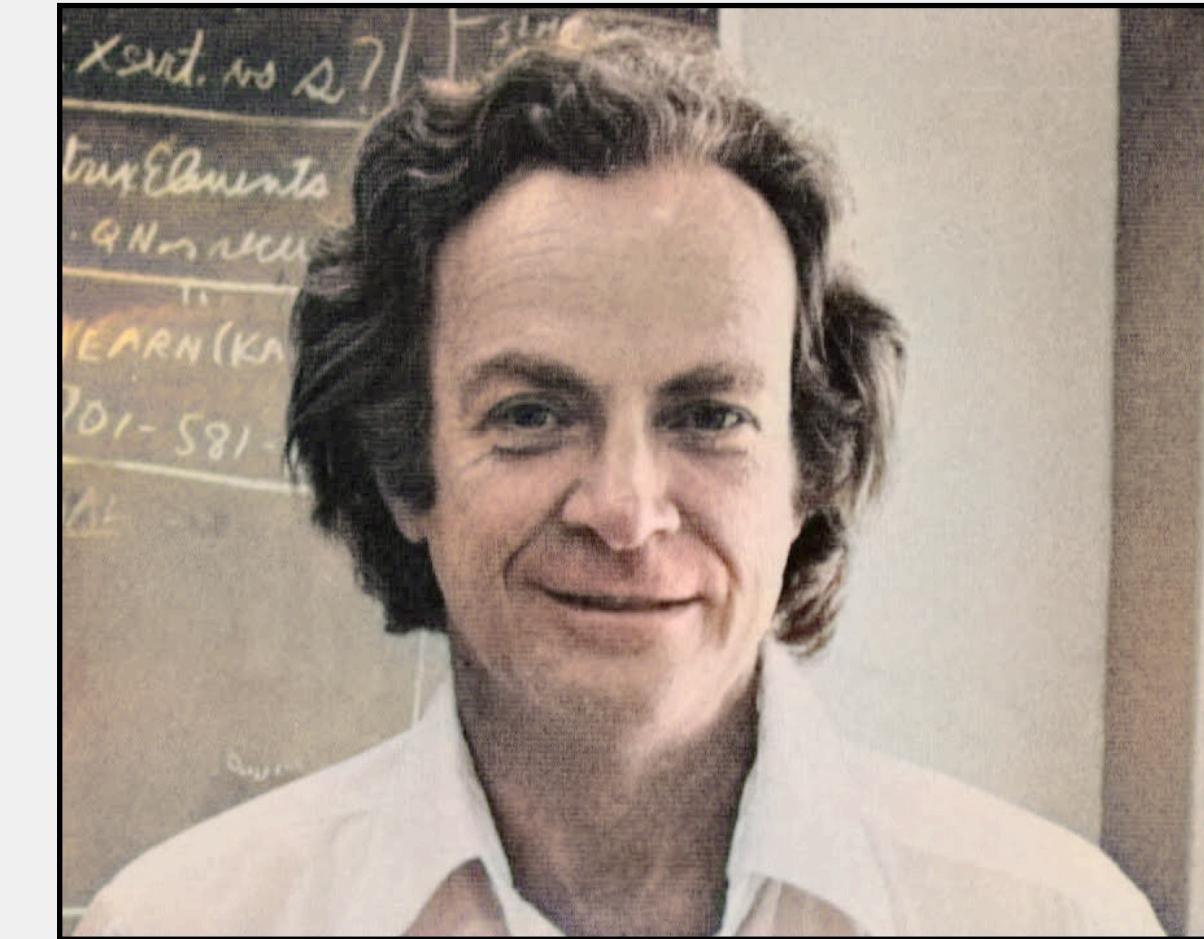
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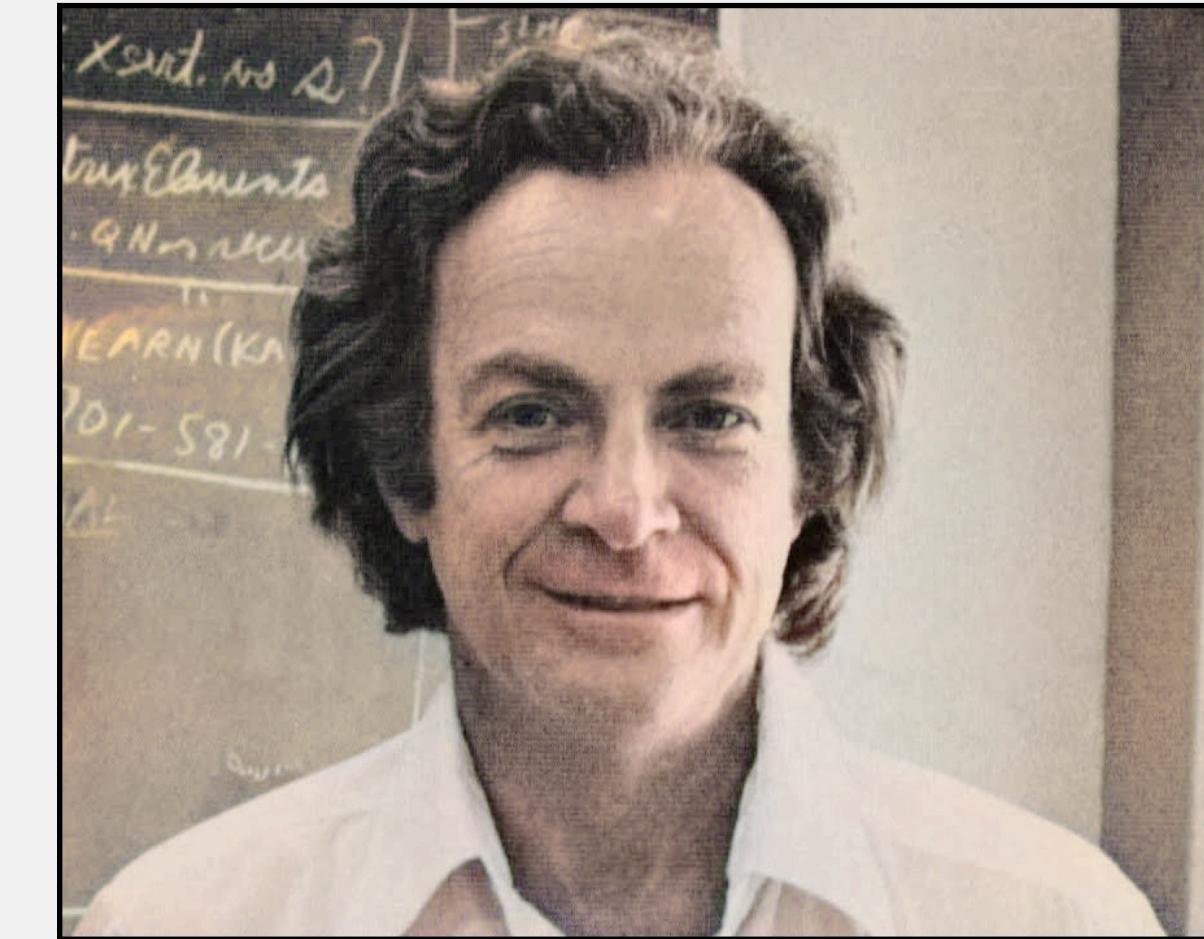
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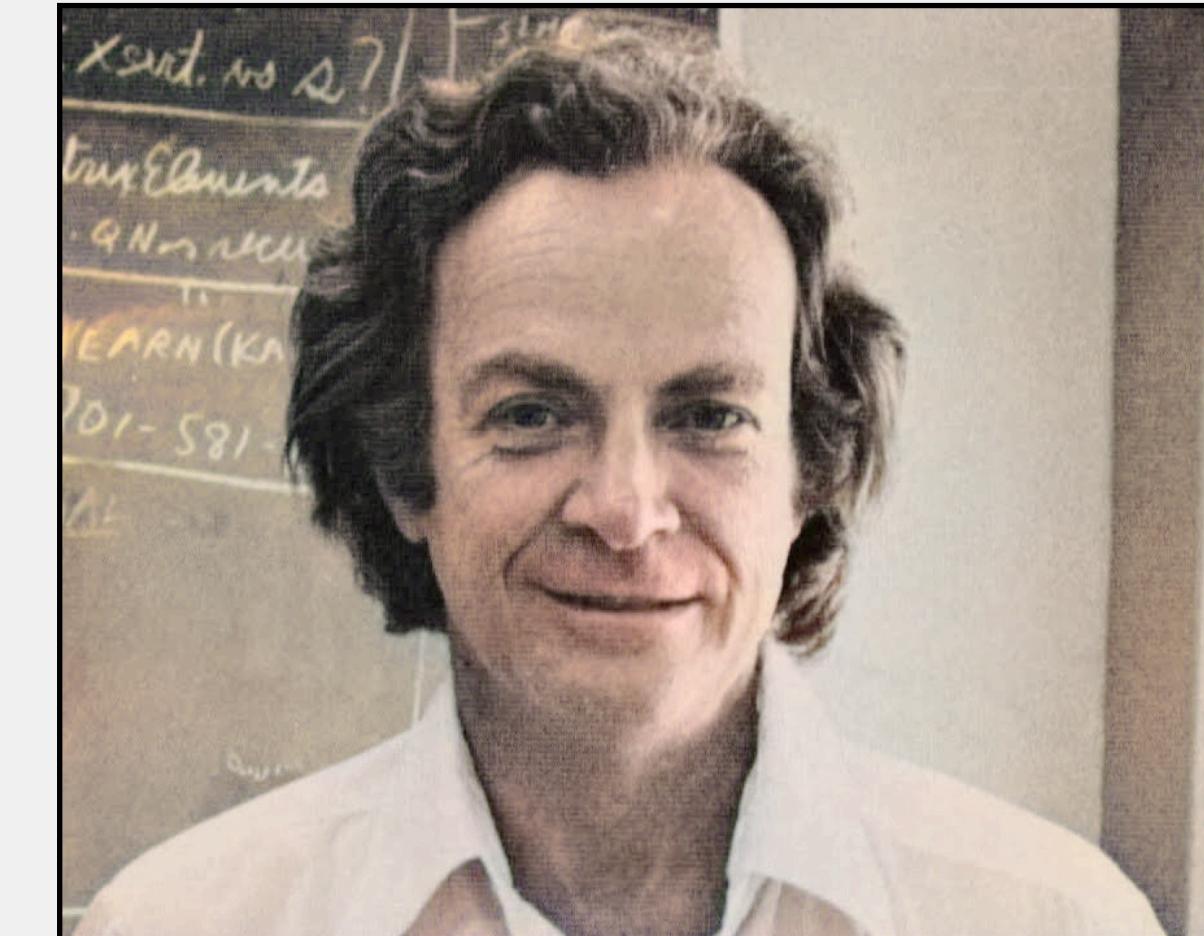
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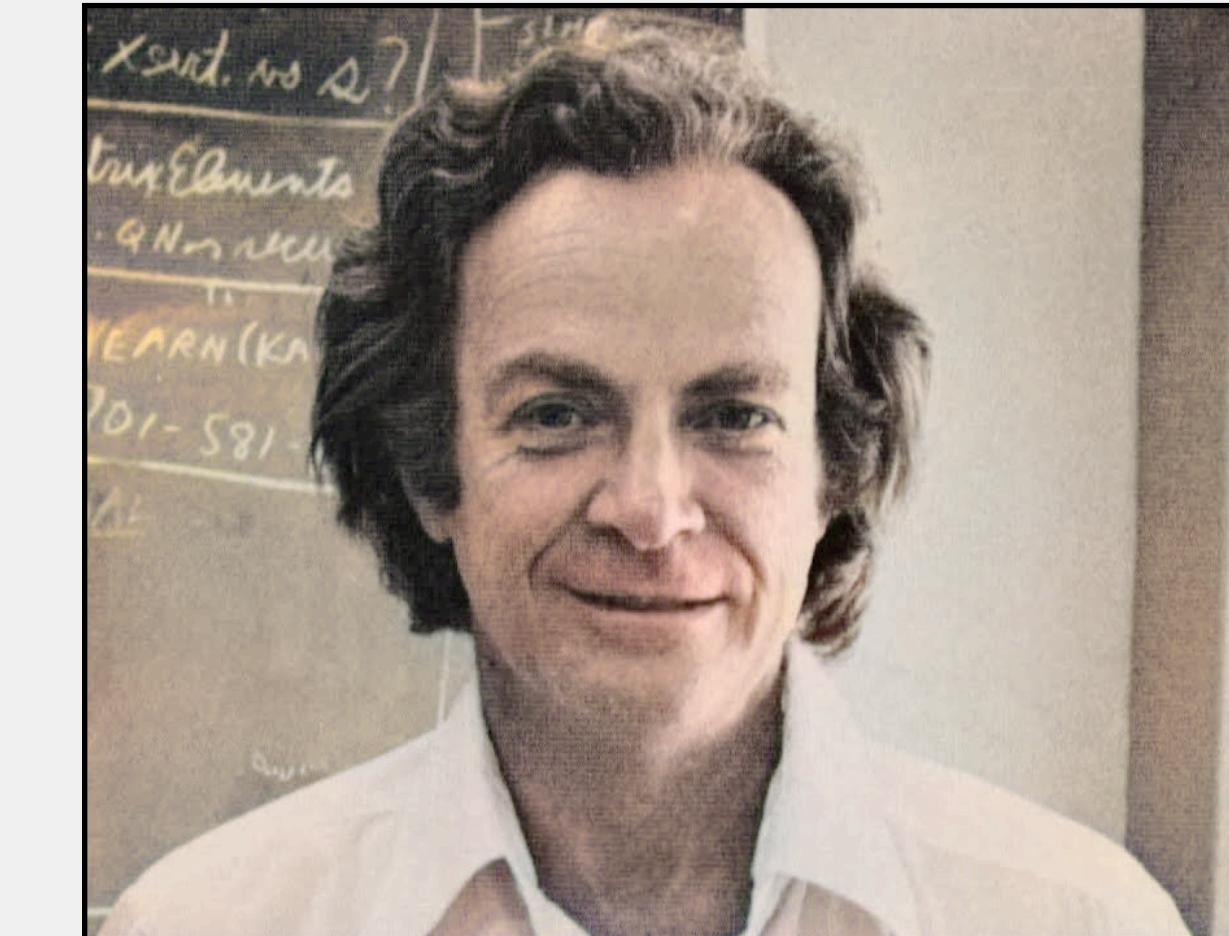
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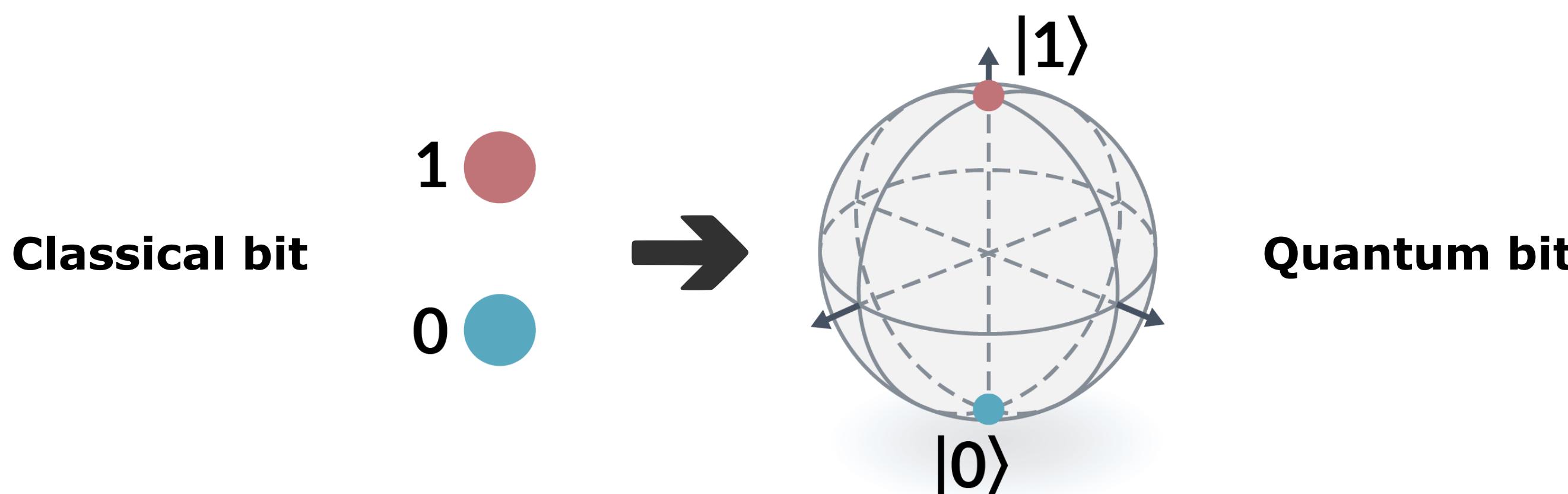
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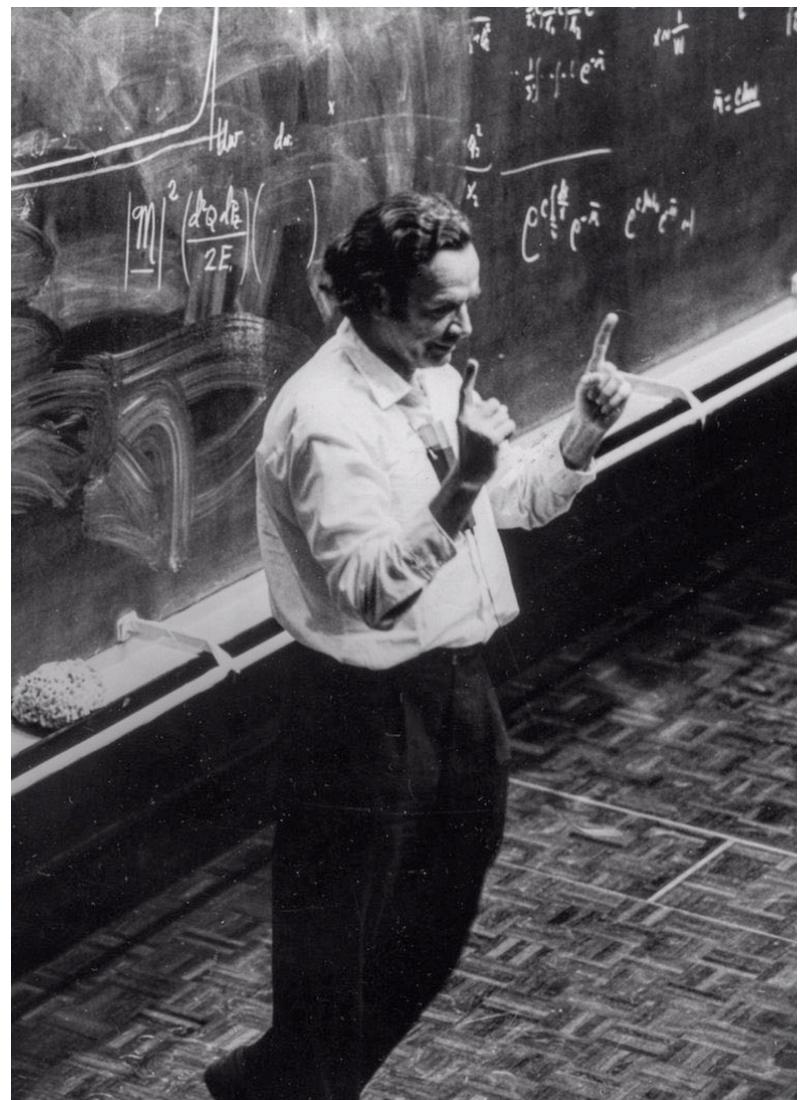


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The development of superconducting circuits

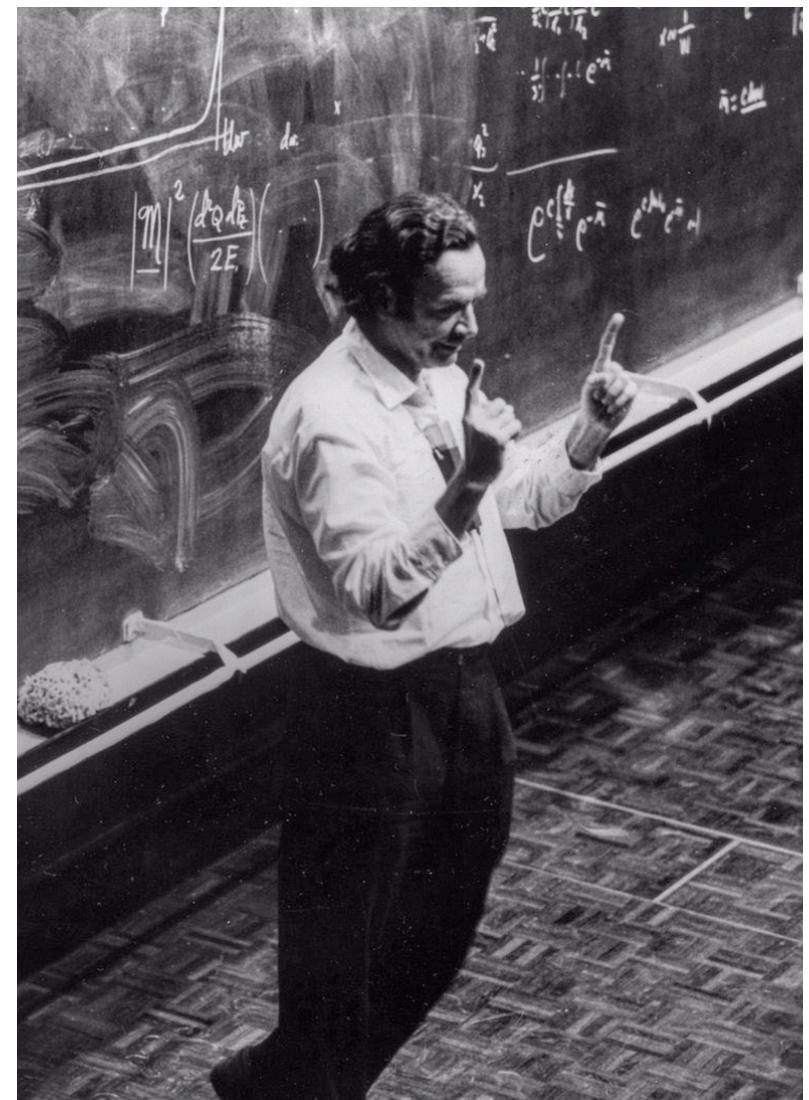
Feynman's talk on QC



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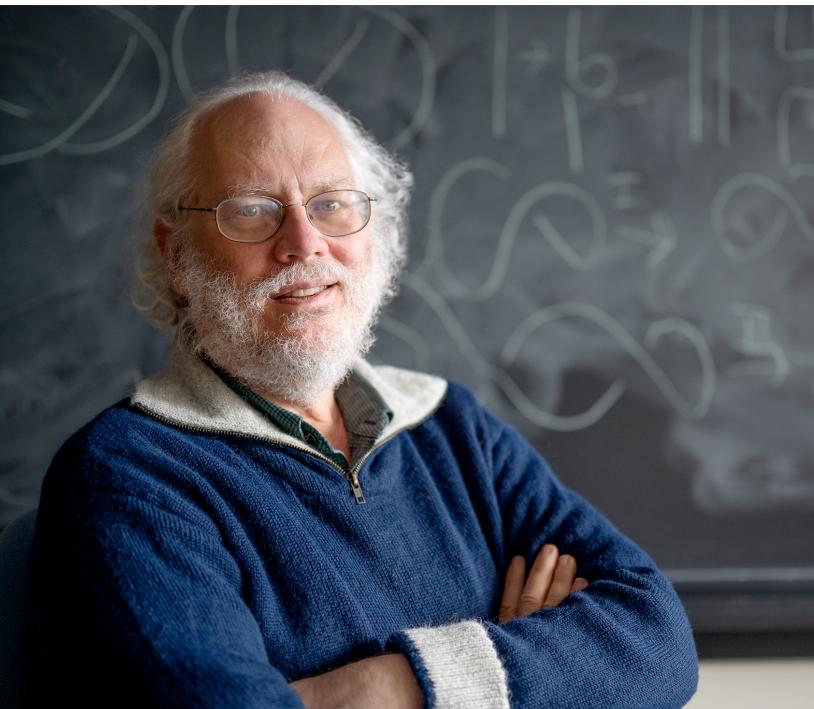
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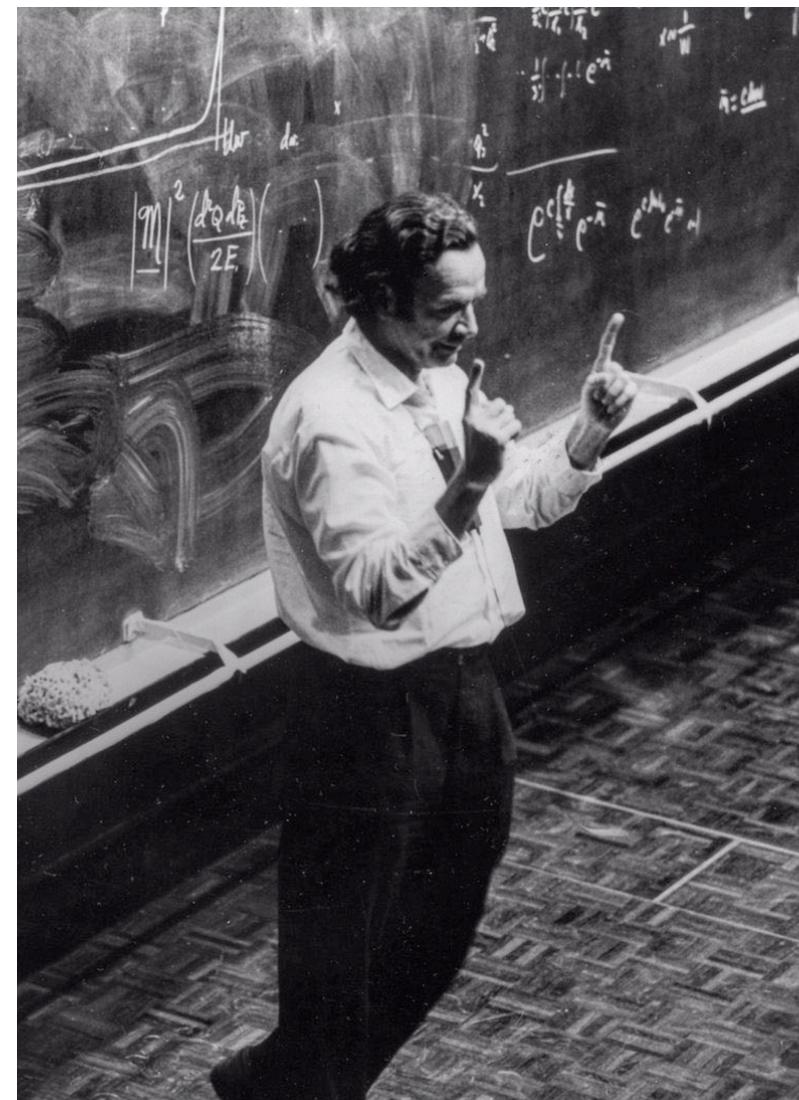
Shor's algorithm



1994

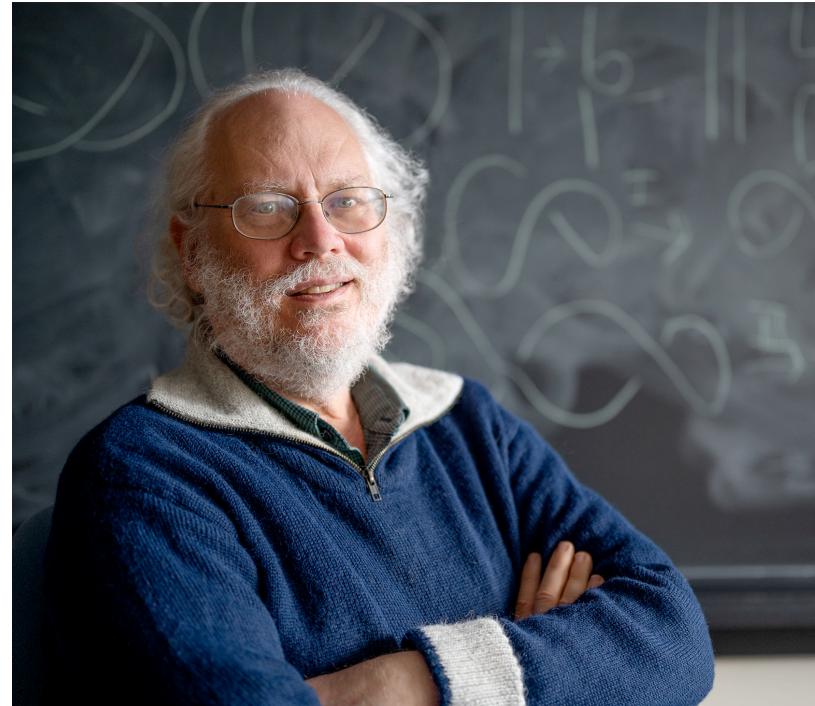
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Quantum Error Correction

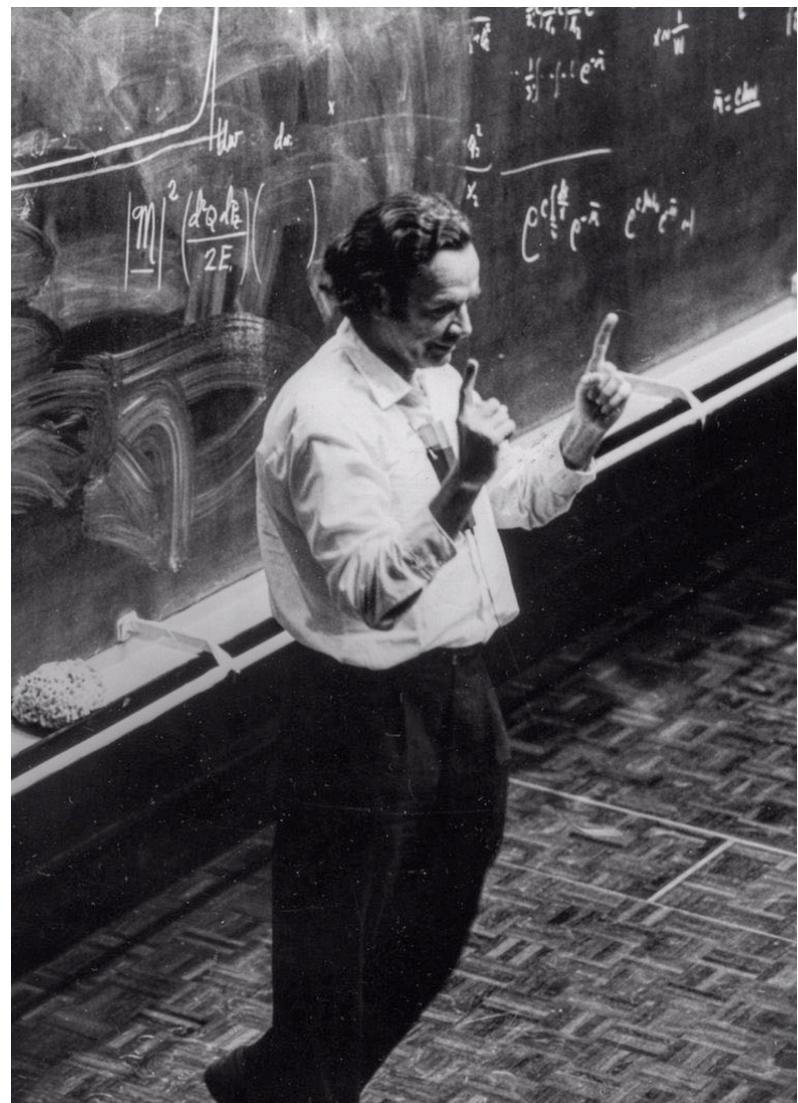


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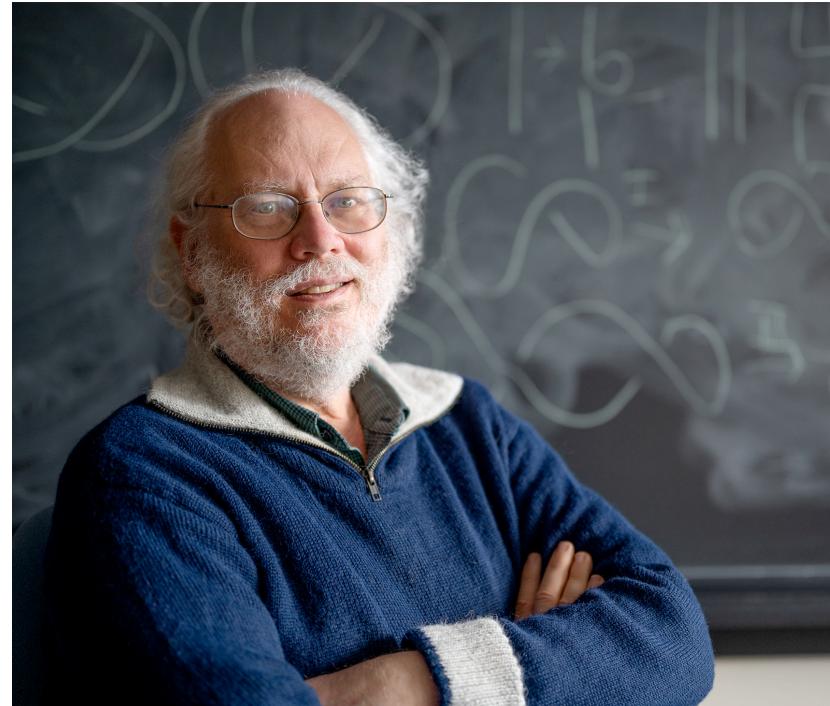
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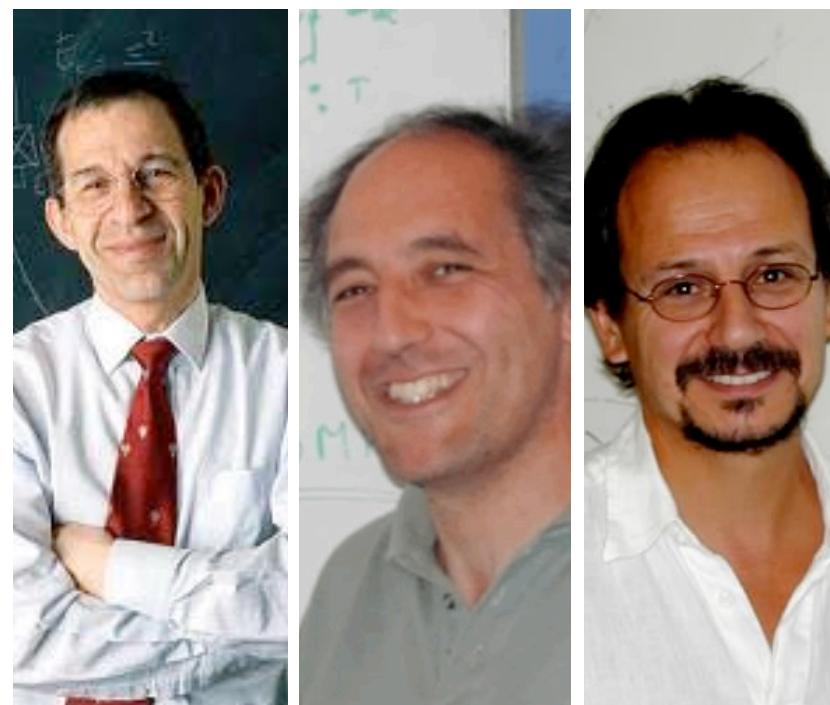
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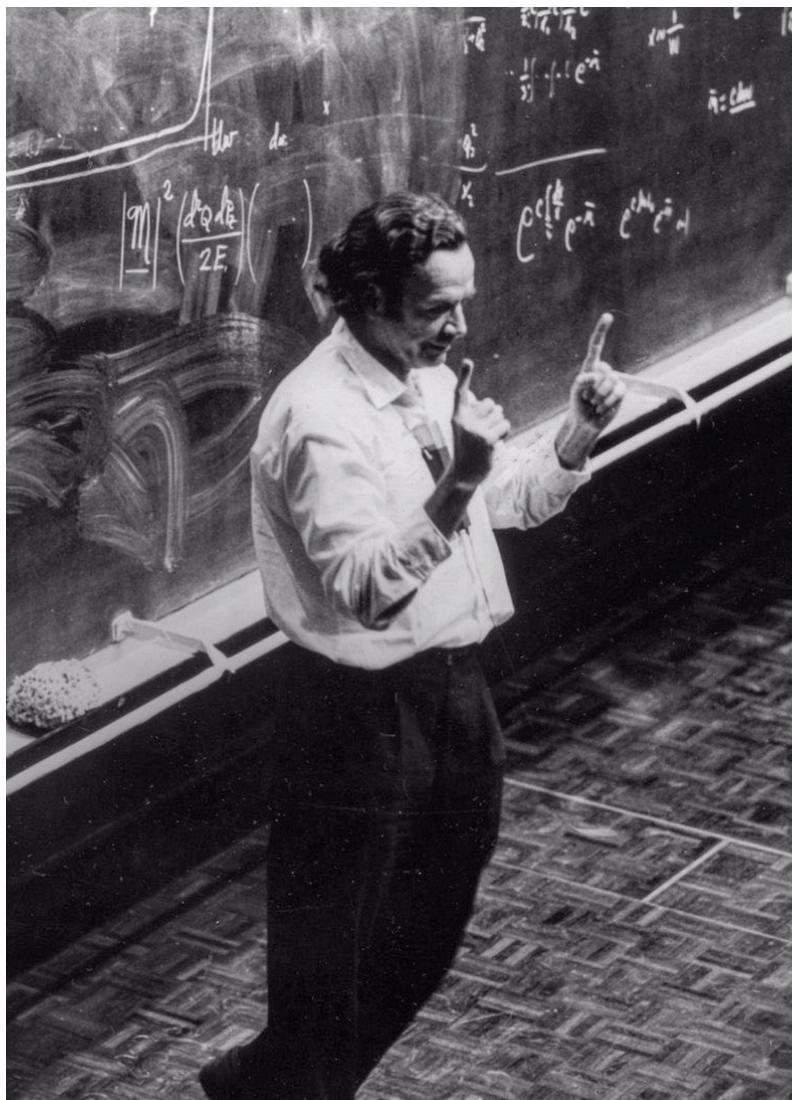
Creation of Quantronics



1985

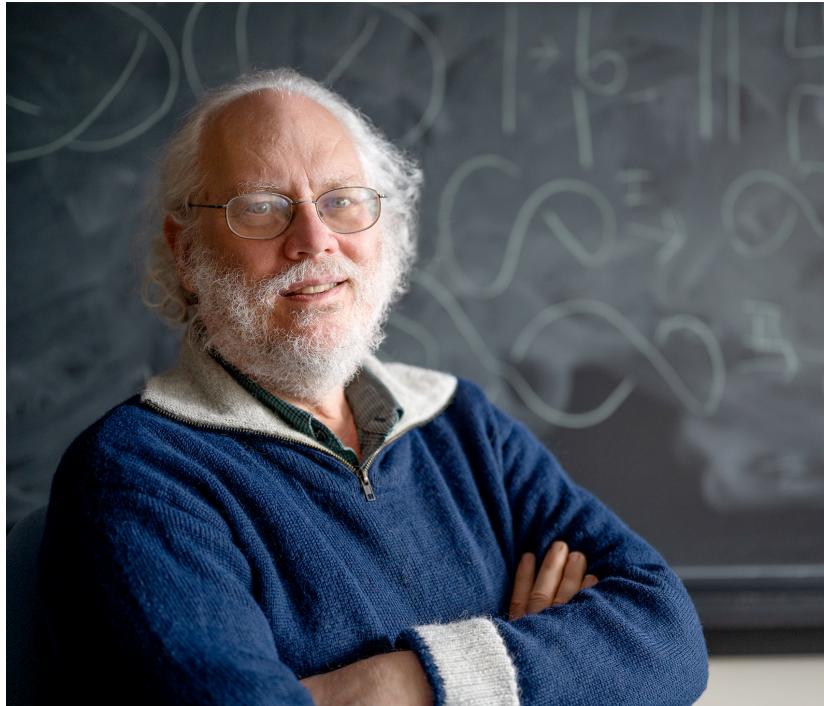
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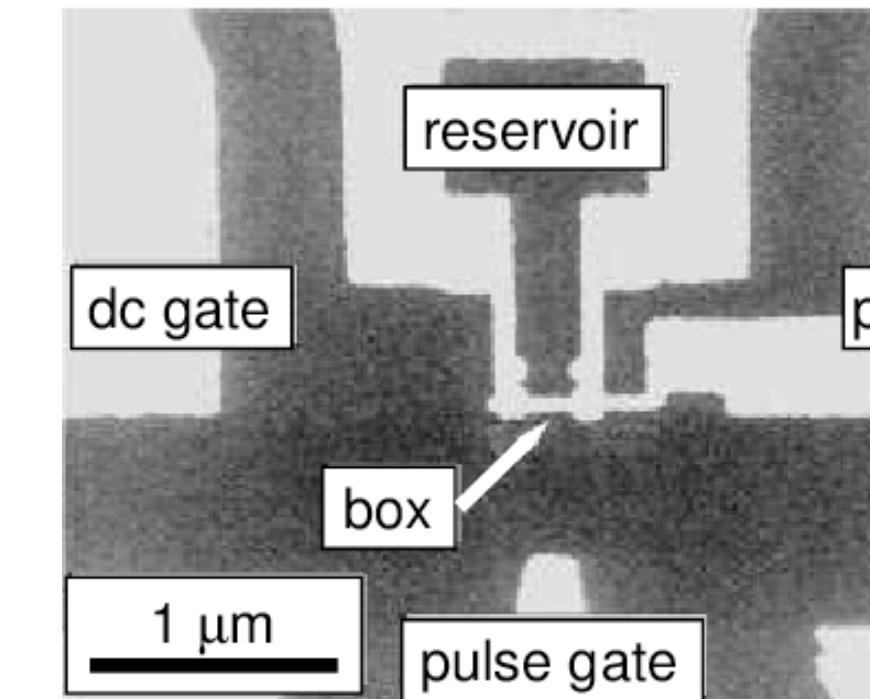
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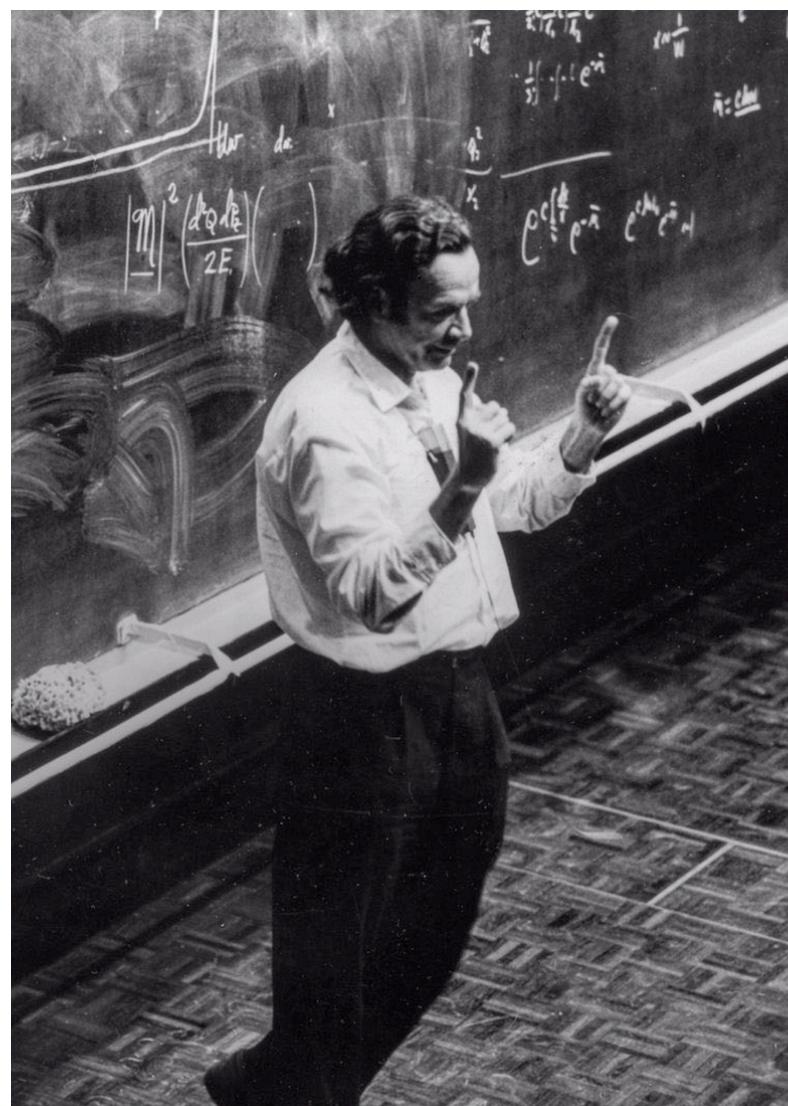
Cooper Pair Box



1999

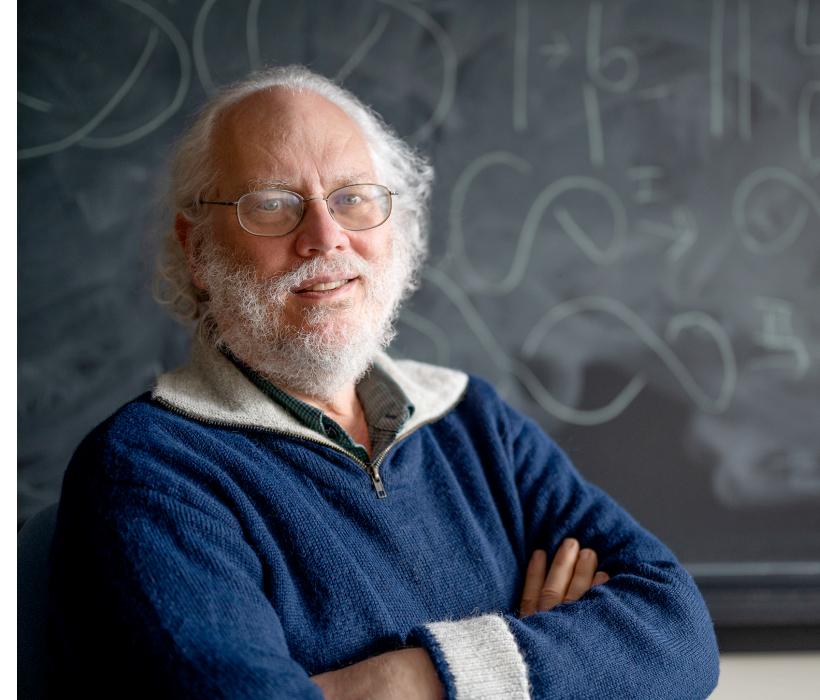
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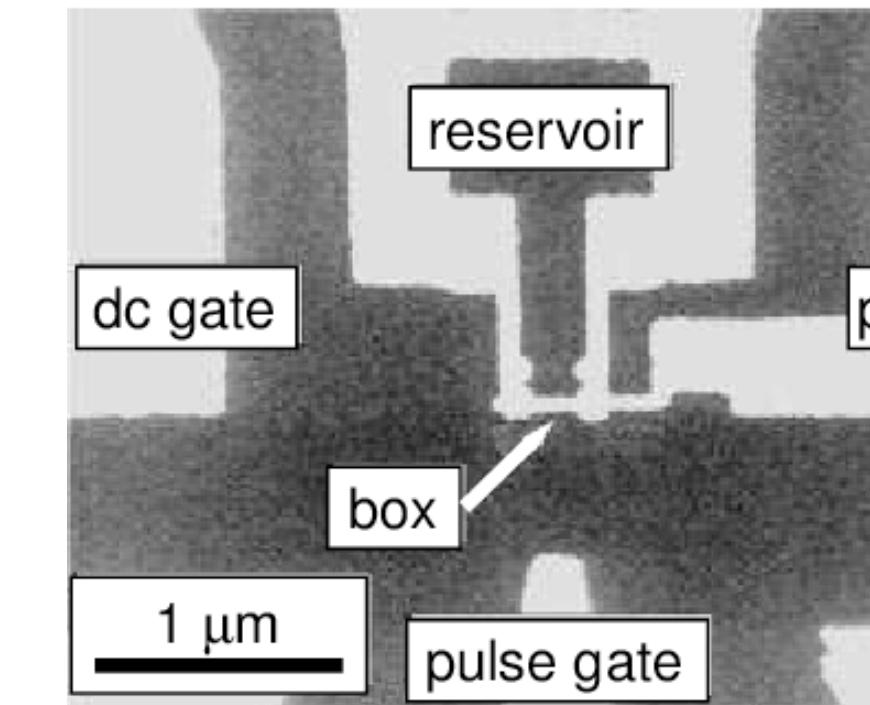
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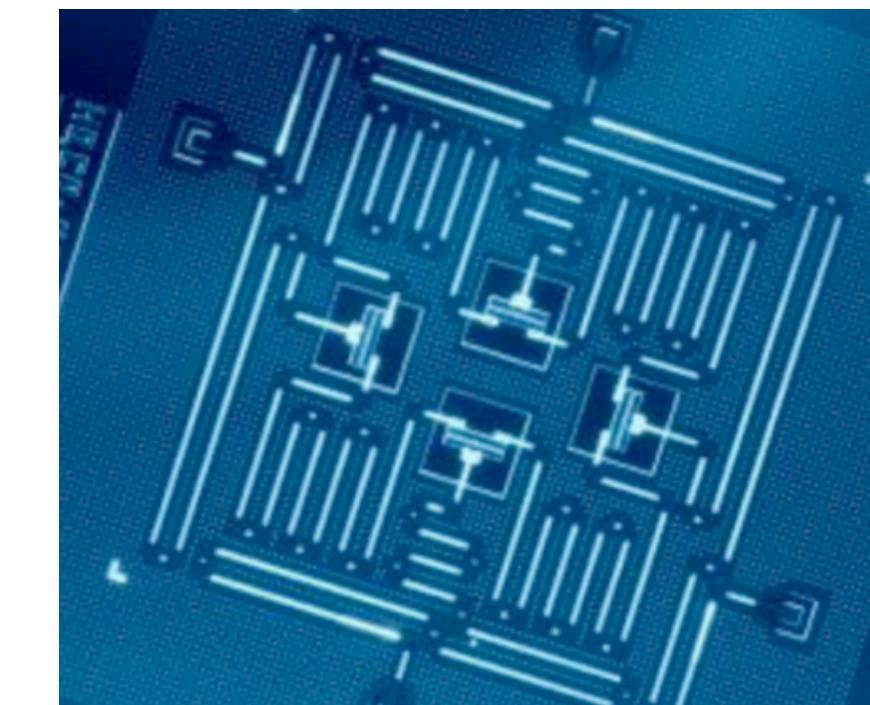
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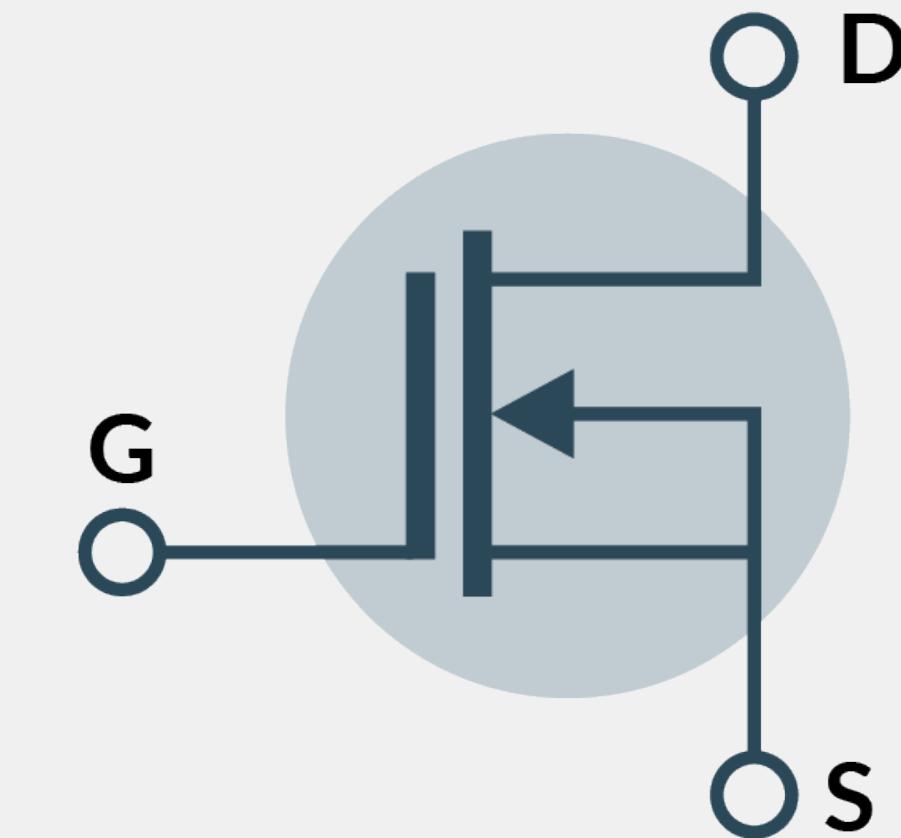
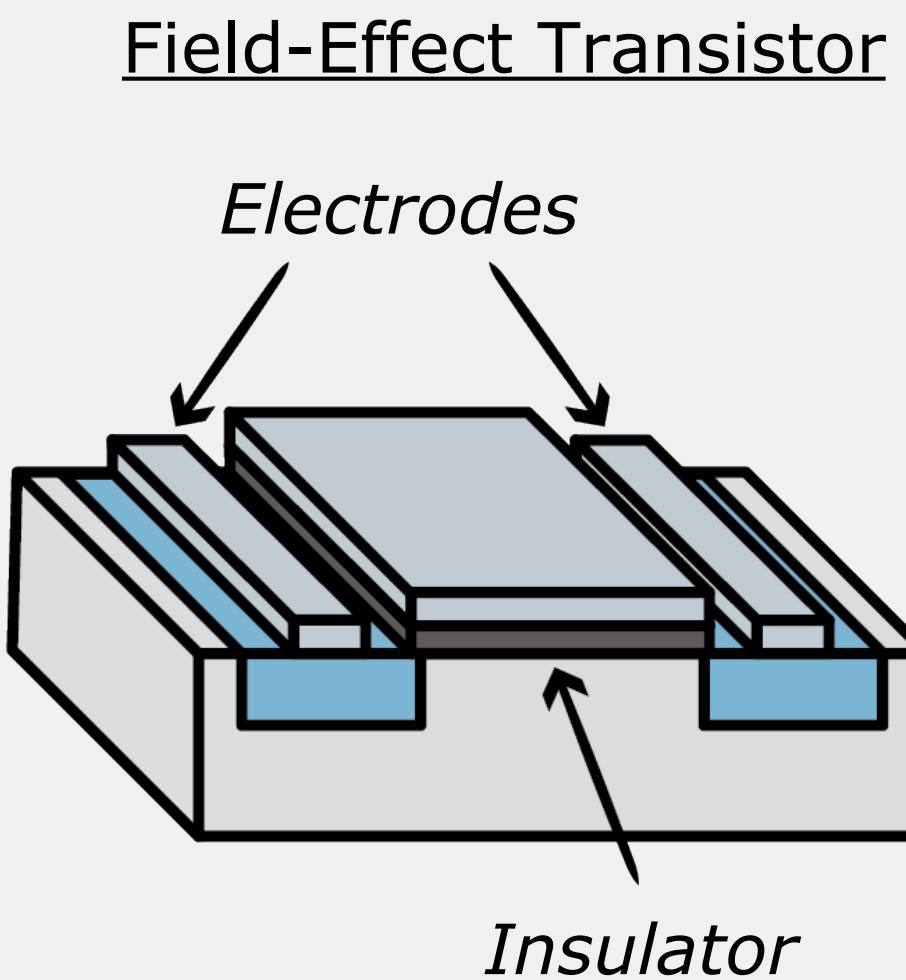


2007

From transistors to transmons

Transistor

1
0

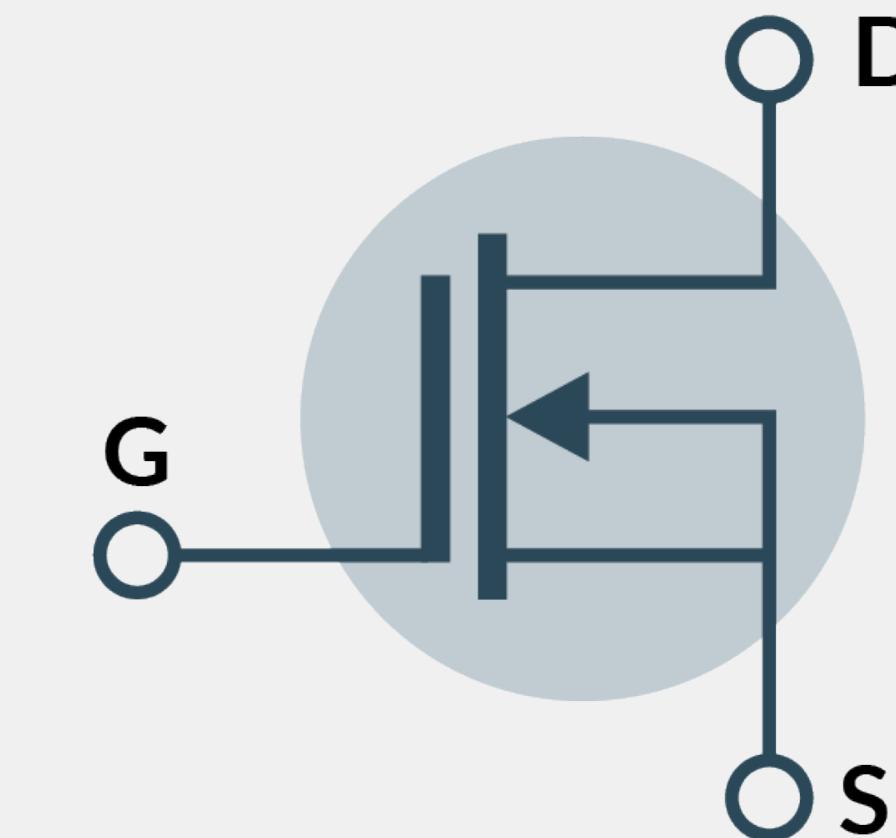
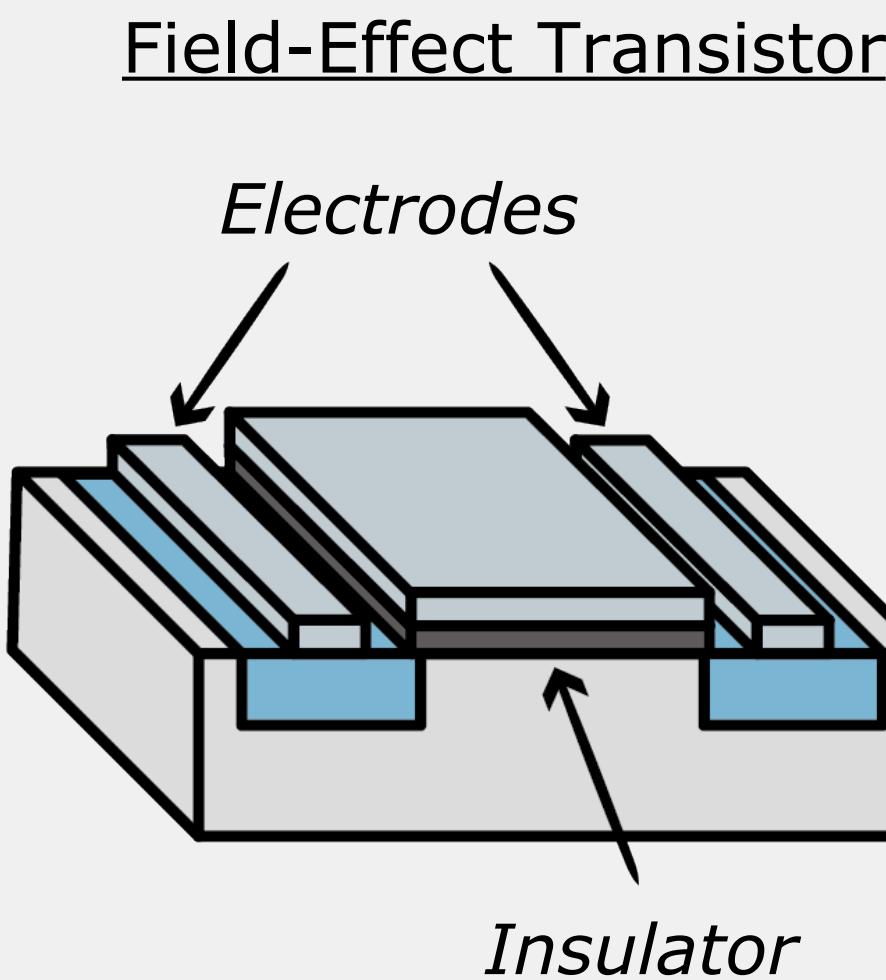


Error per operation
 $\sim 10^{-20} - 10^{-22}$

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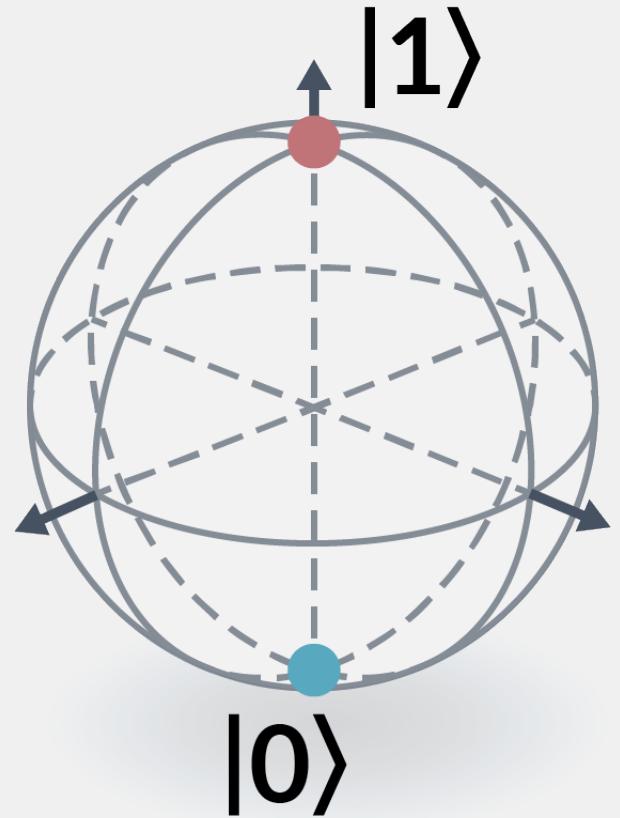
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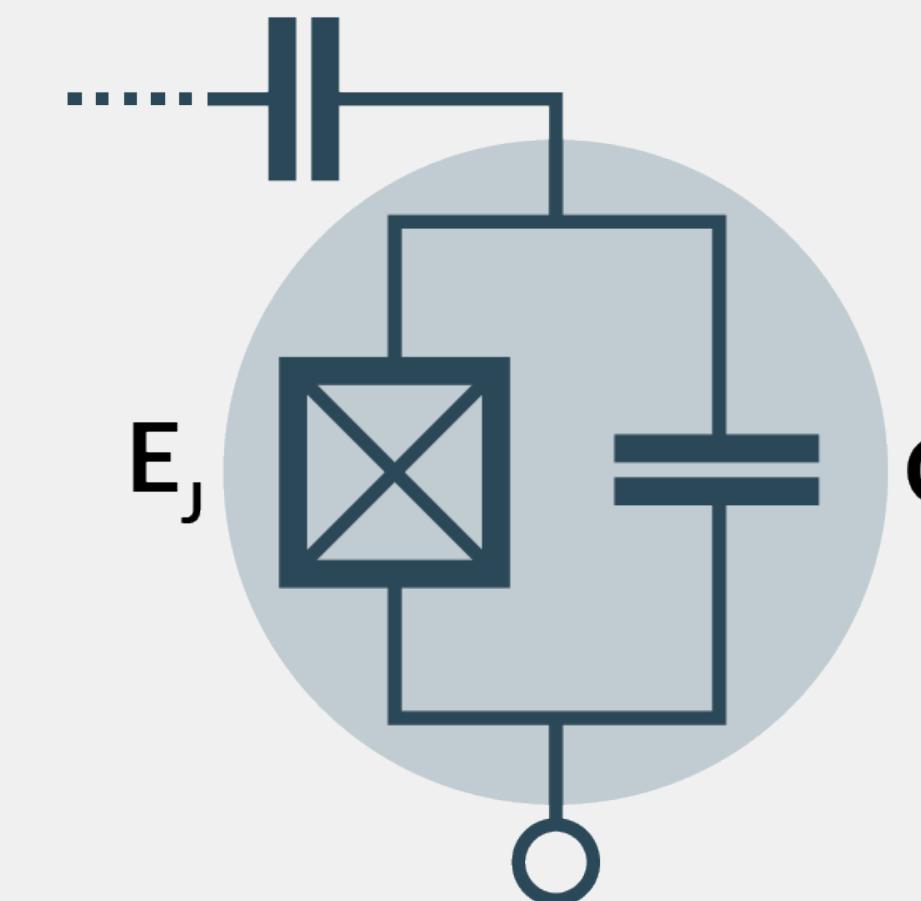
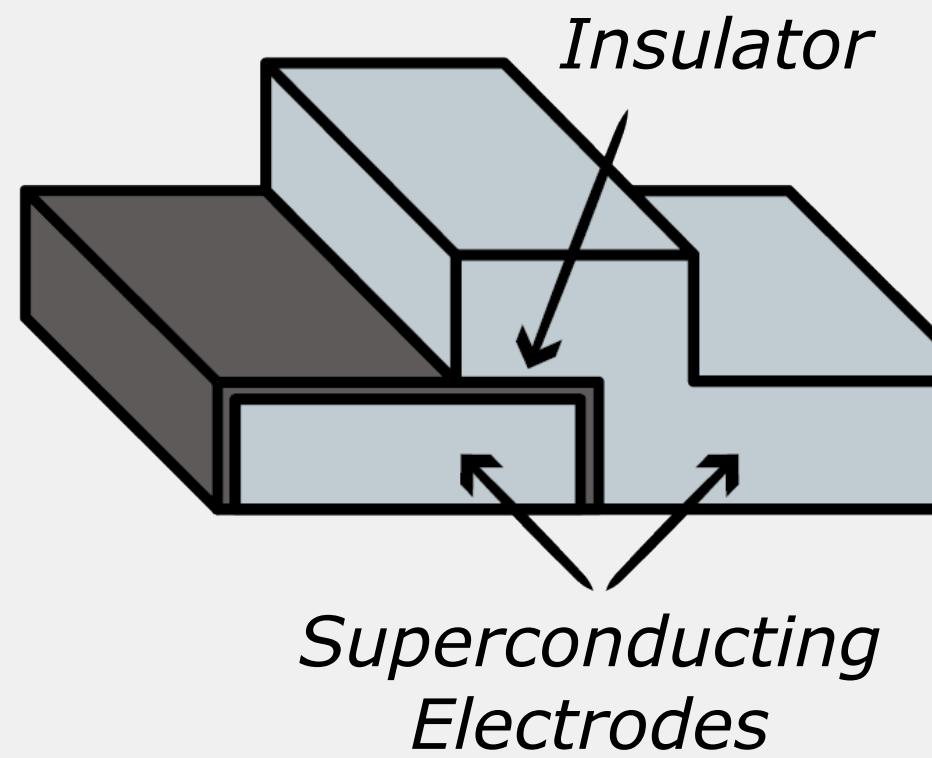


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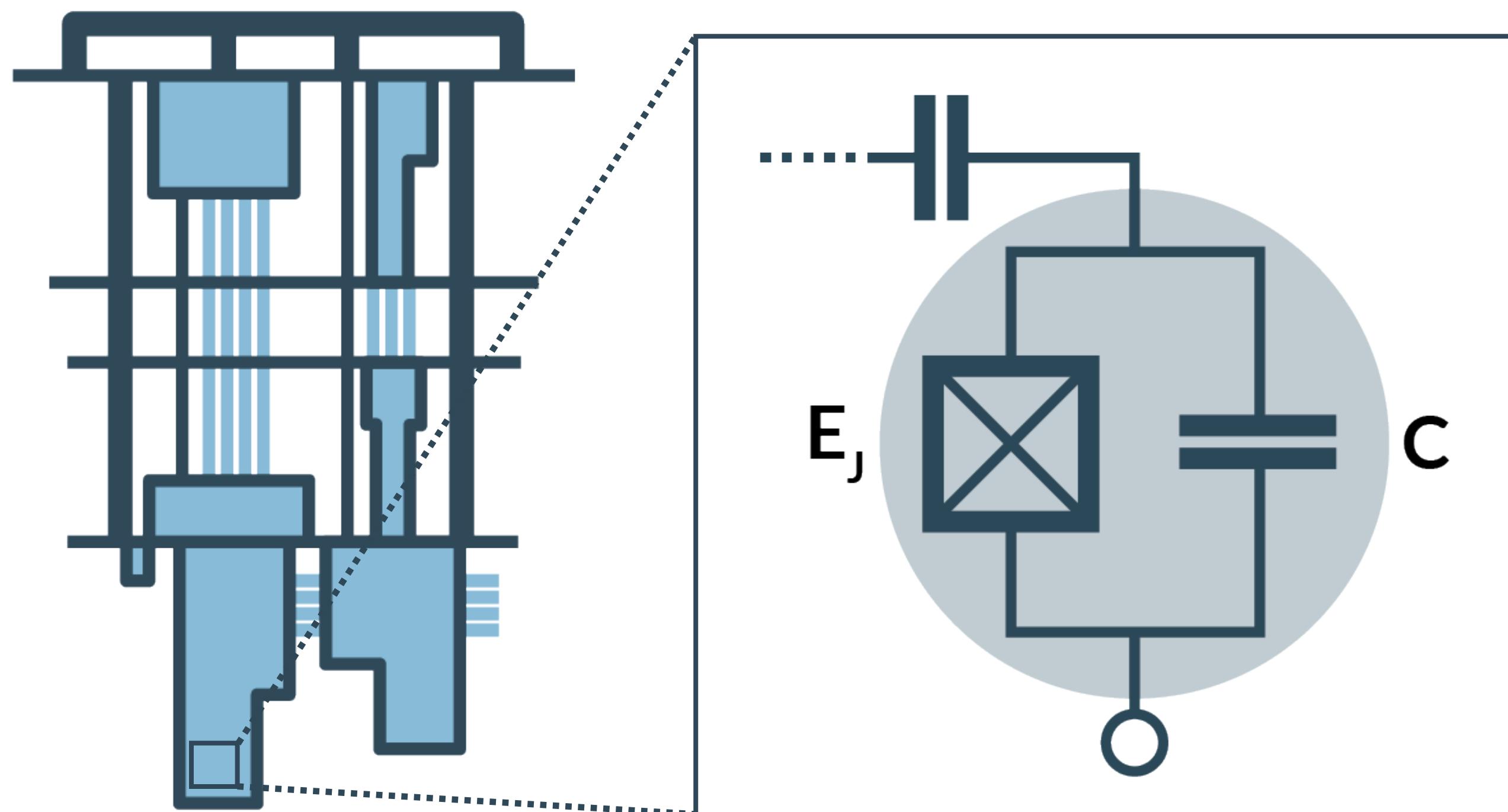
Josephson Junction



Error per operation
 $\sim 10^{-2} - 10^{-4}$

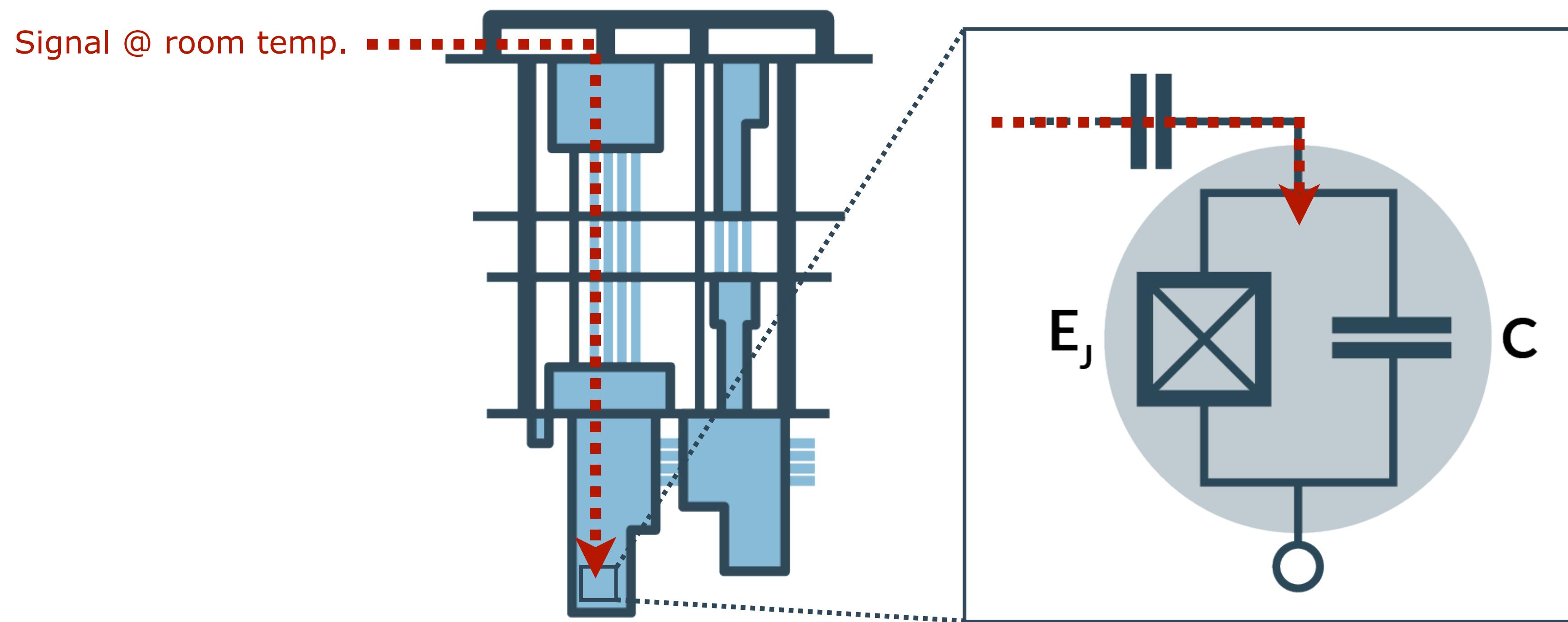
A fundamental predicament

High controllability \longleftrightarrow Long lifetime



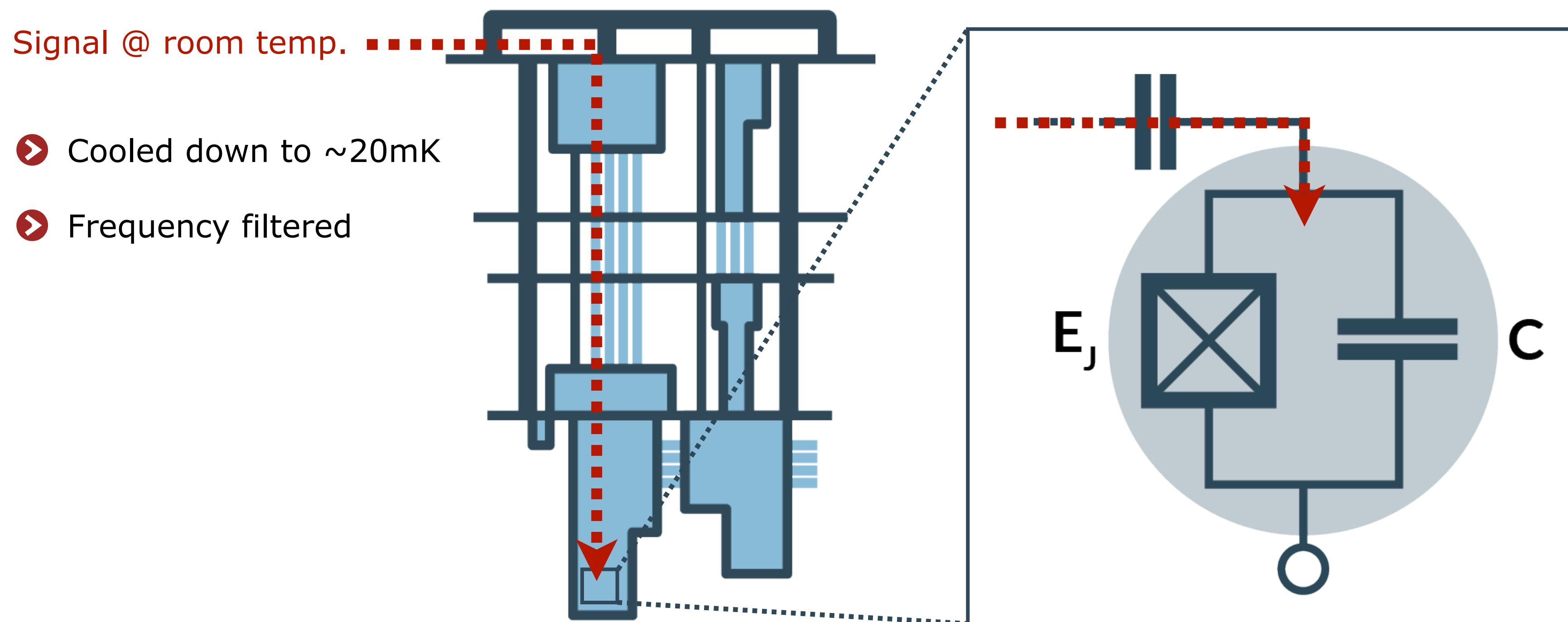
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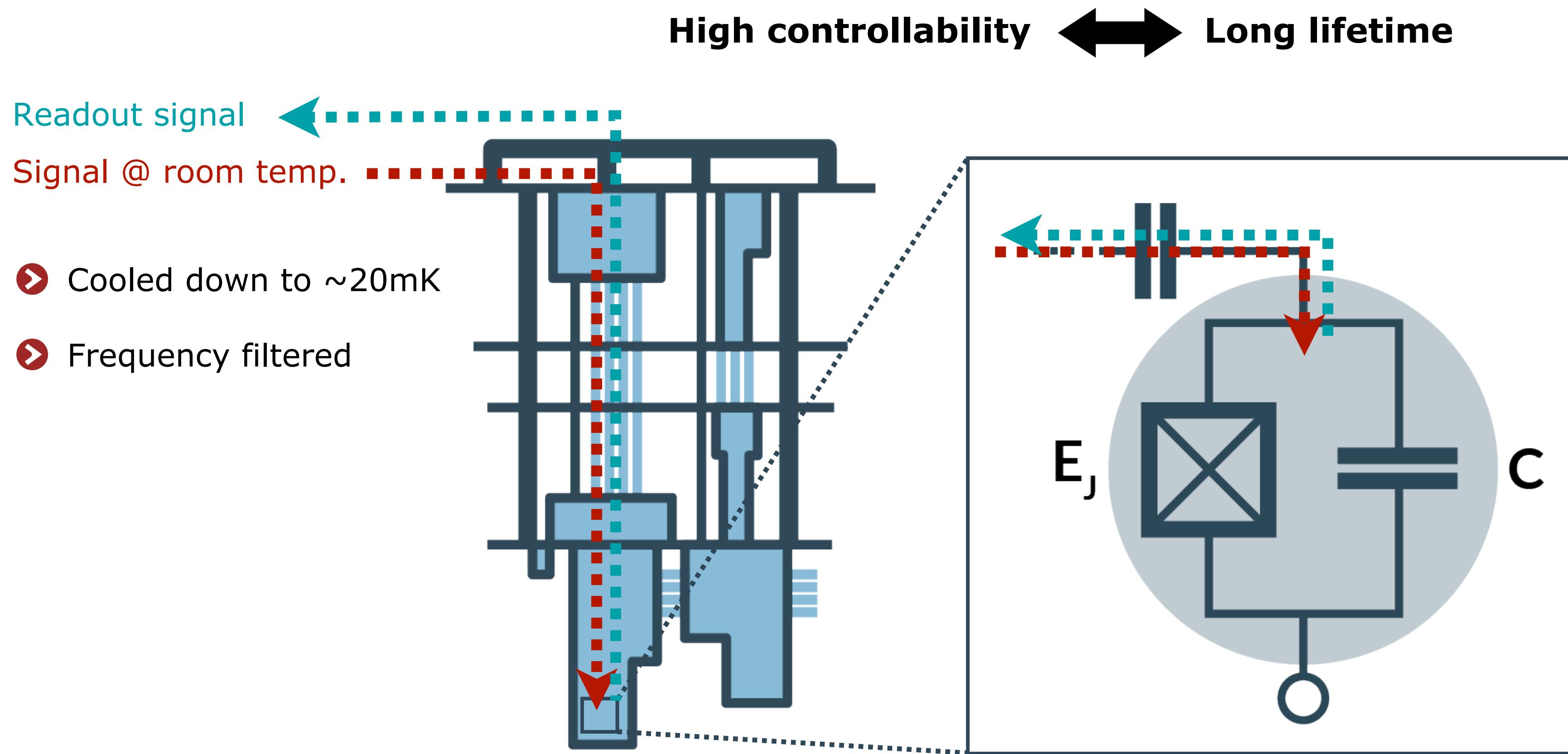


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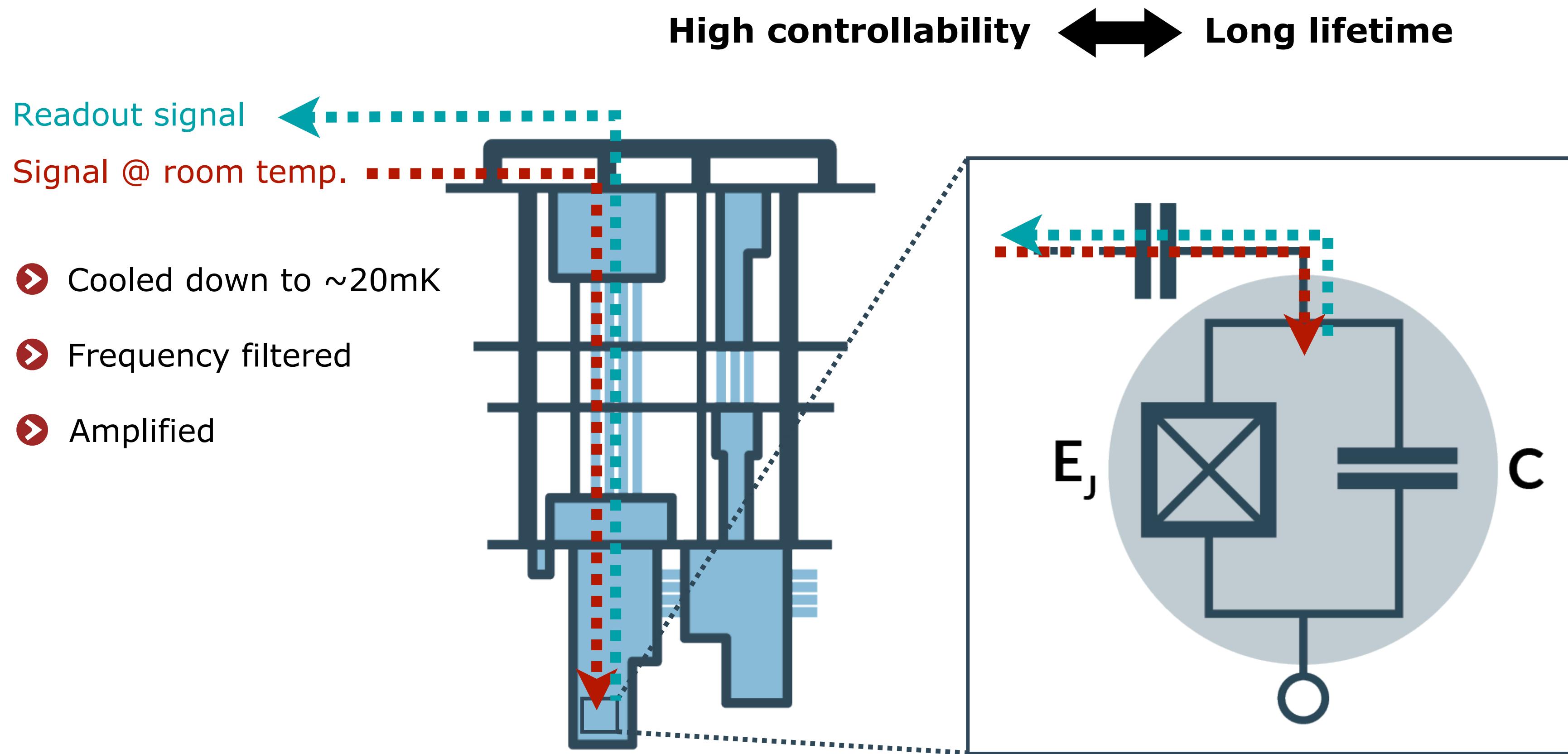
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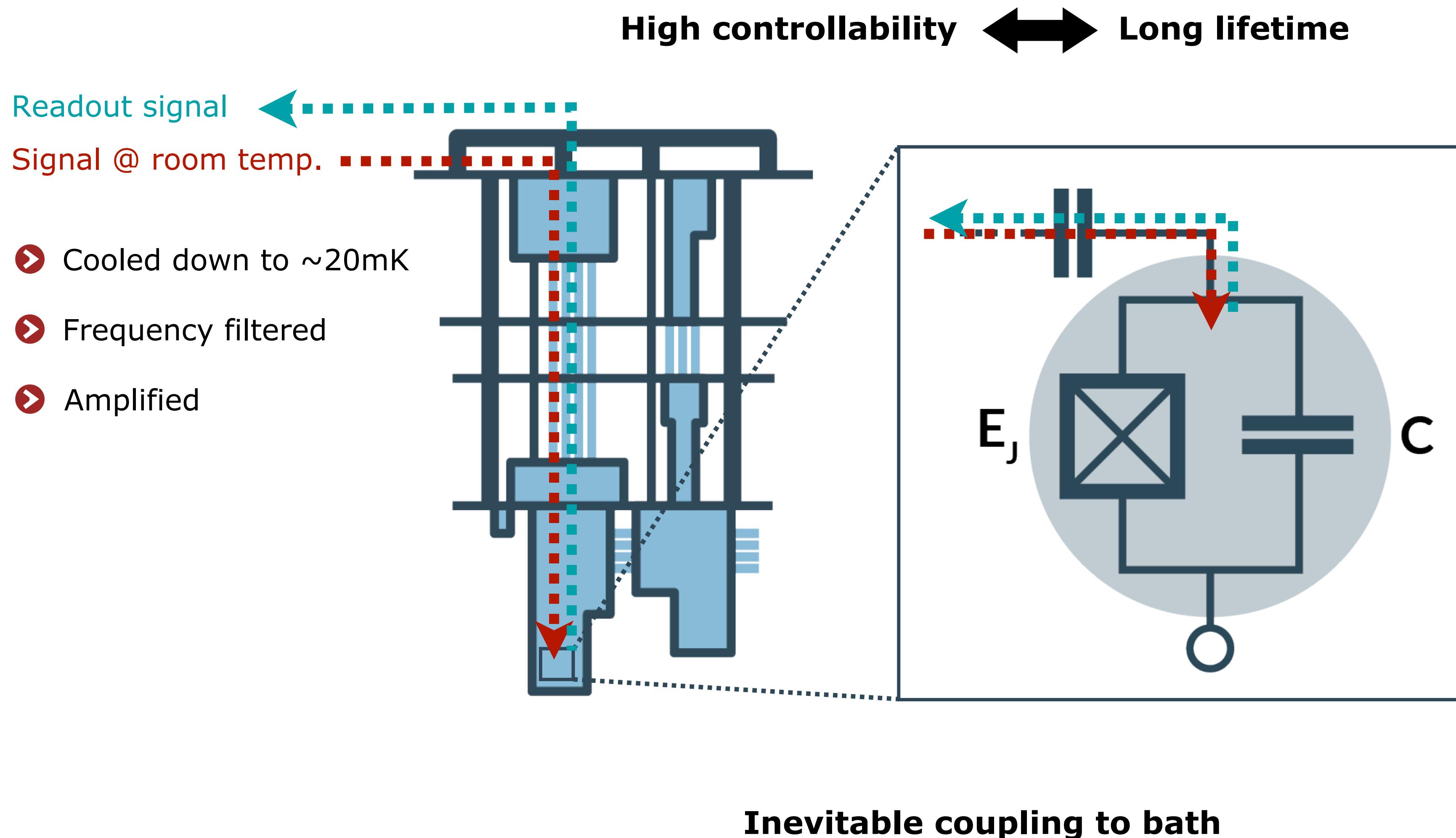
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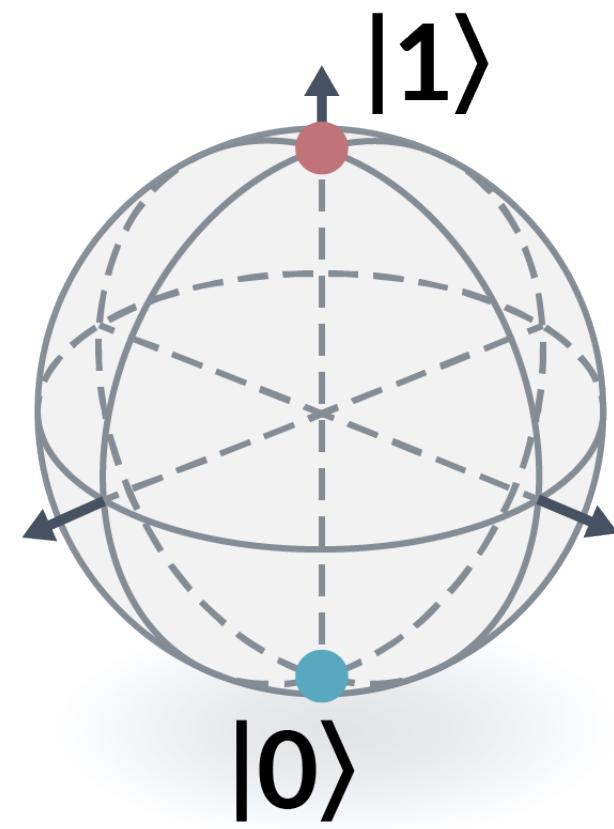
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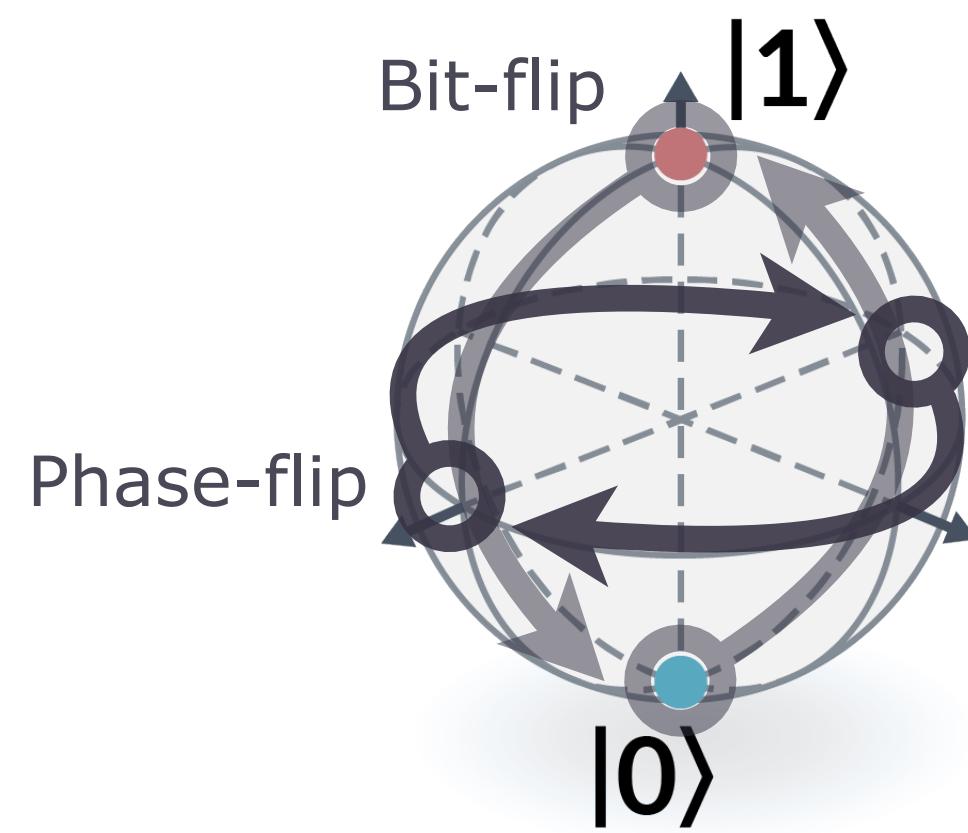
Quantum error correction



Error discretisation theorem

Correcting Pauli errors = correcting arbitrary errors

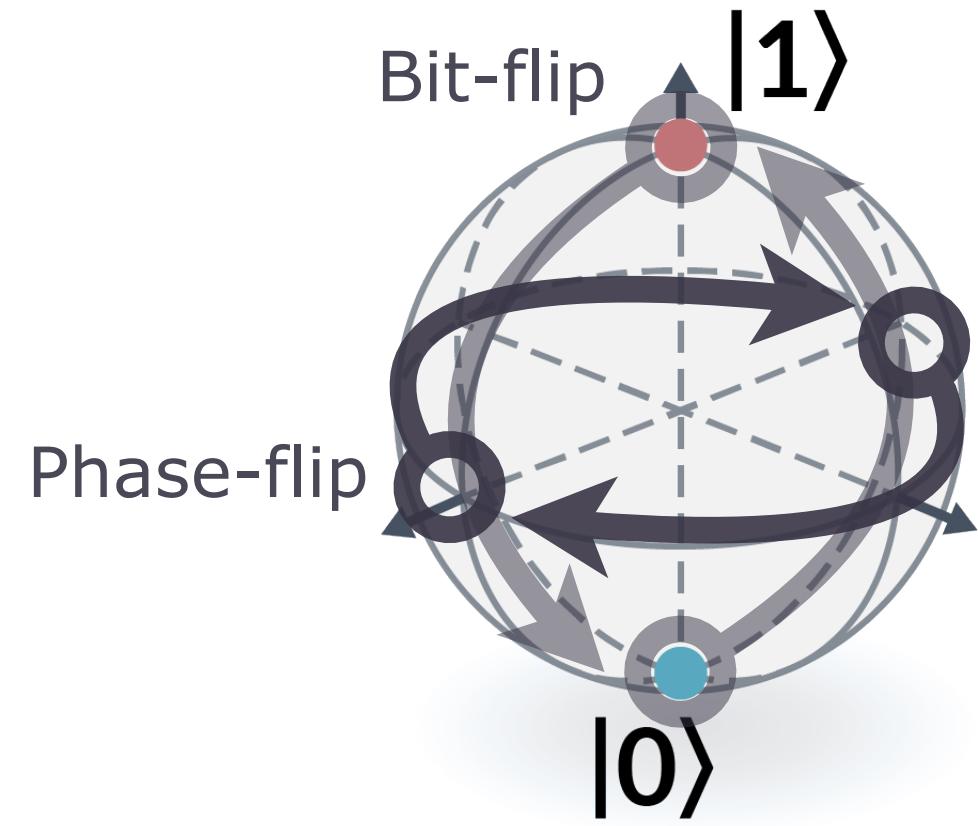
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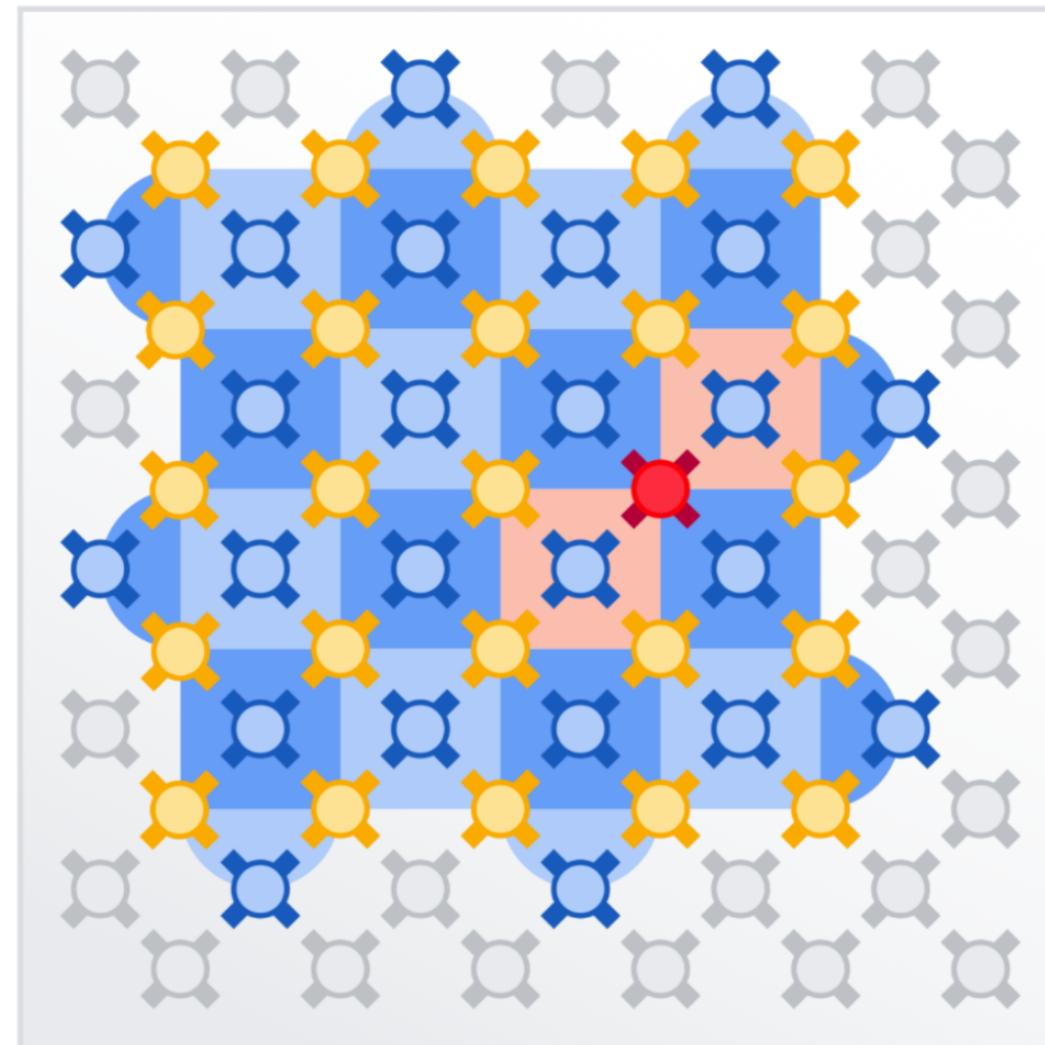
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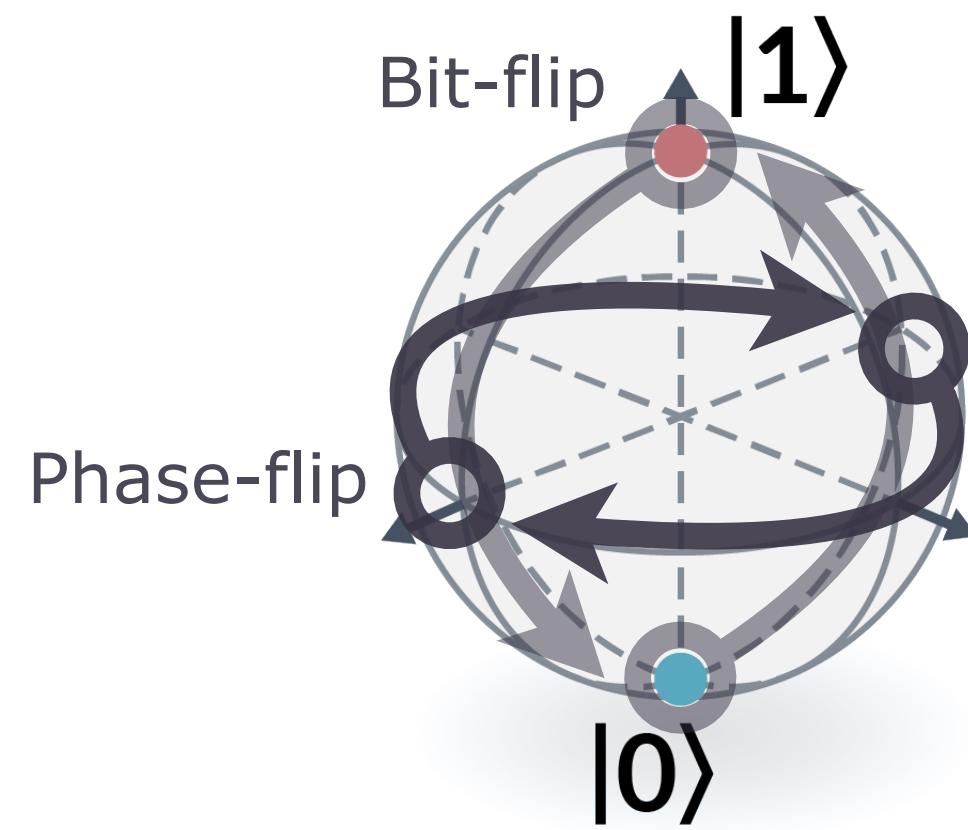
Discrete qubit codes



■ Data qubit ■ Measure qubit ■ Data qubit with error

Google Quantum AI, Nature 2022

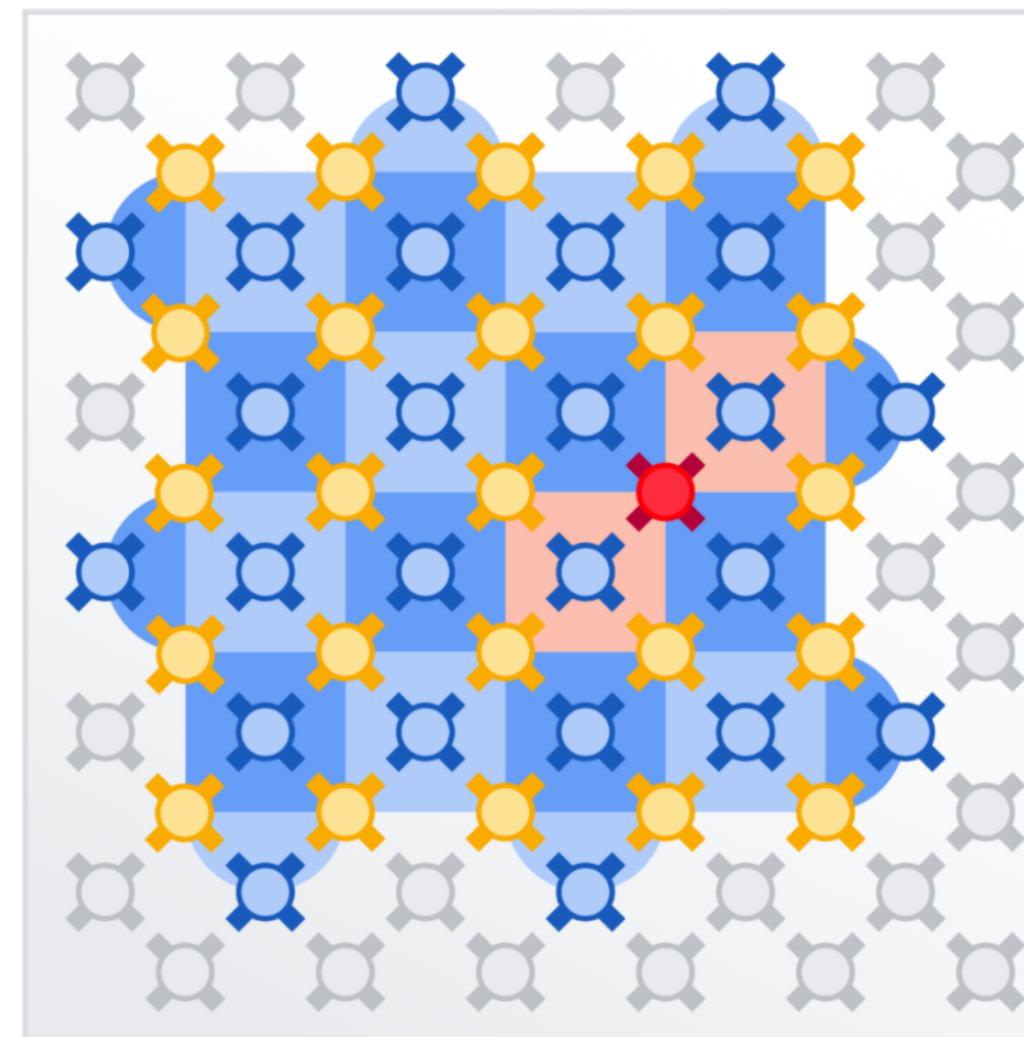
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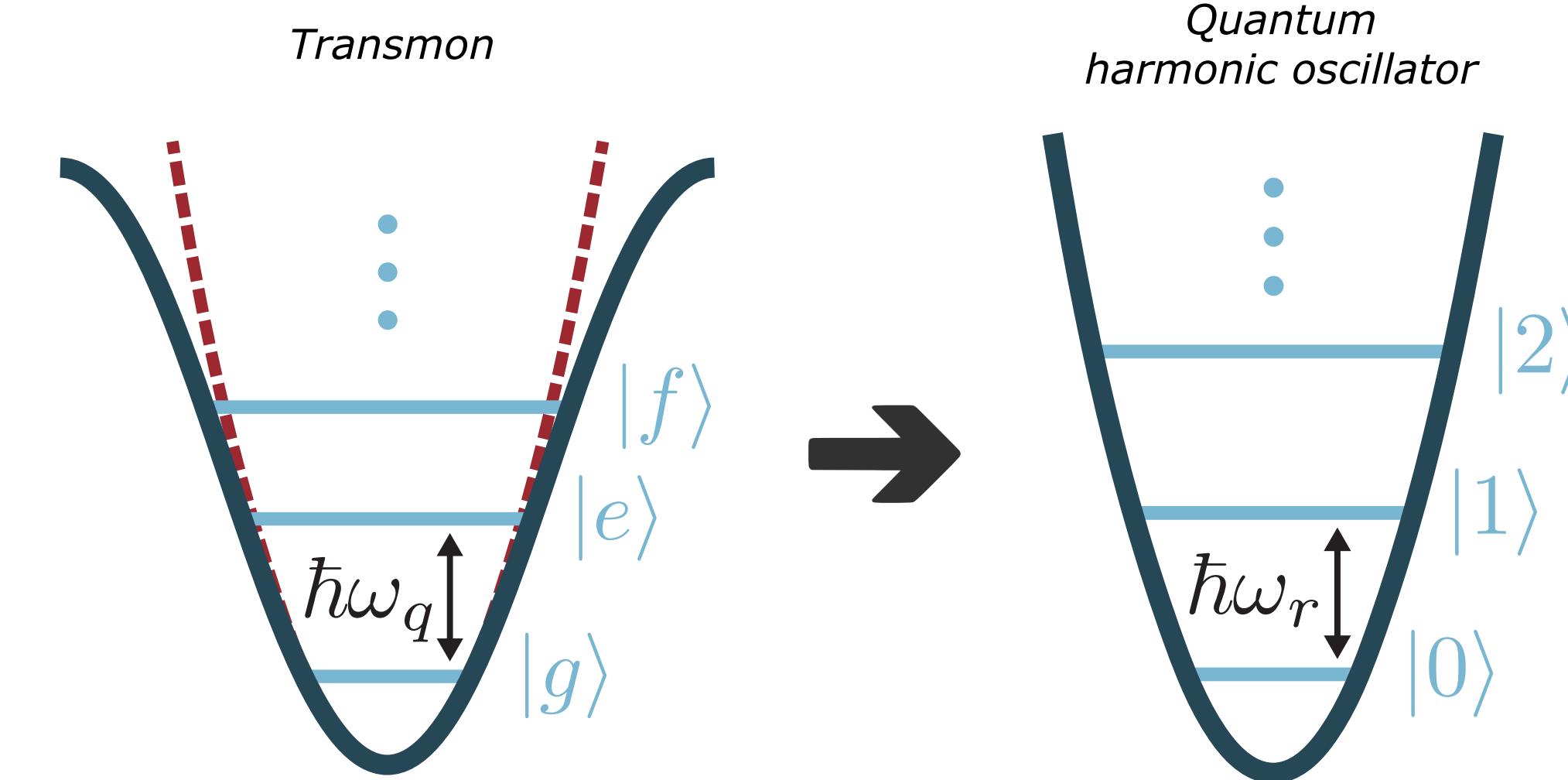
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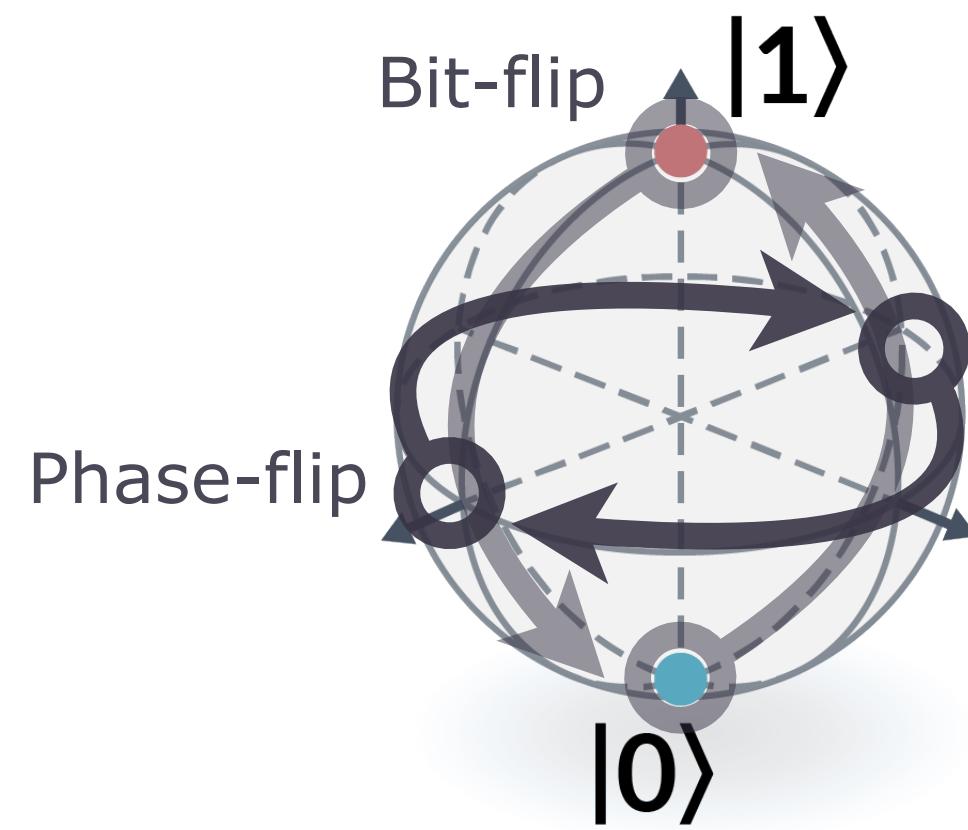
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Bosonic codes



Google Quantum AI, Nature 2022

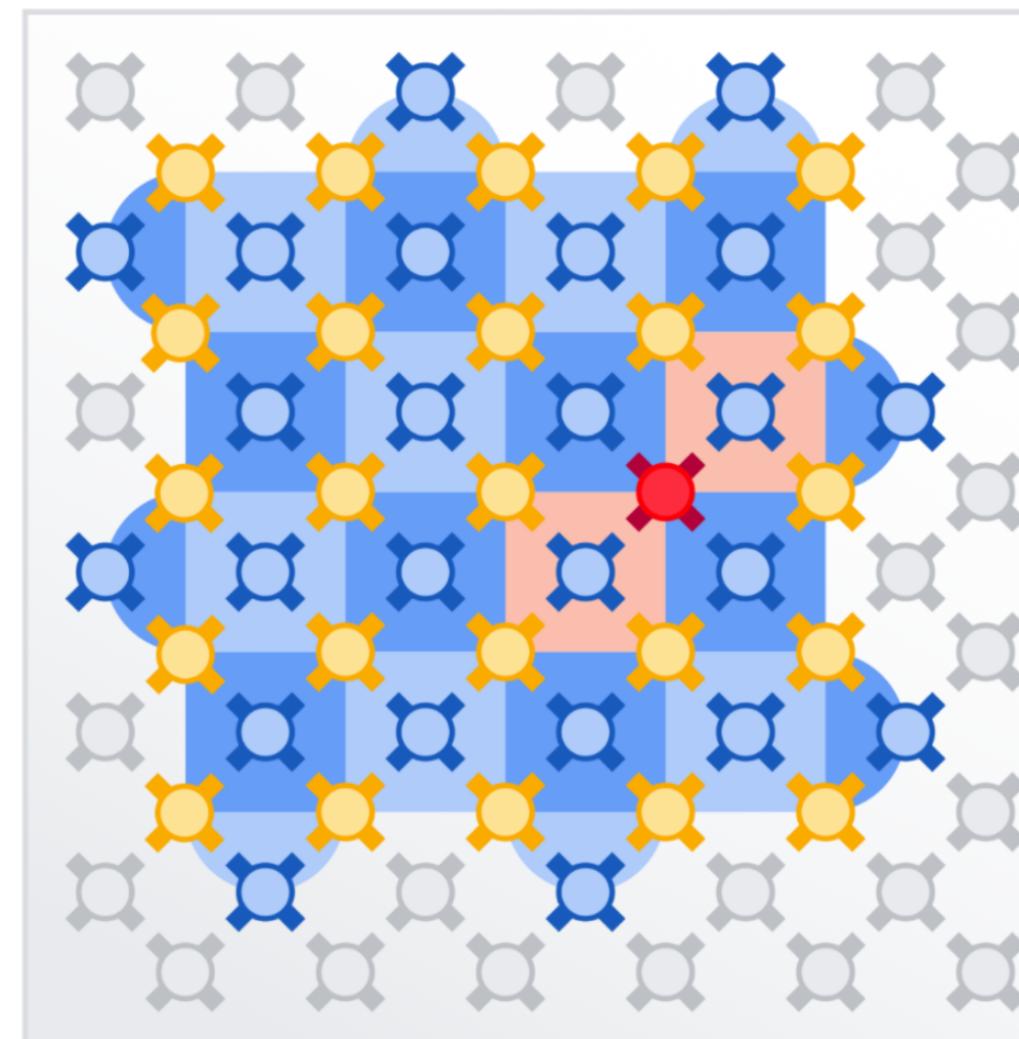
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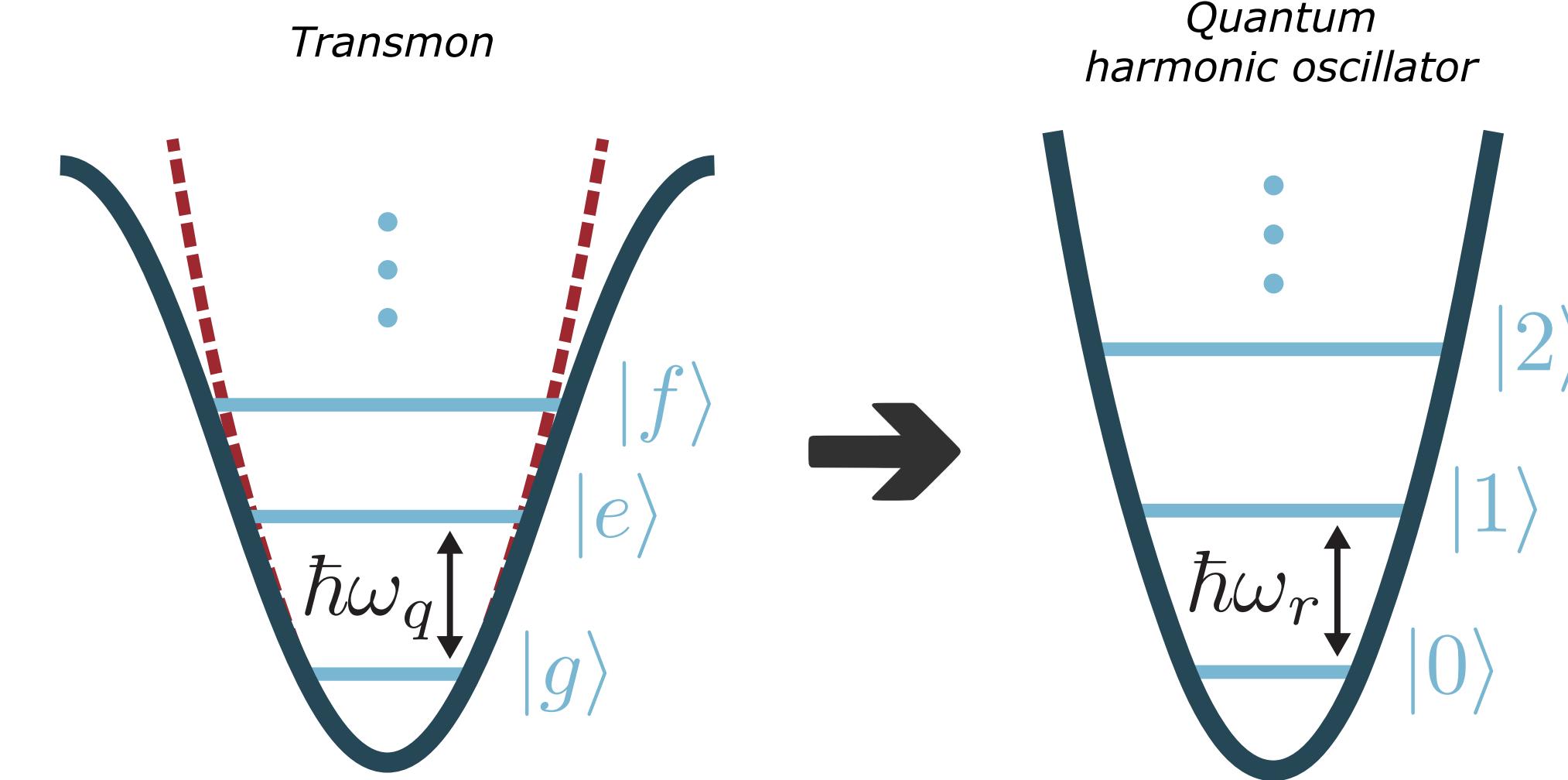
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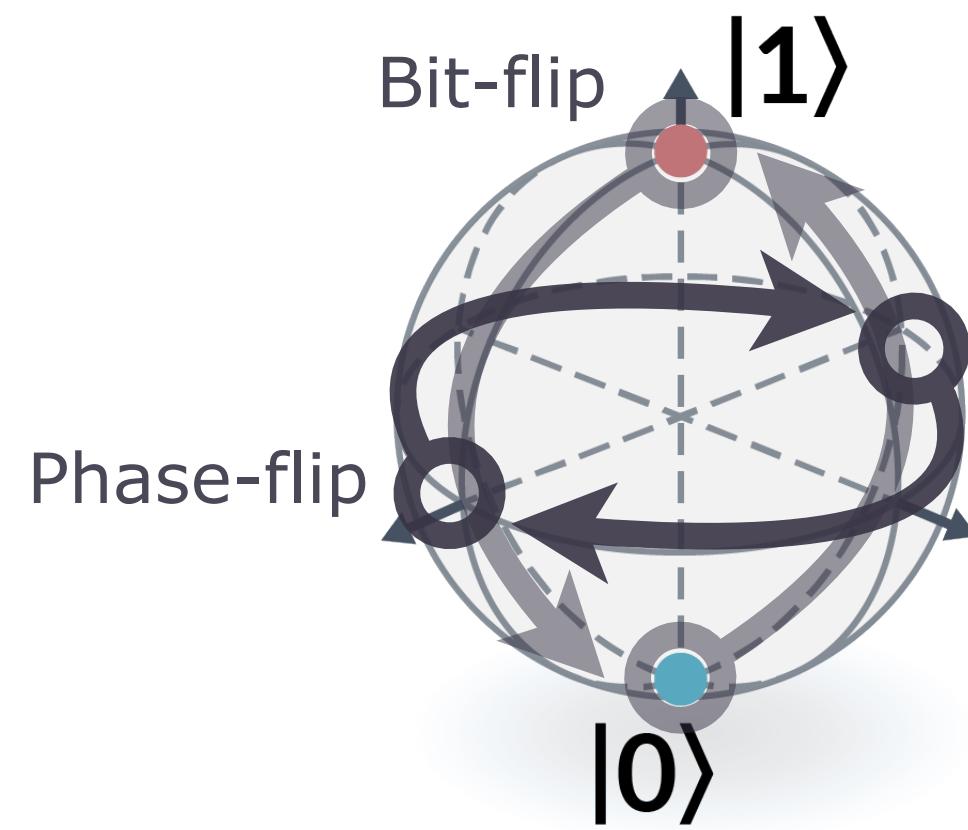
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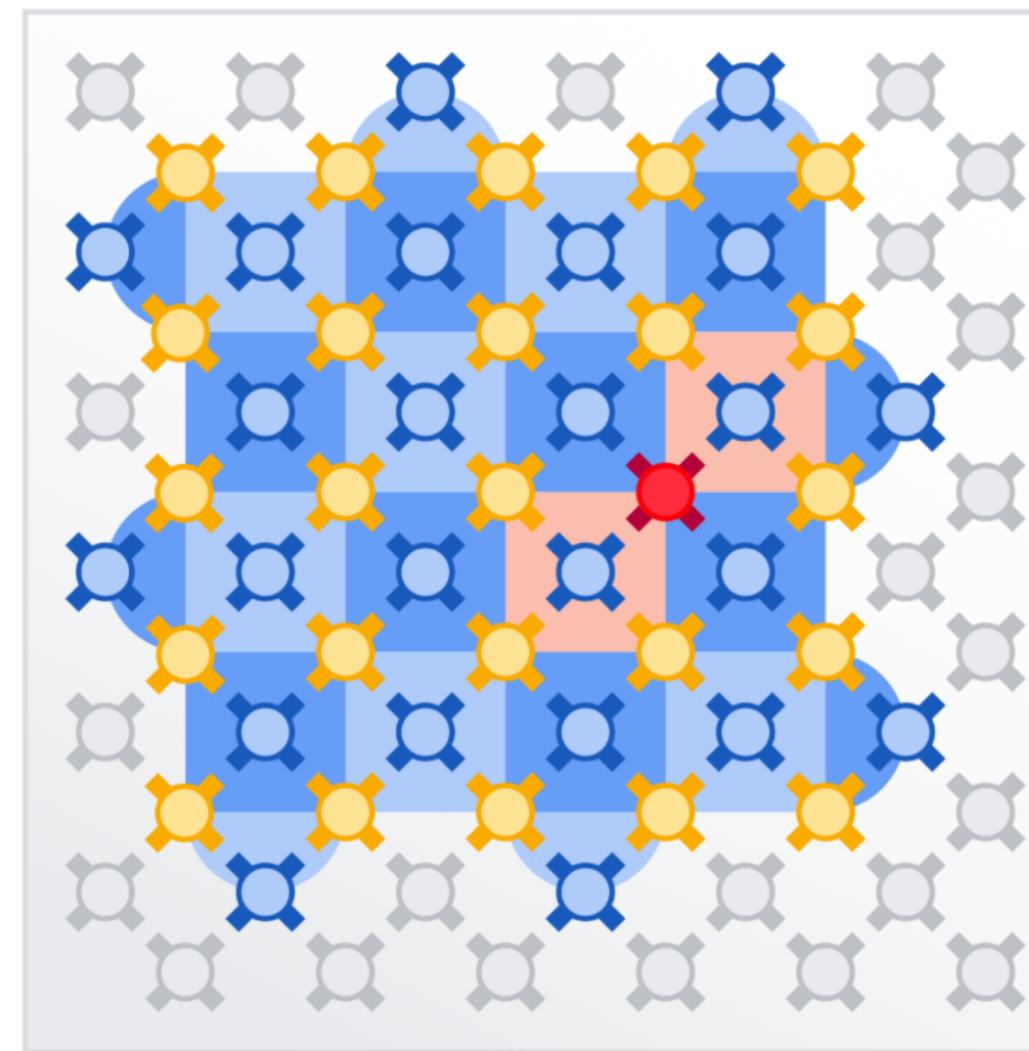
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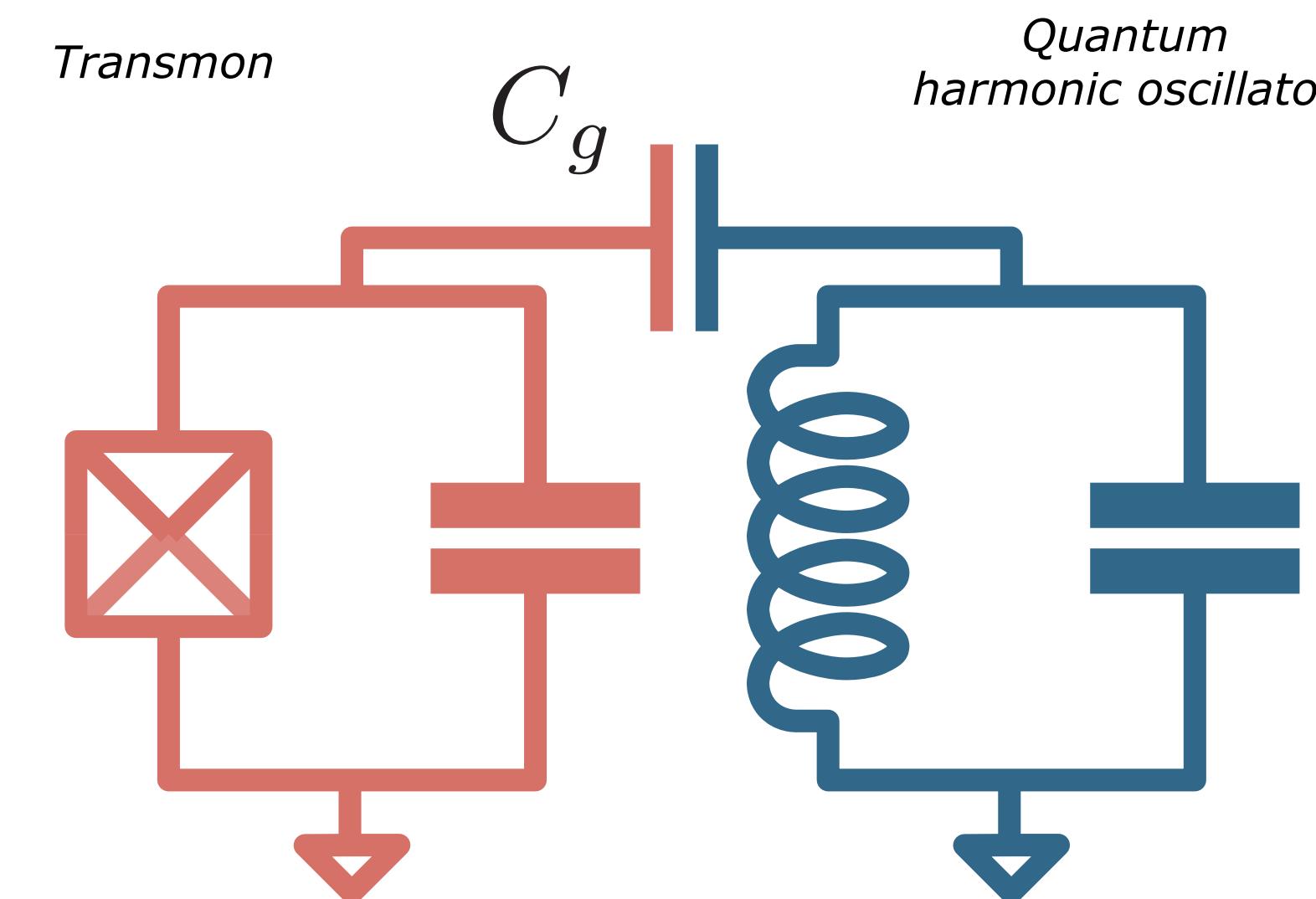
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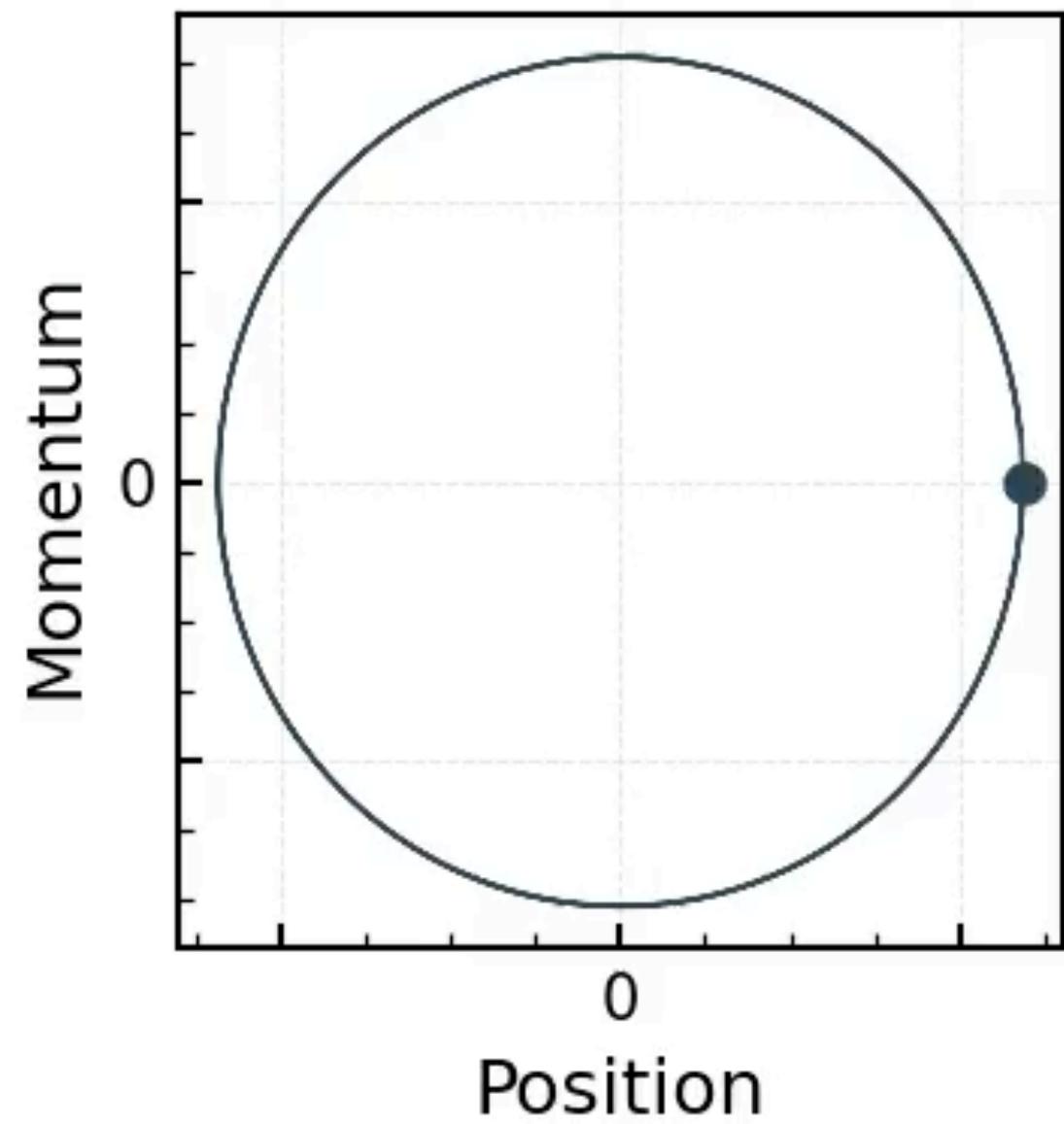
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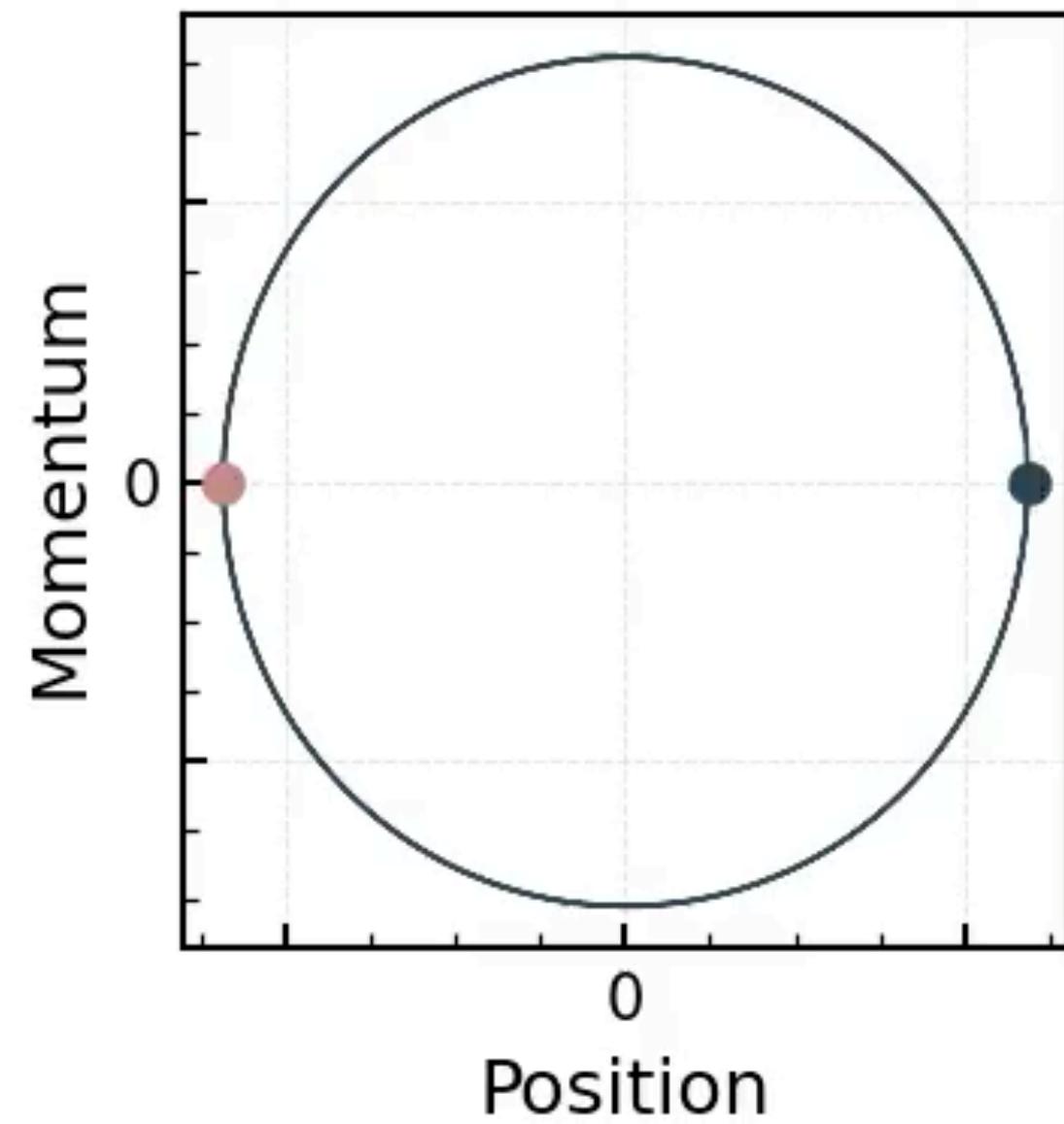
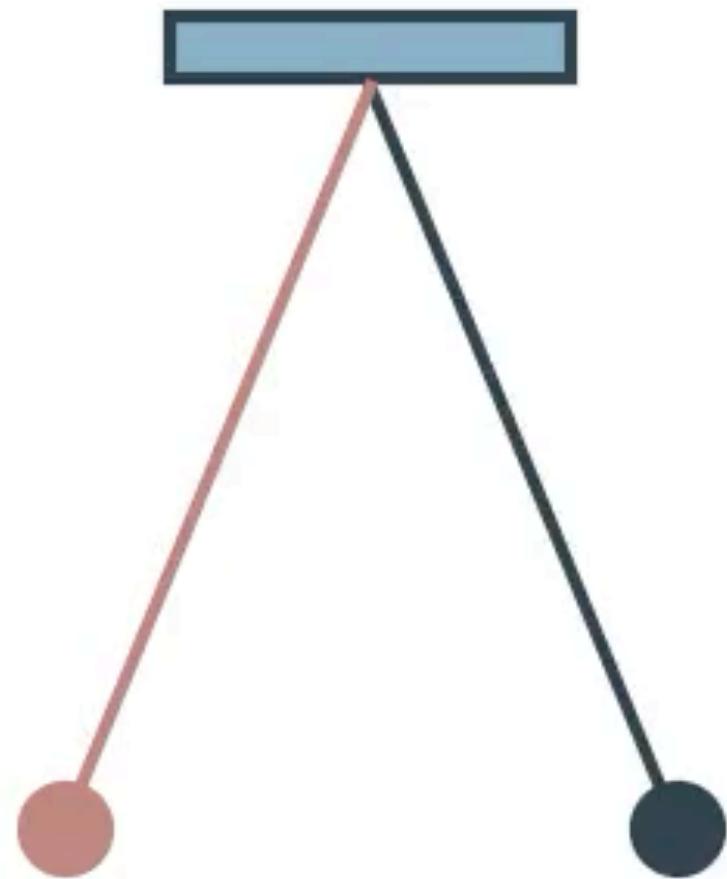
Encoding harmonic oscillators

**How can we encode a
harmonic oscillator?**



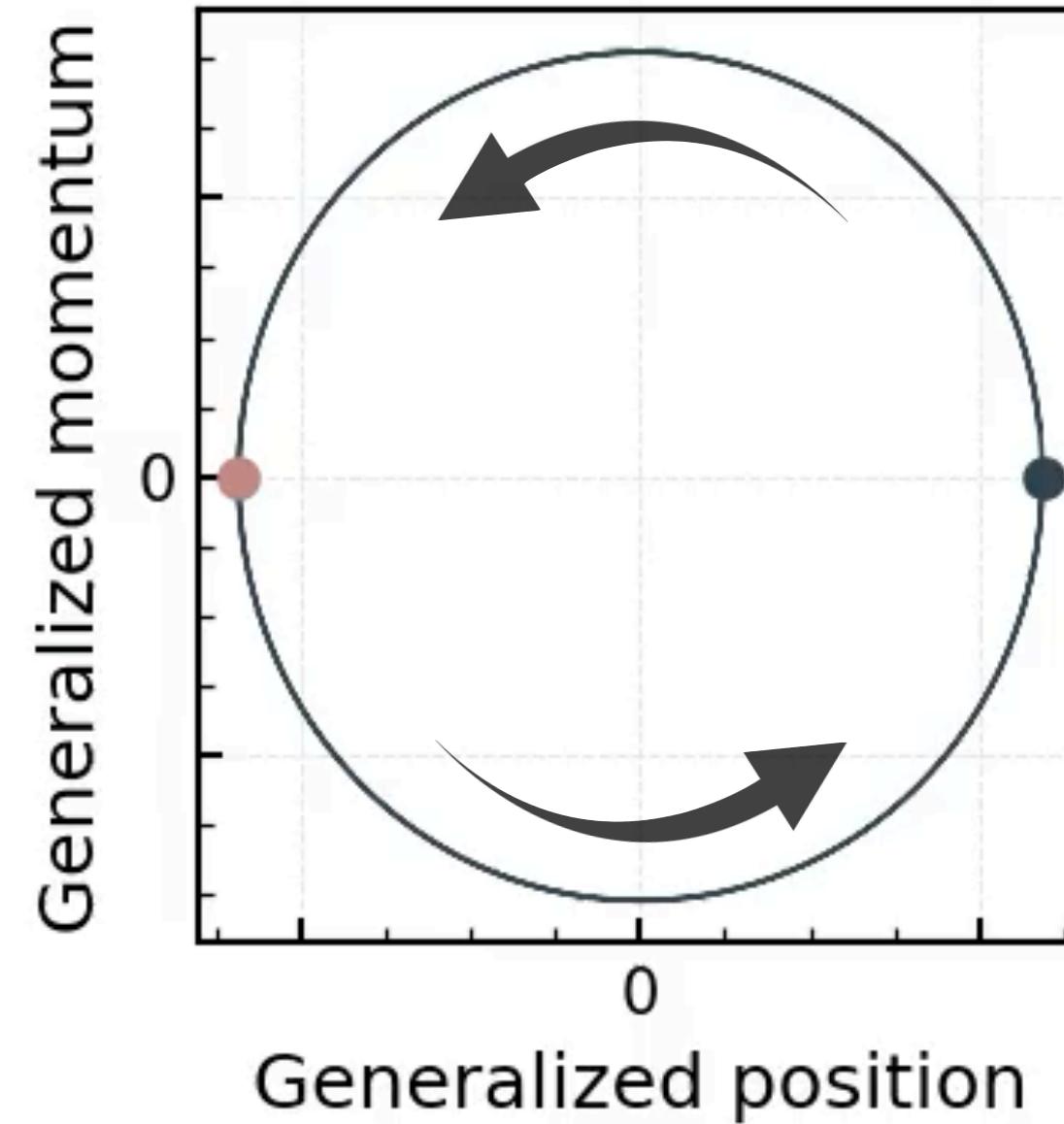
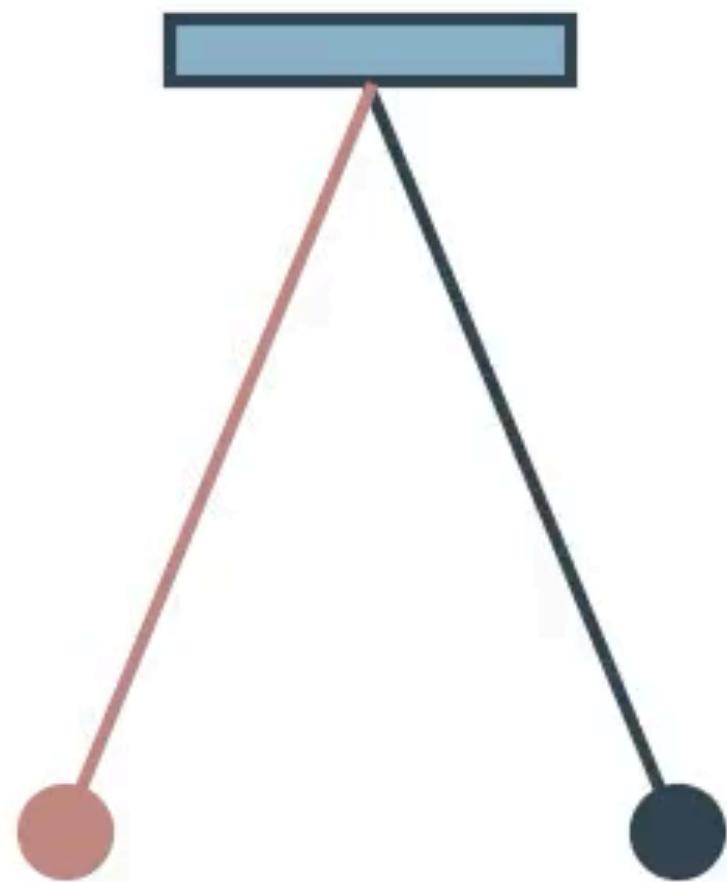
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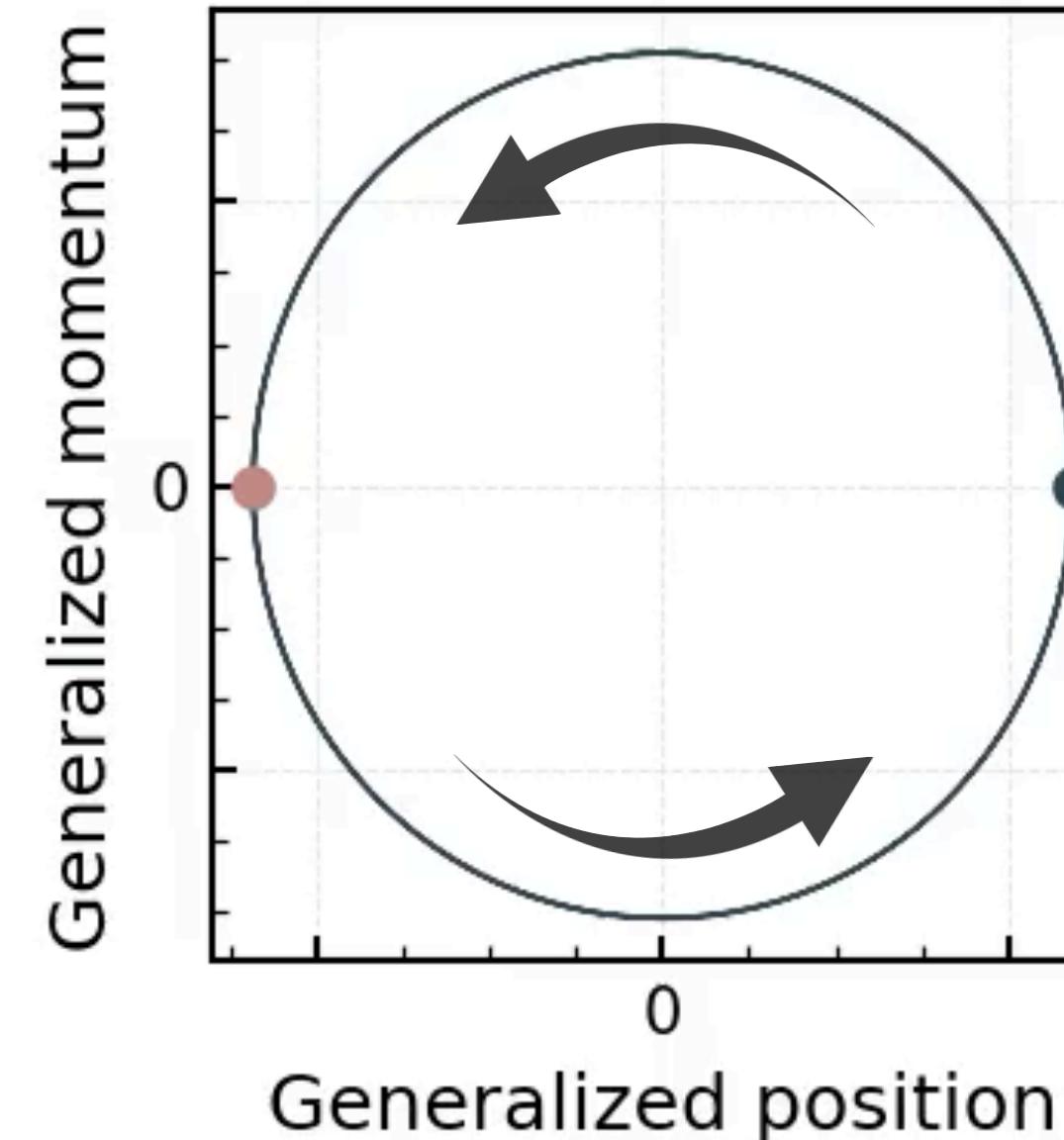
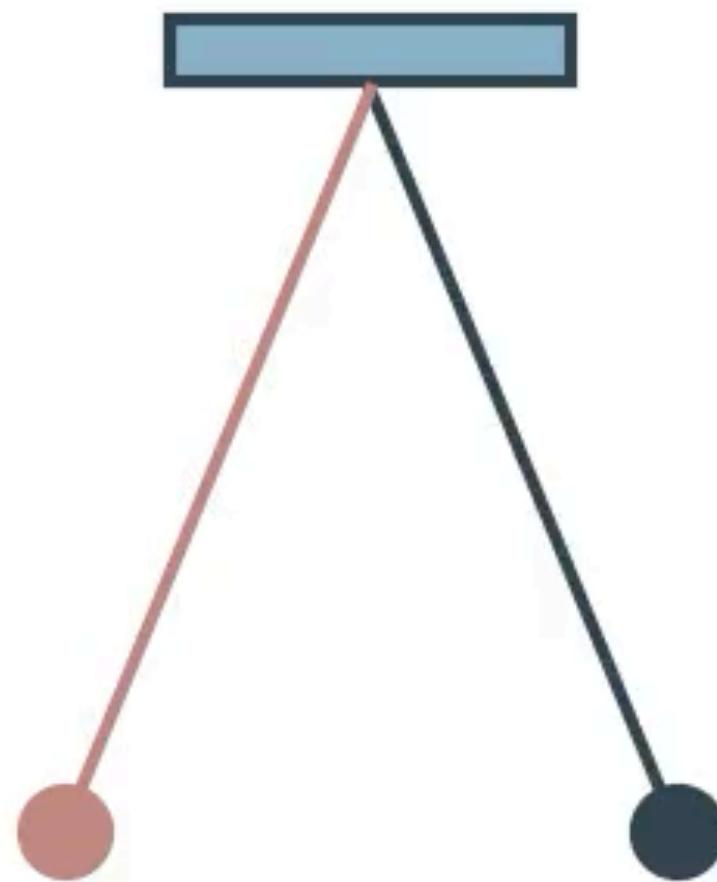
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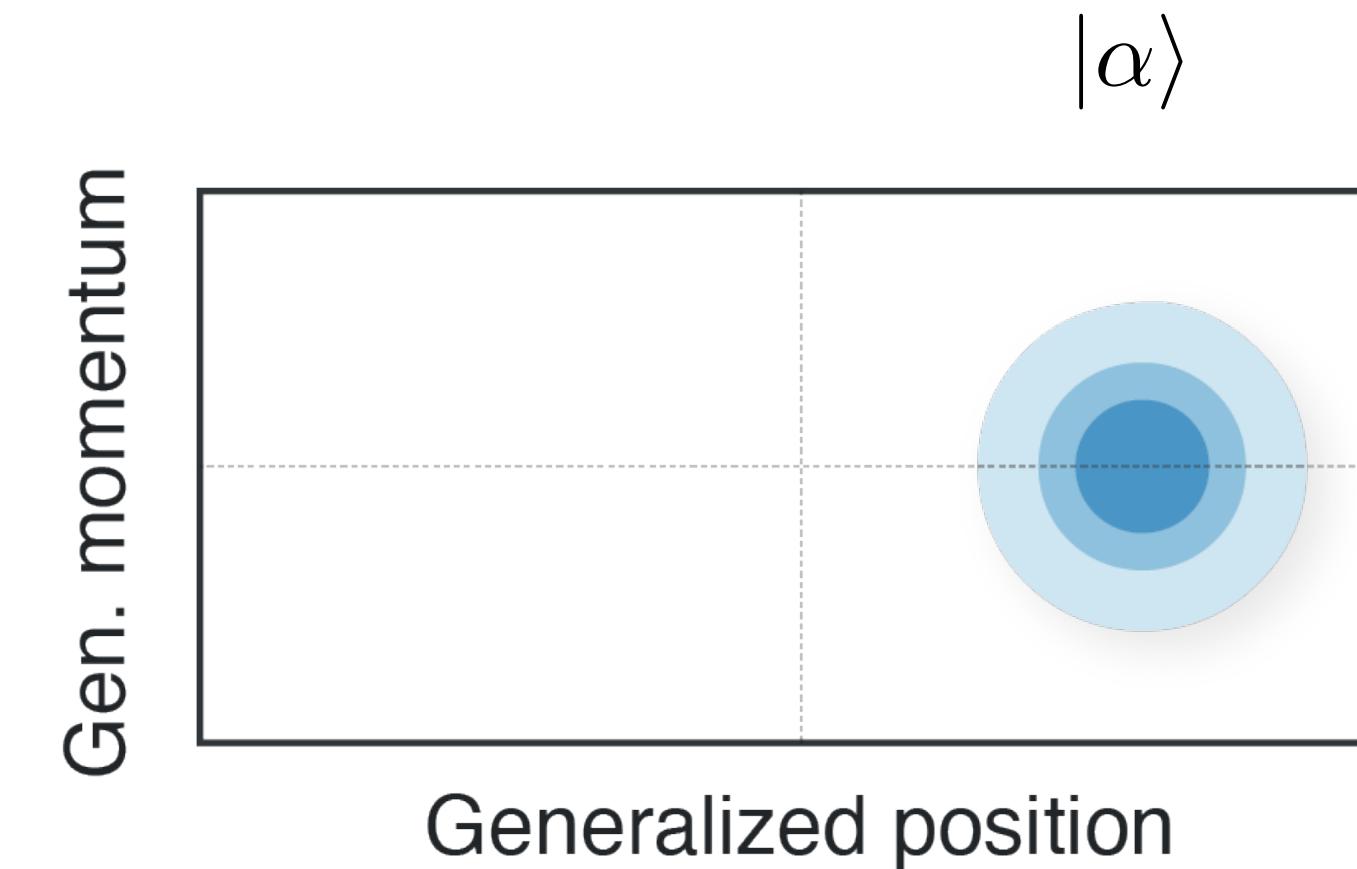


Encoding harmonic oscillators

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How can we encode a quantum harmonic oscillator?

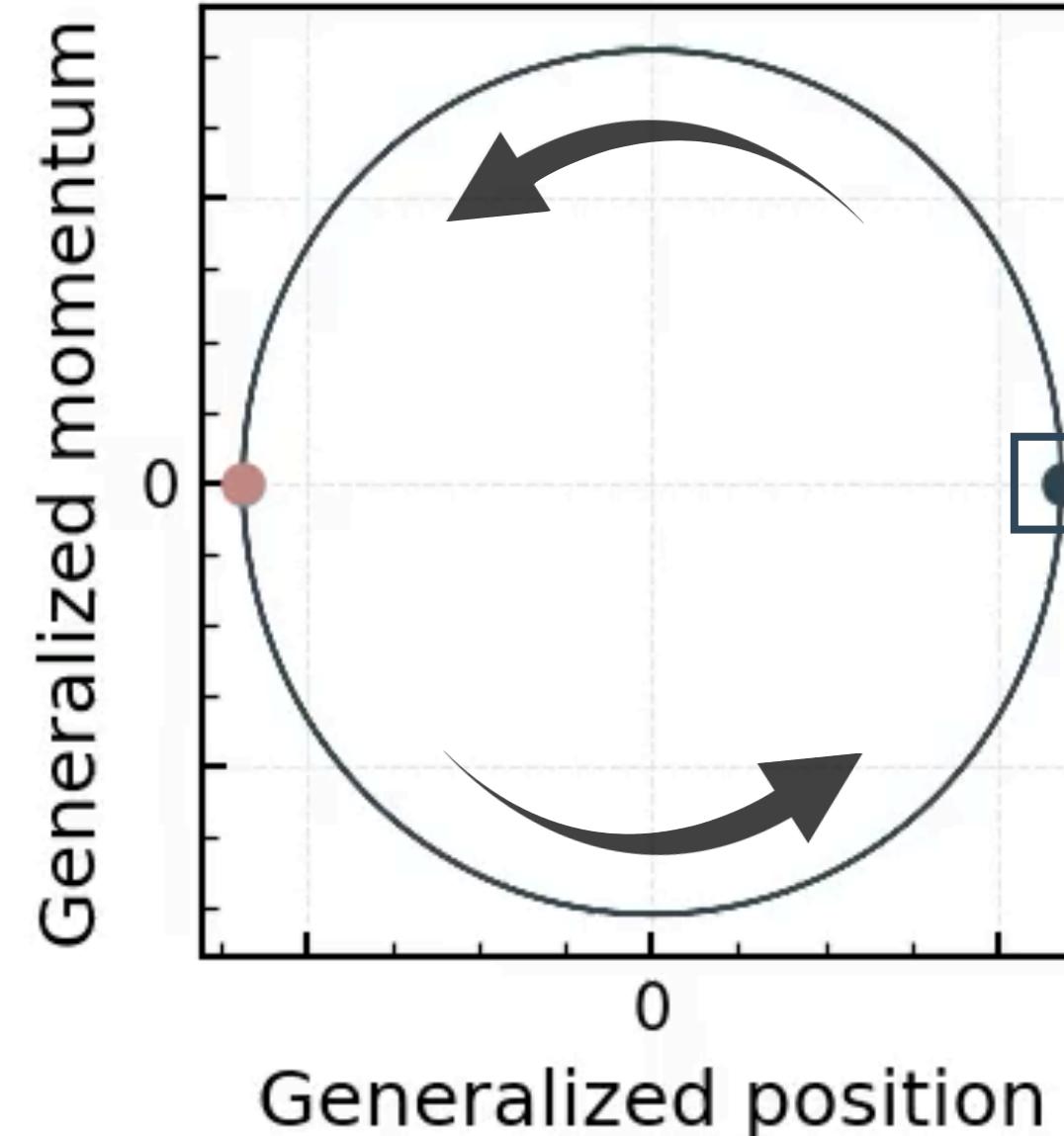


where $\hat{a}|\pm\alpha\rangle = \pm|\pm\alpha\rangle$

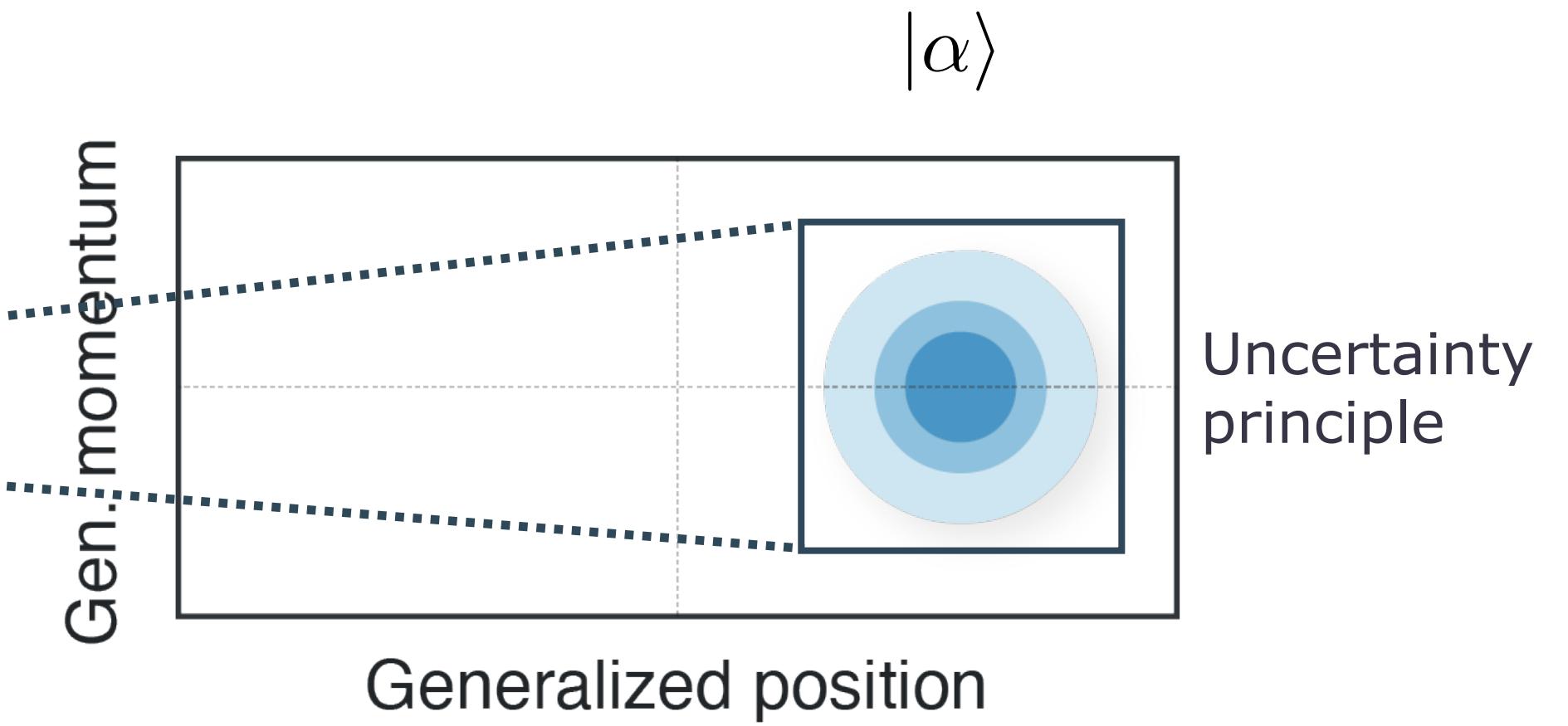
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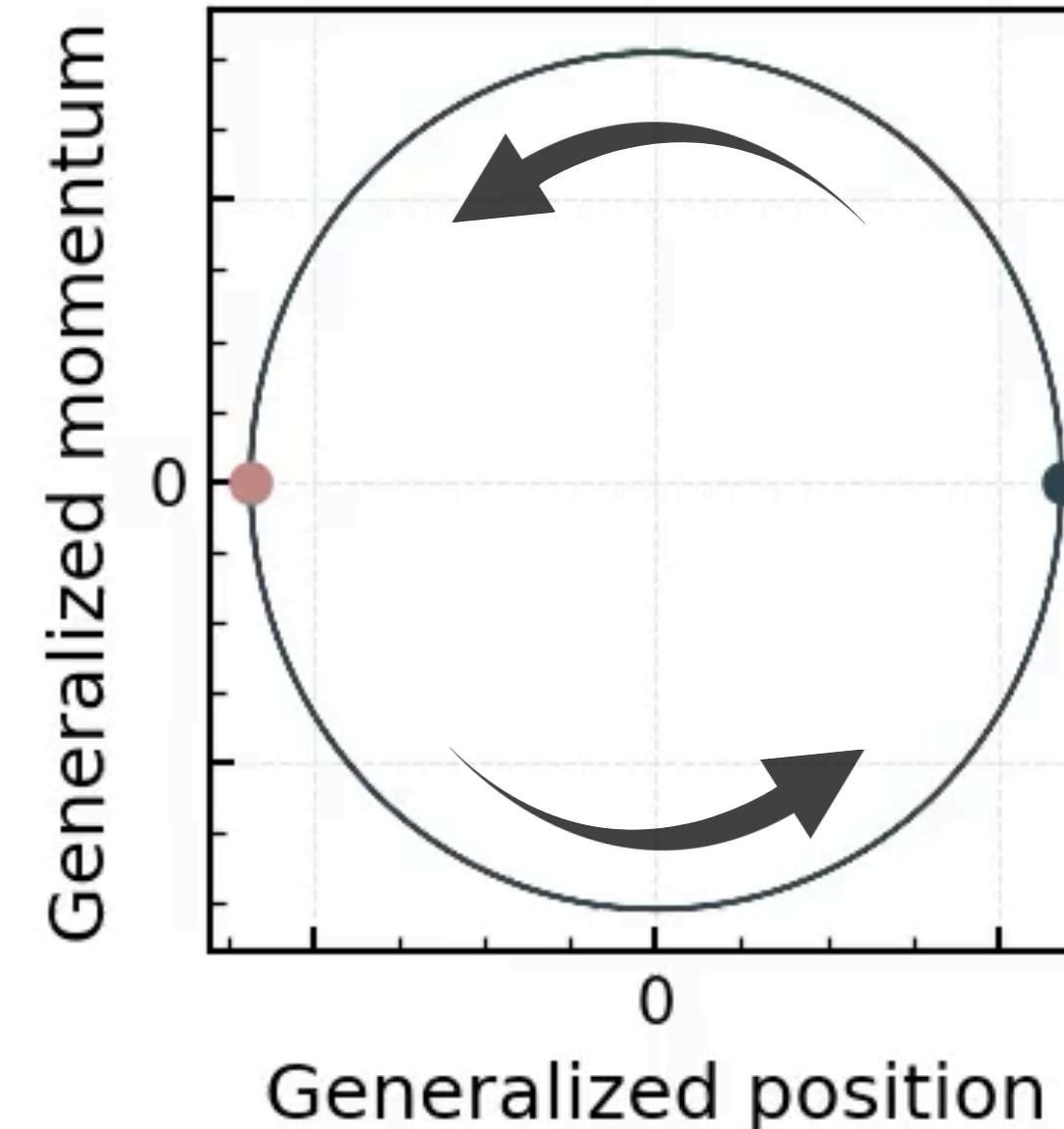
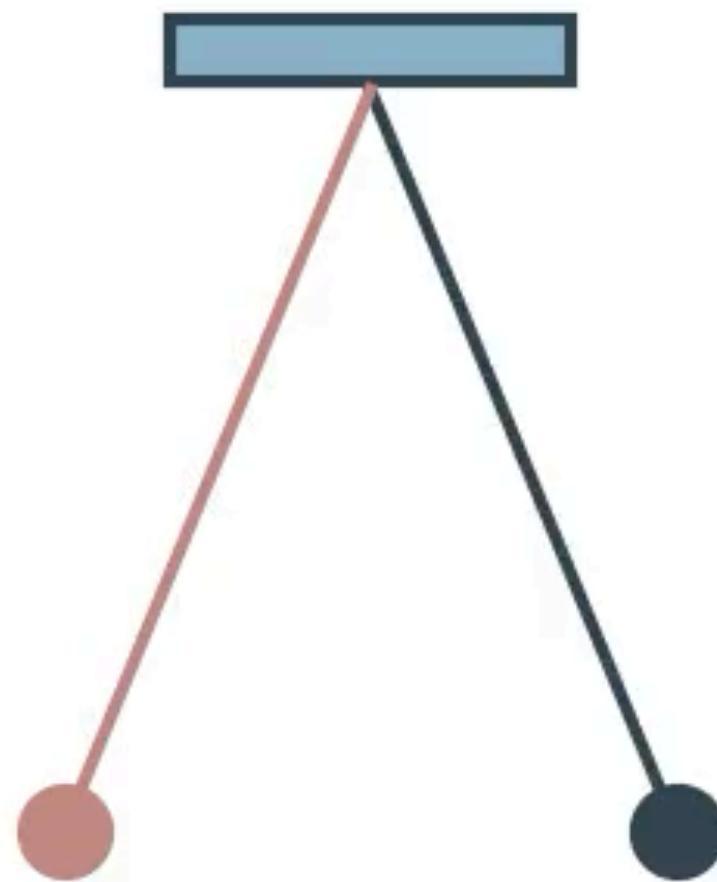
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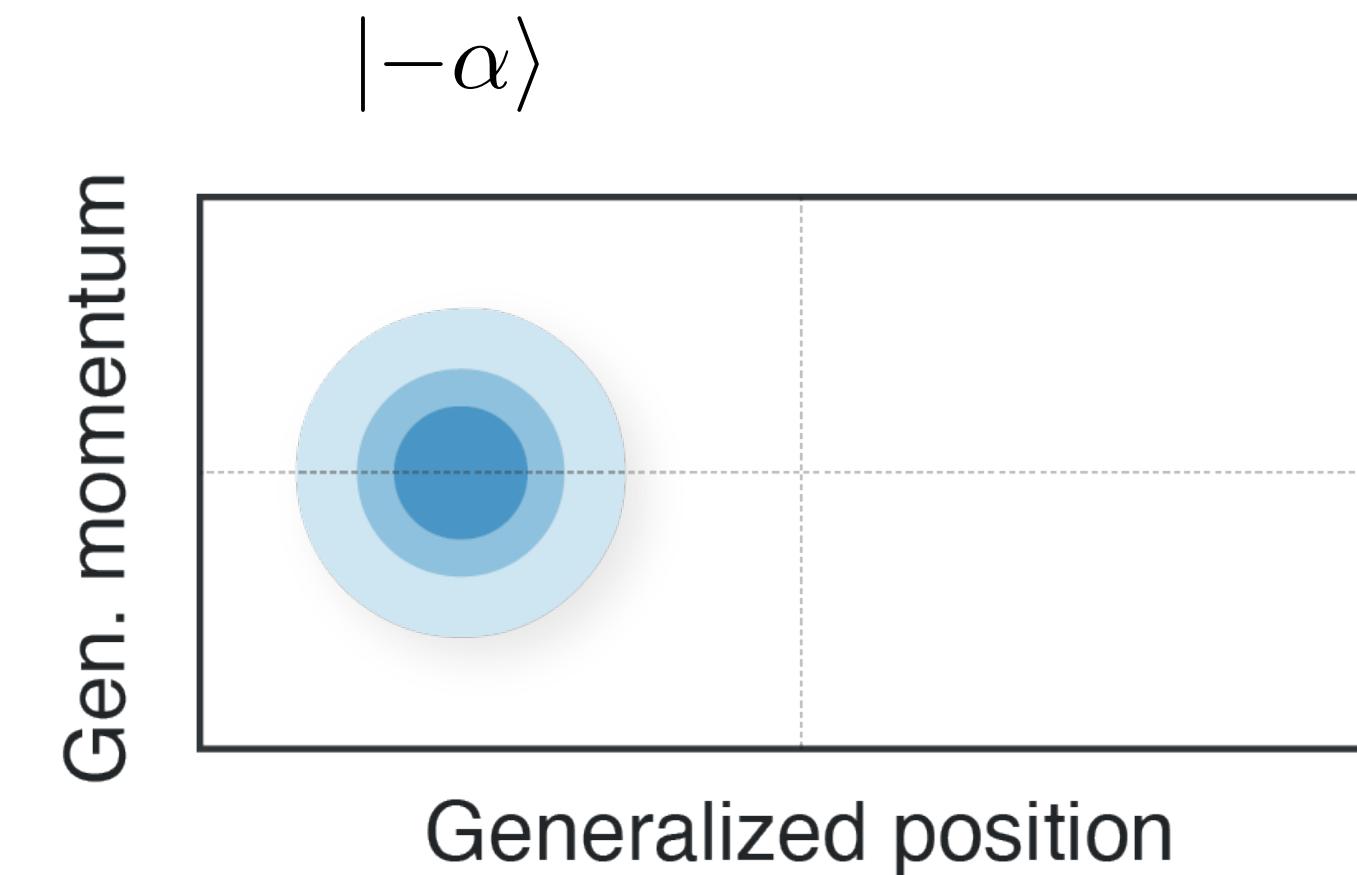
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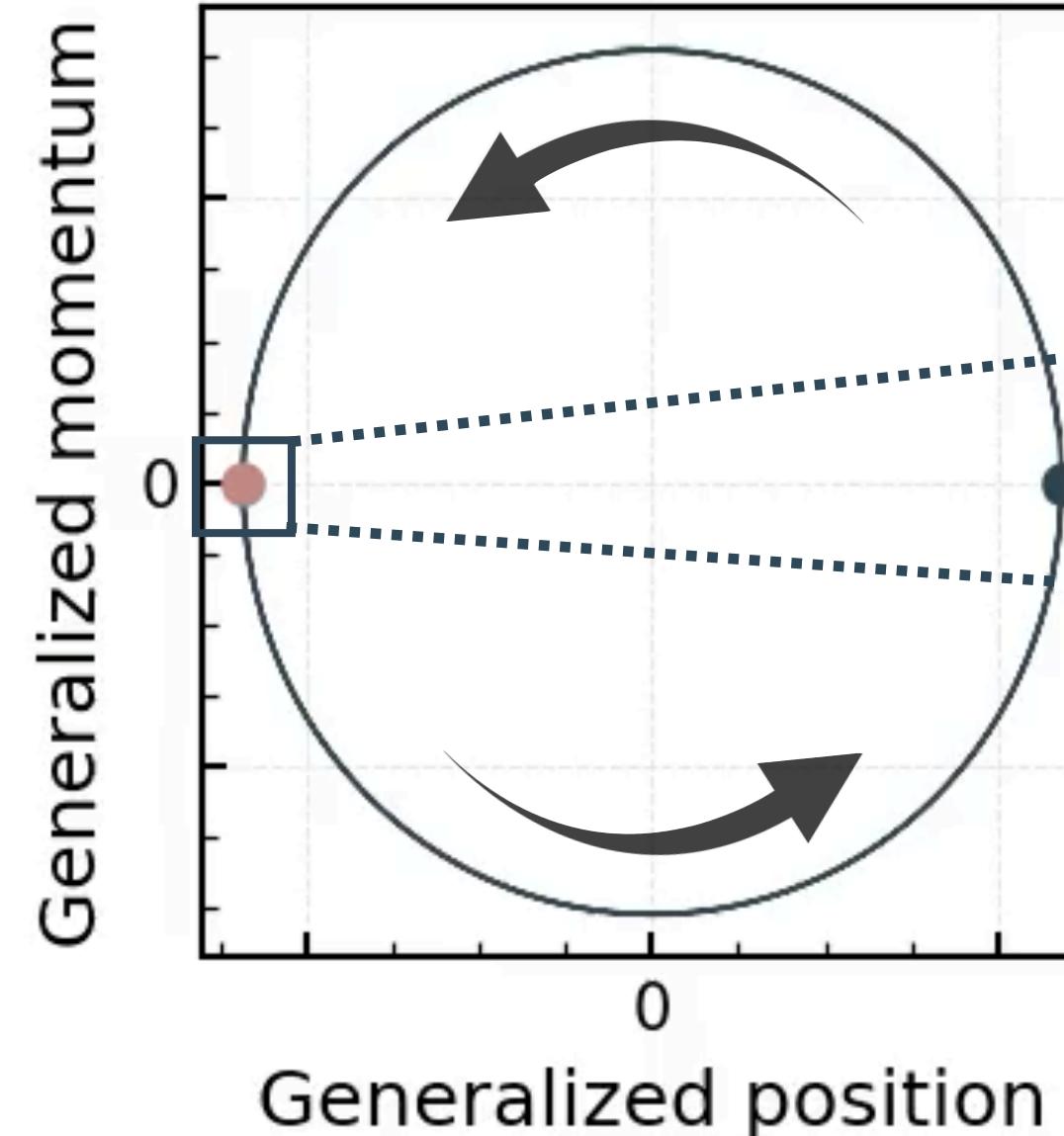


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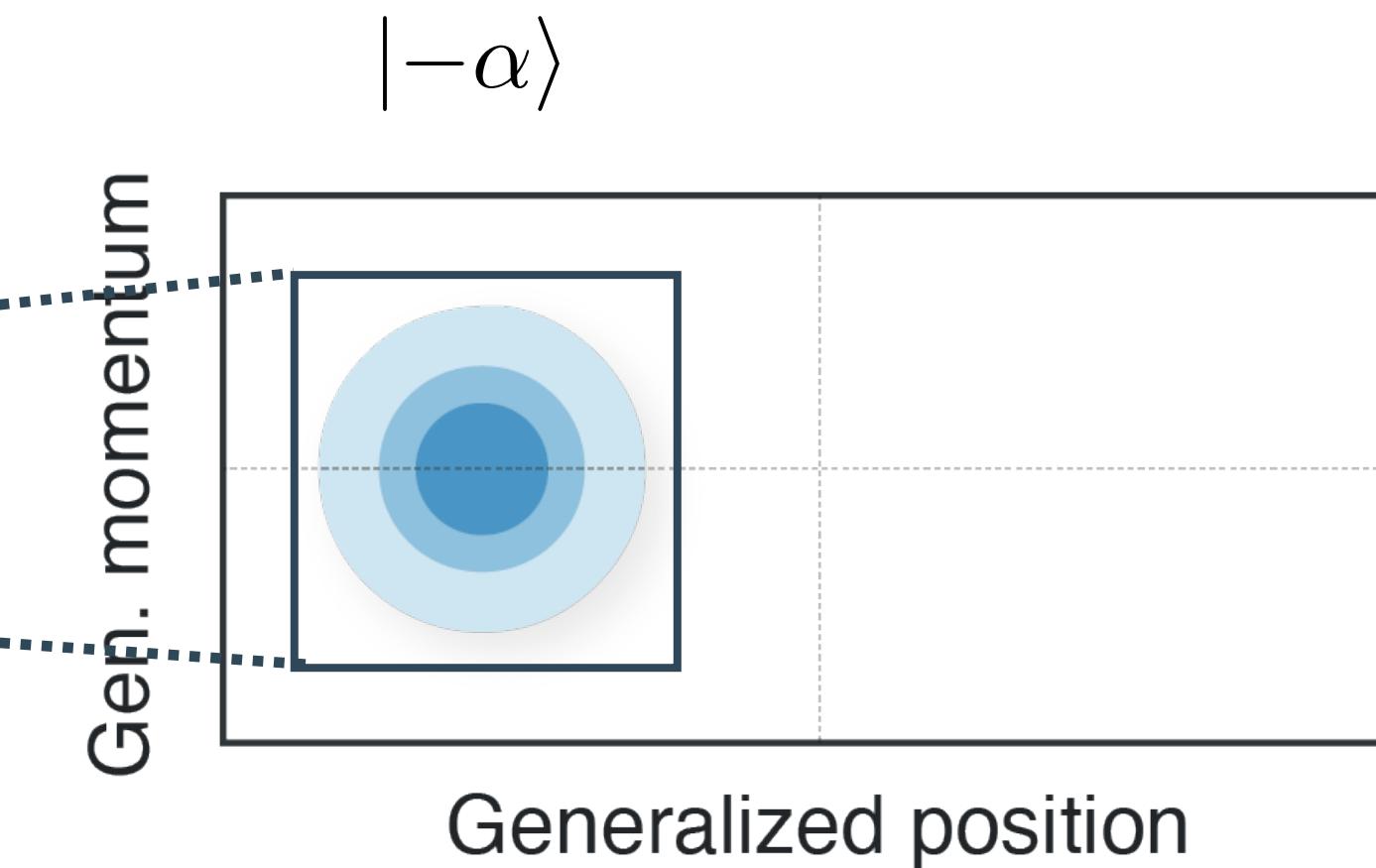
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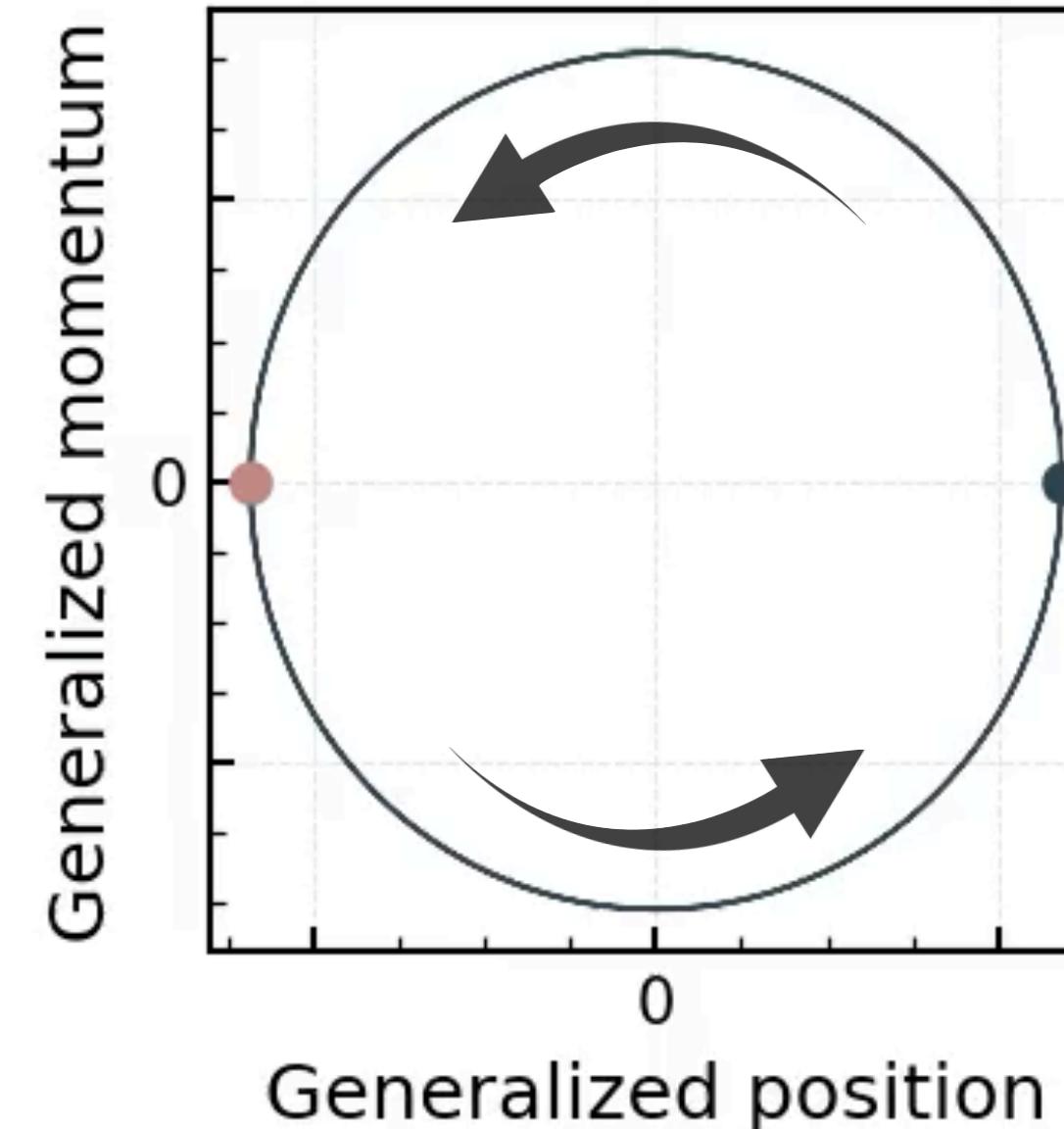
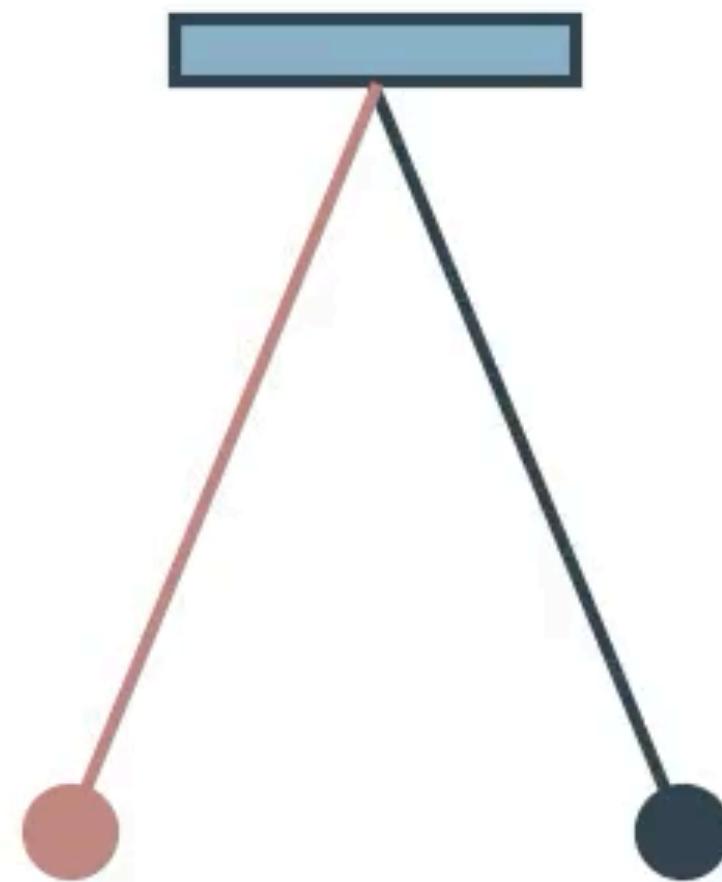
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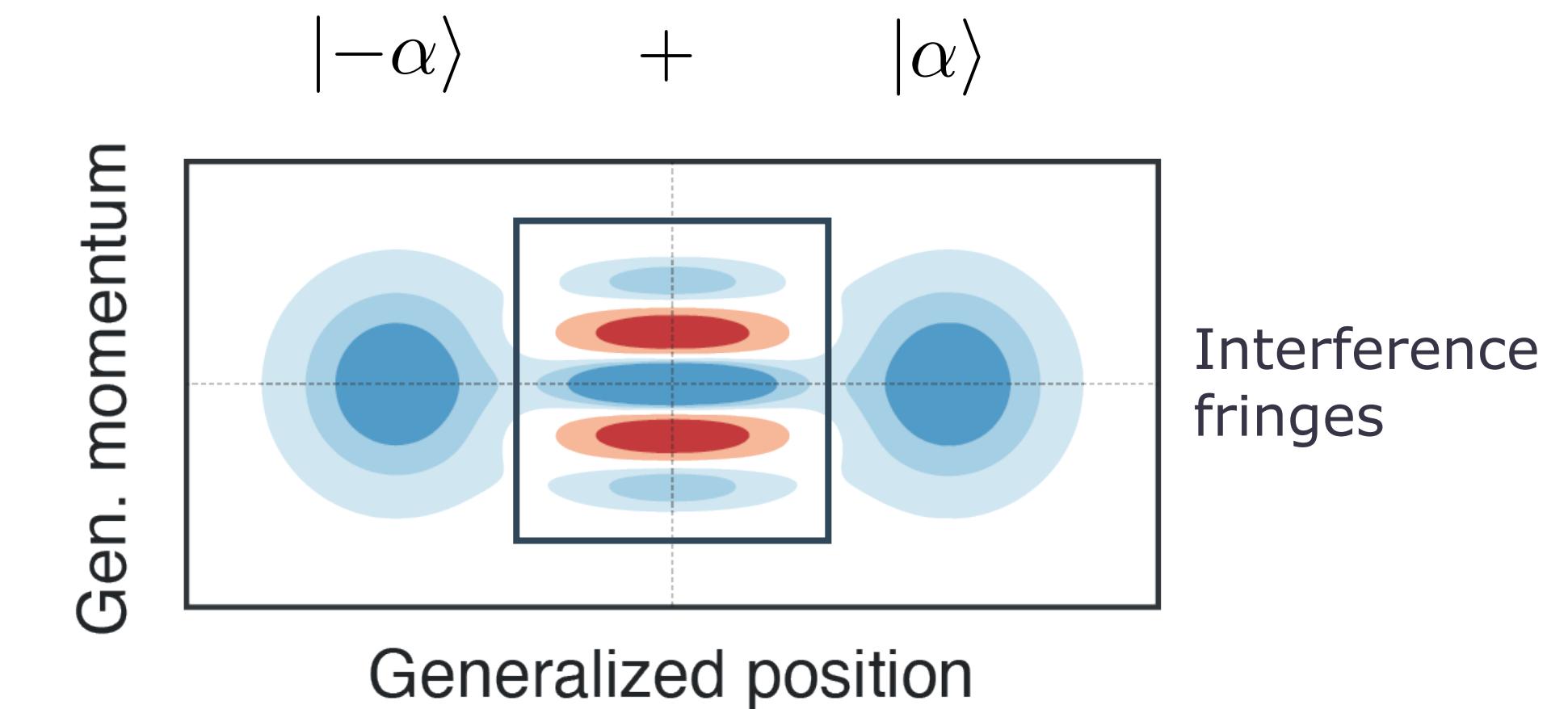
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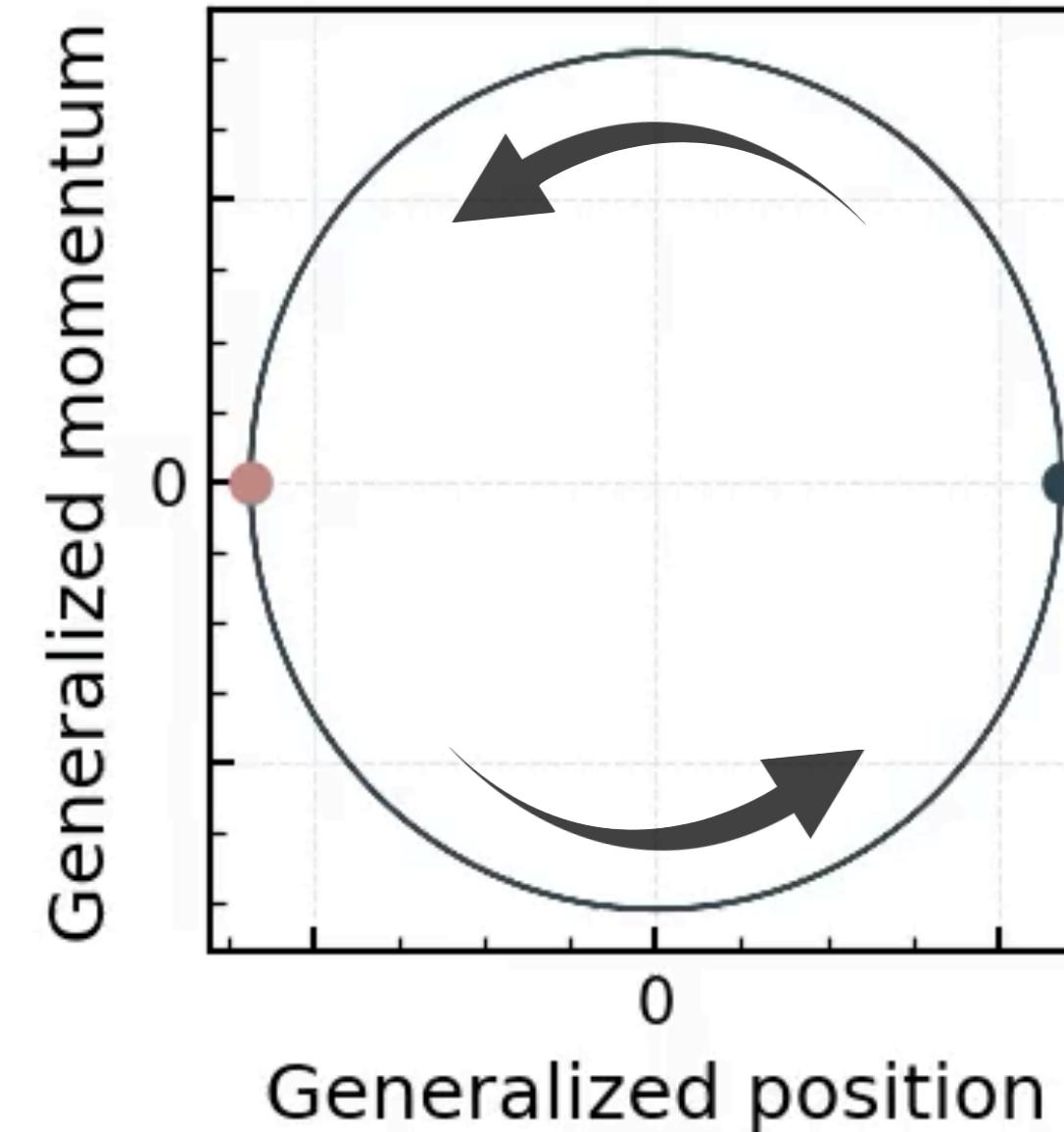
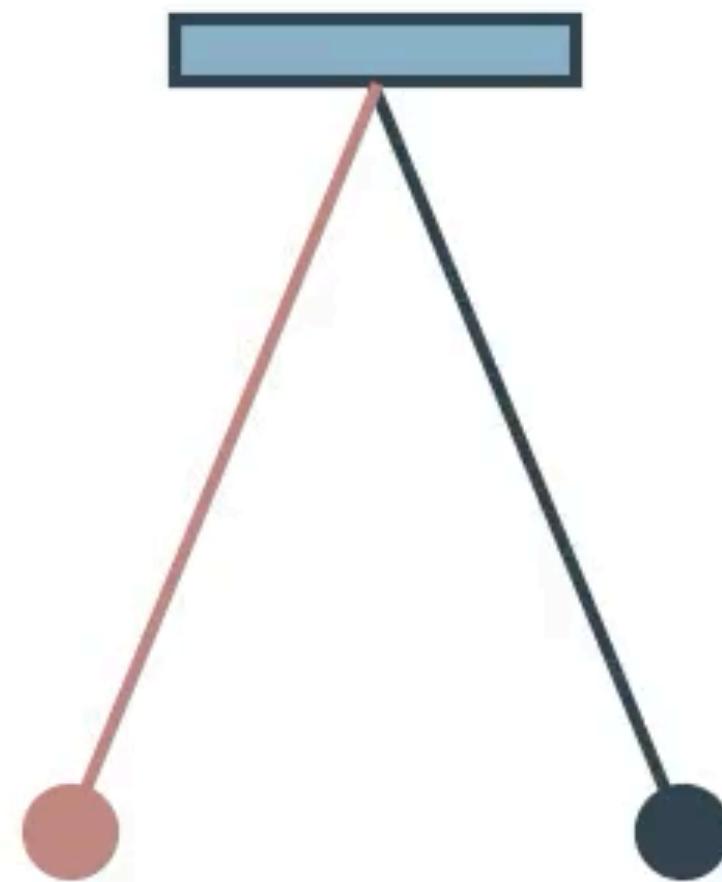


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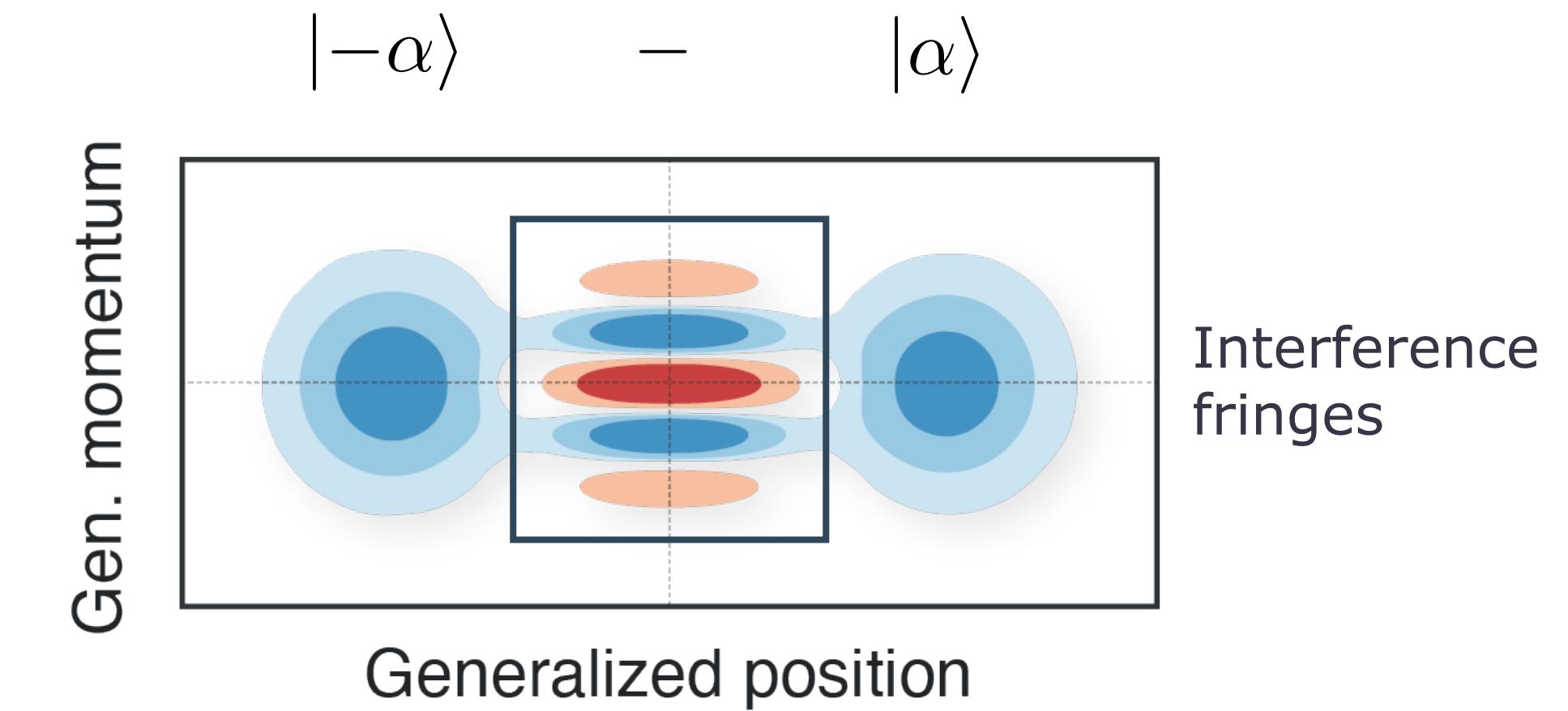
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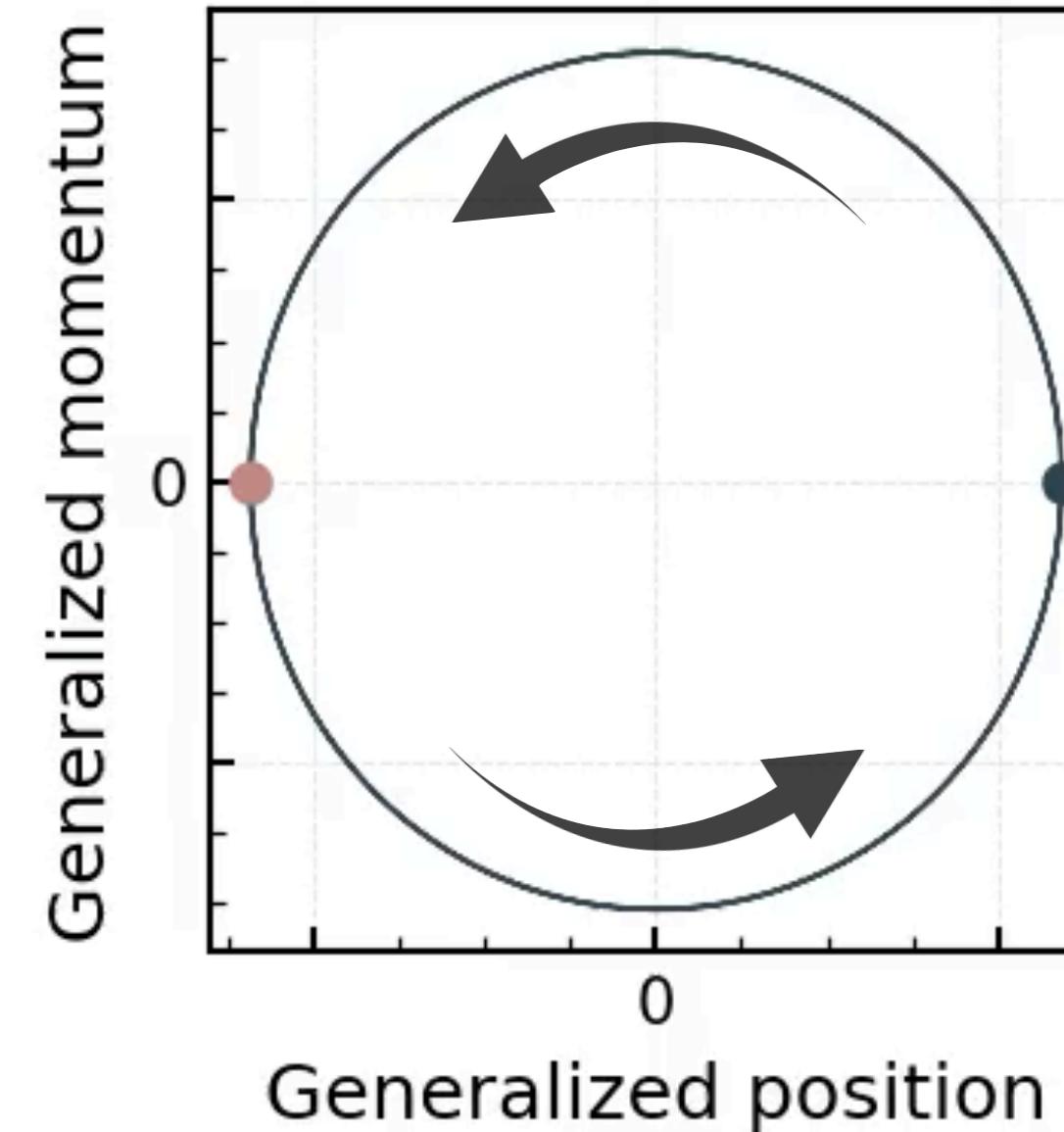
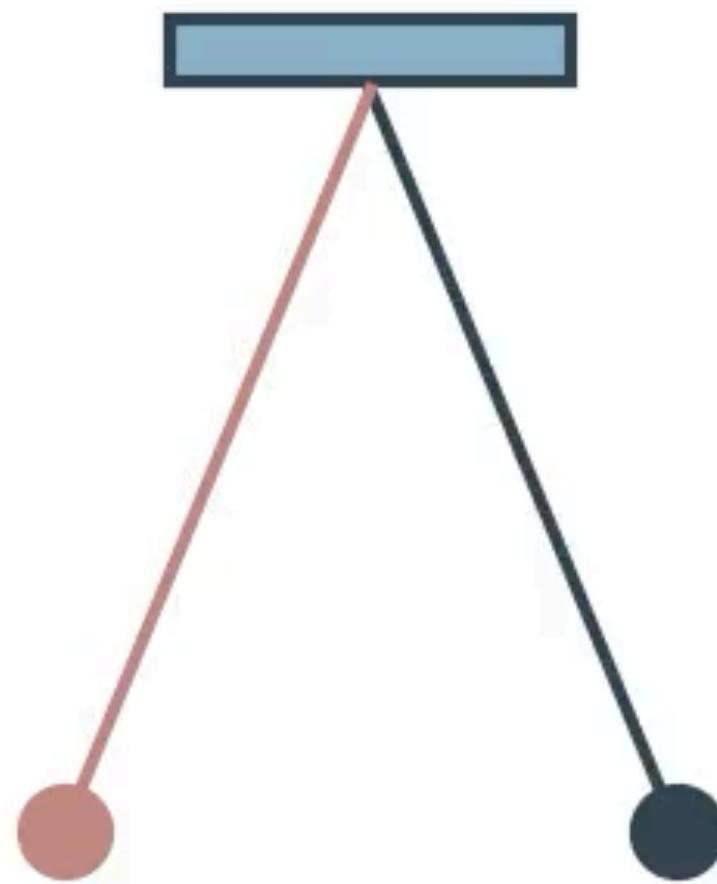


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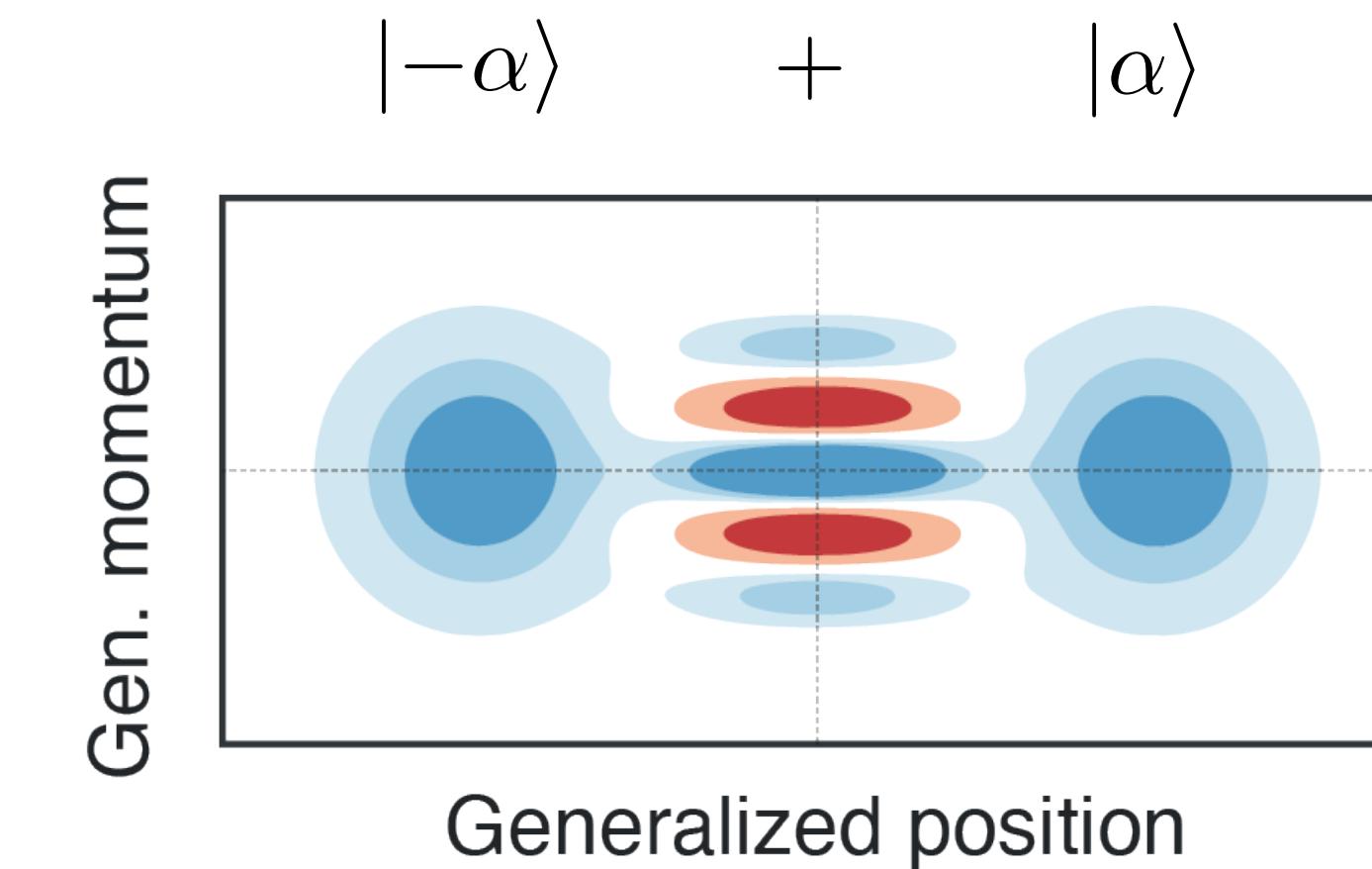
with $\hat{a} = \hat{x} + i\hat{p}$

Encoding harmonic oscillators

How can we encode a harmonic oscillator?



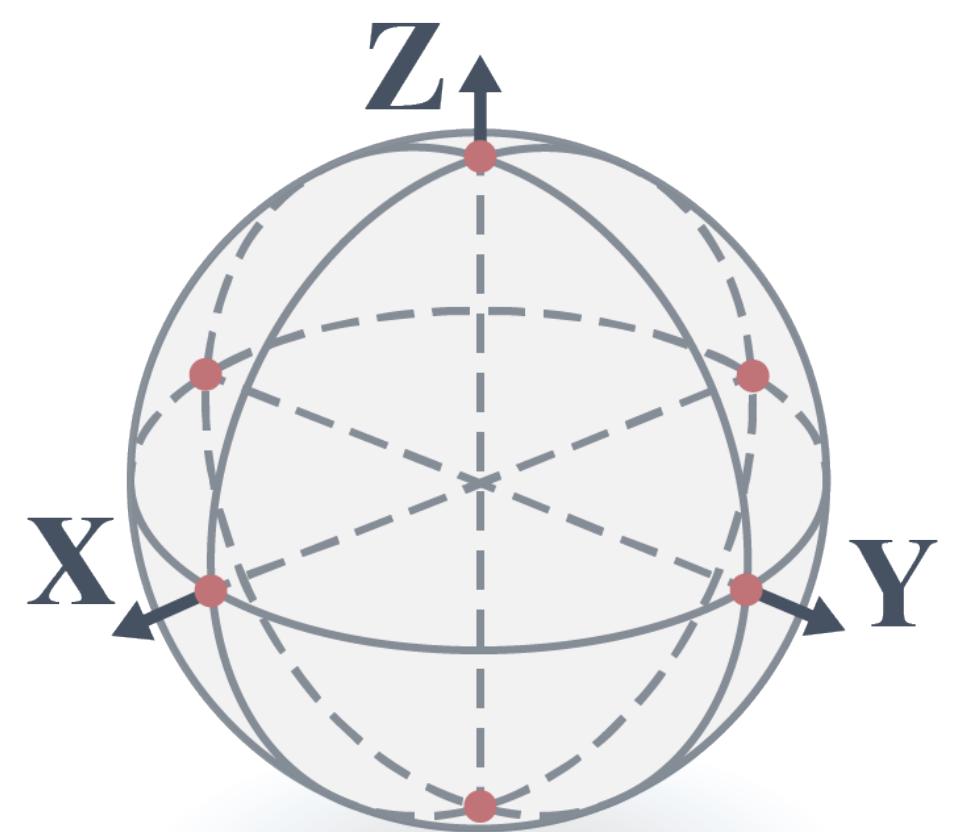
How can we encode a quantum harmonic oscillator?



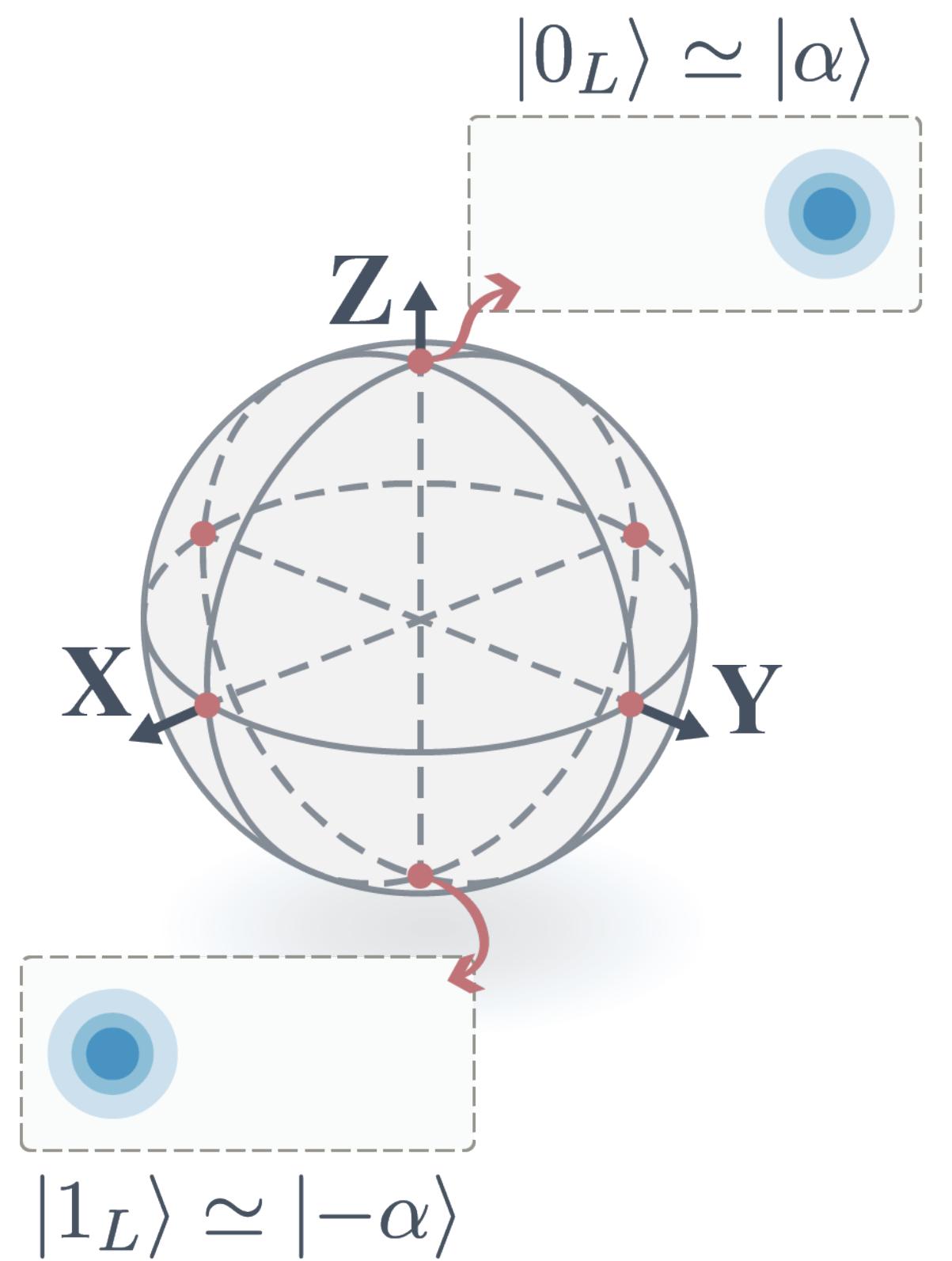
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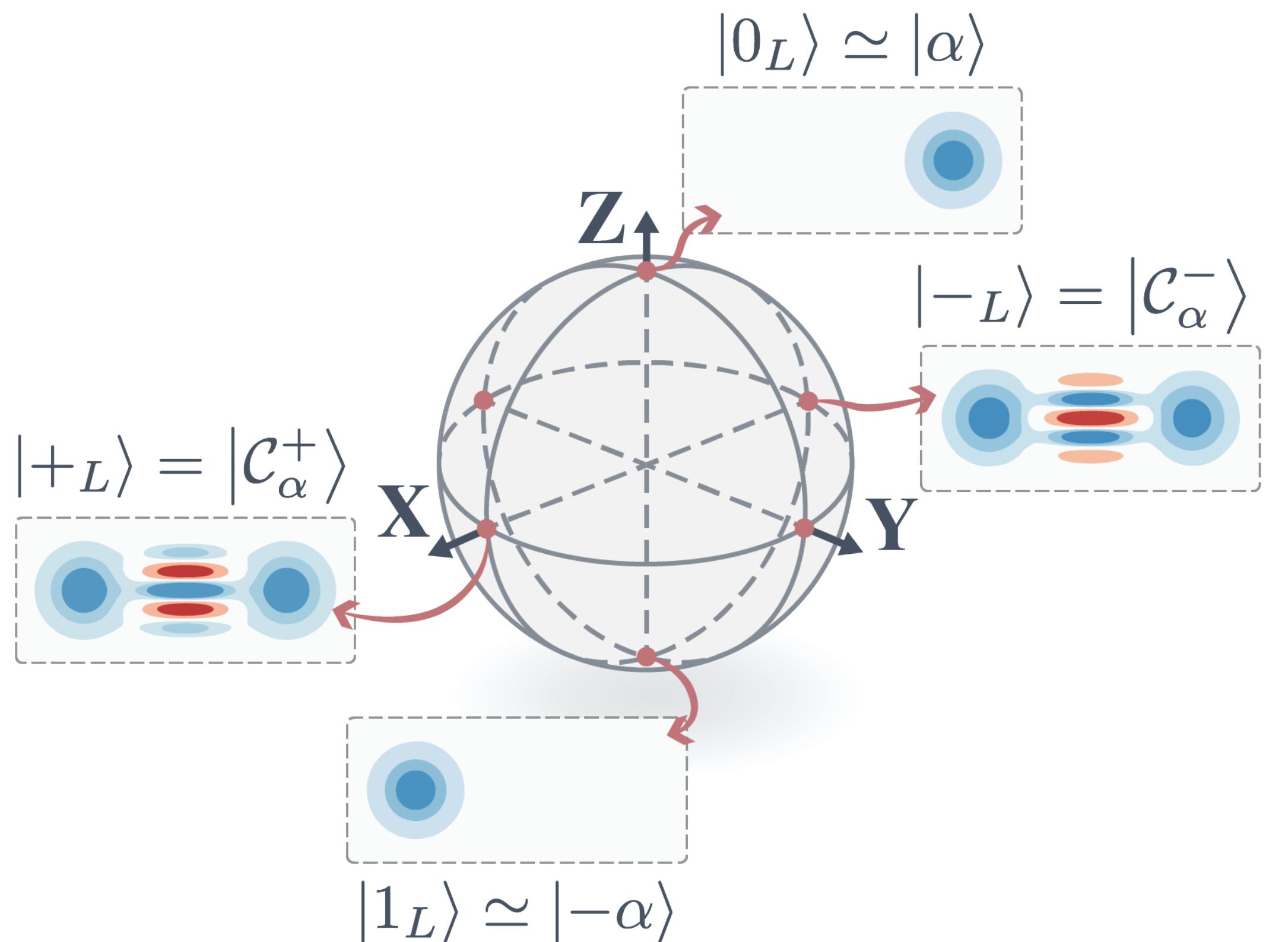
Cat qubits



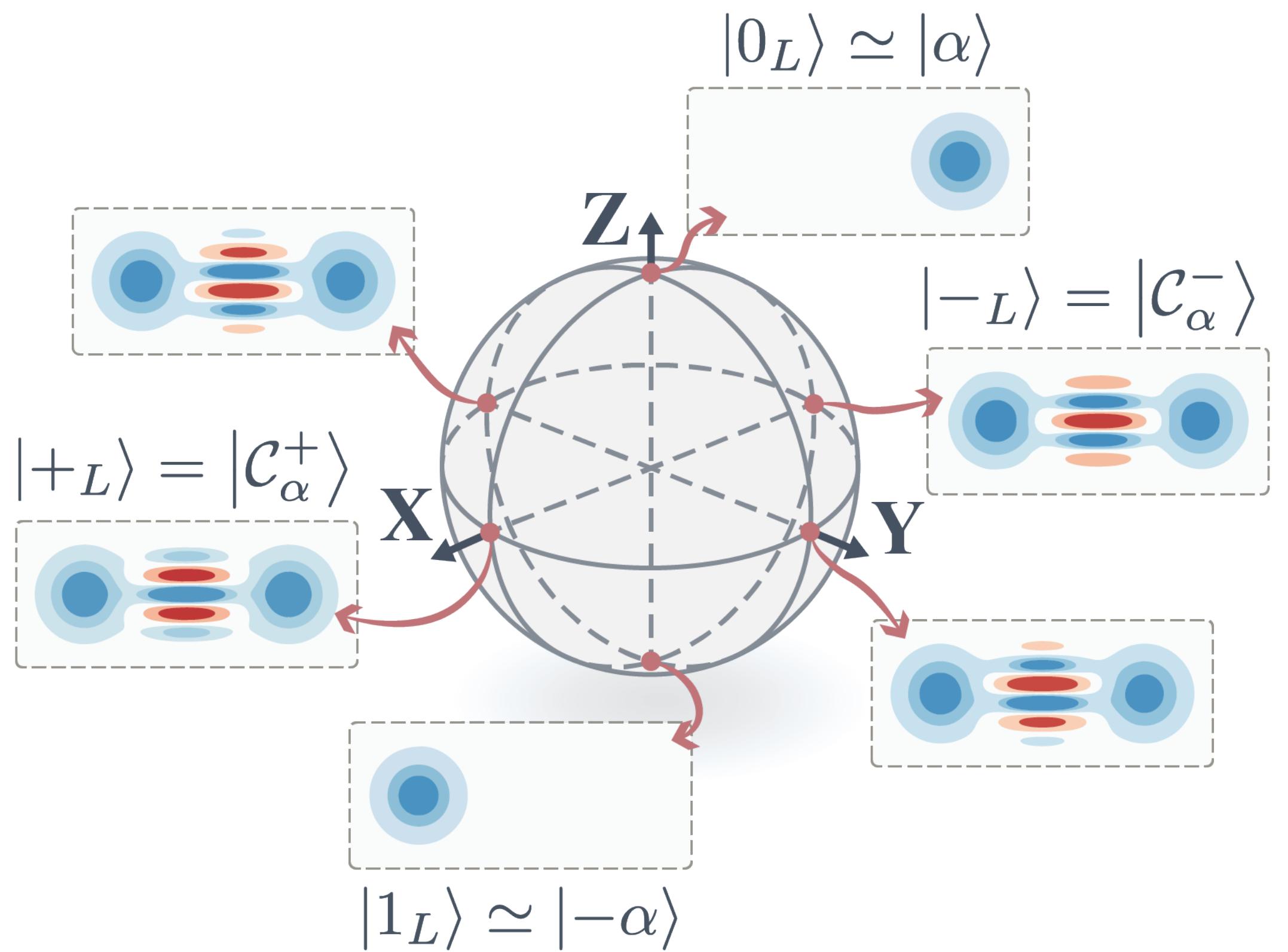
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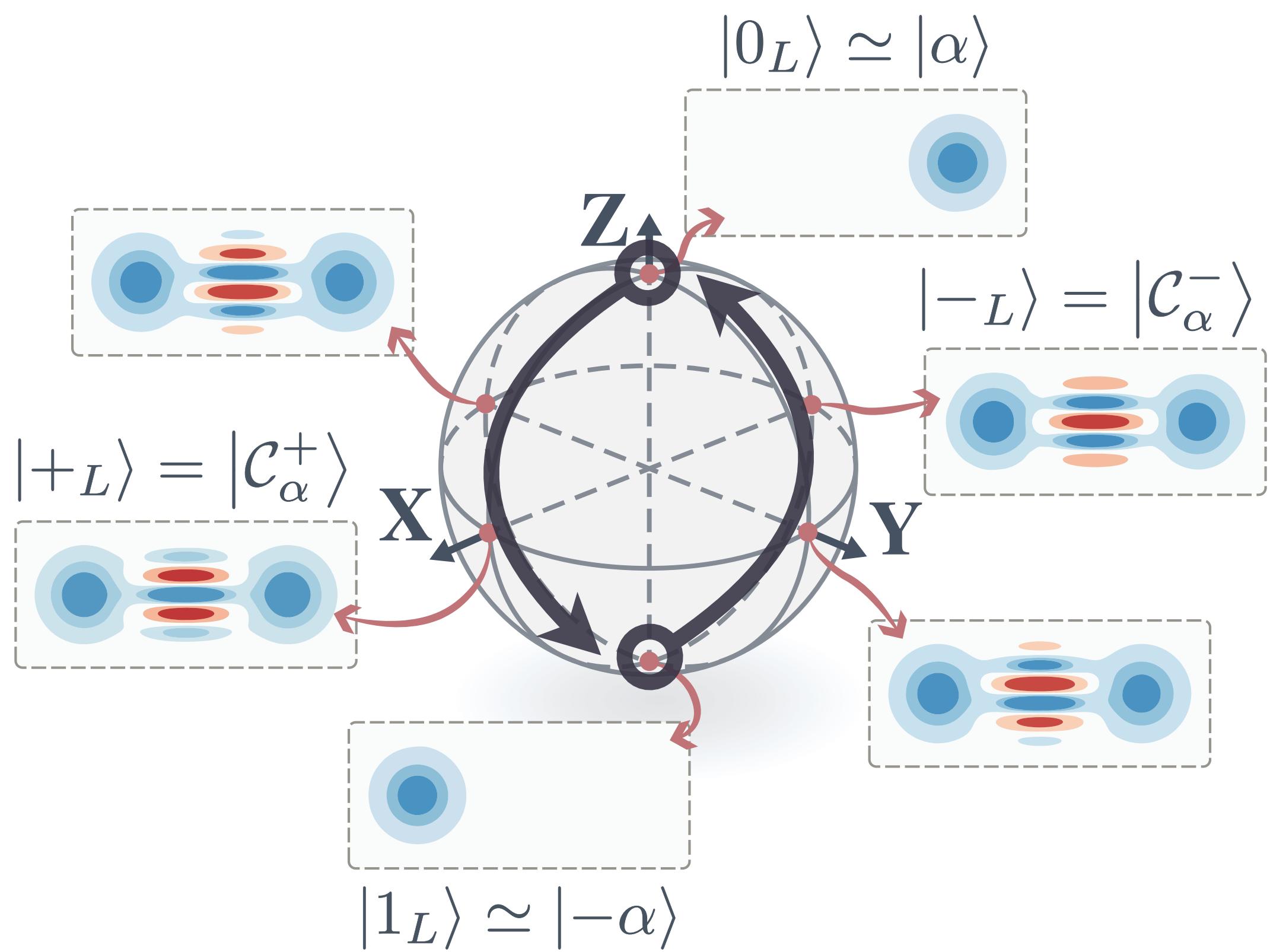
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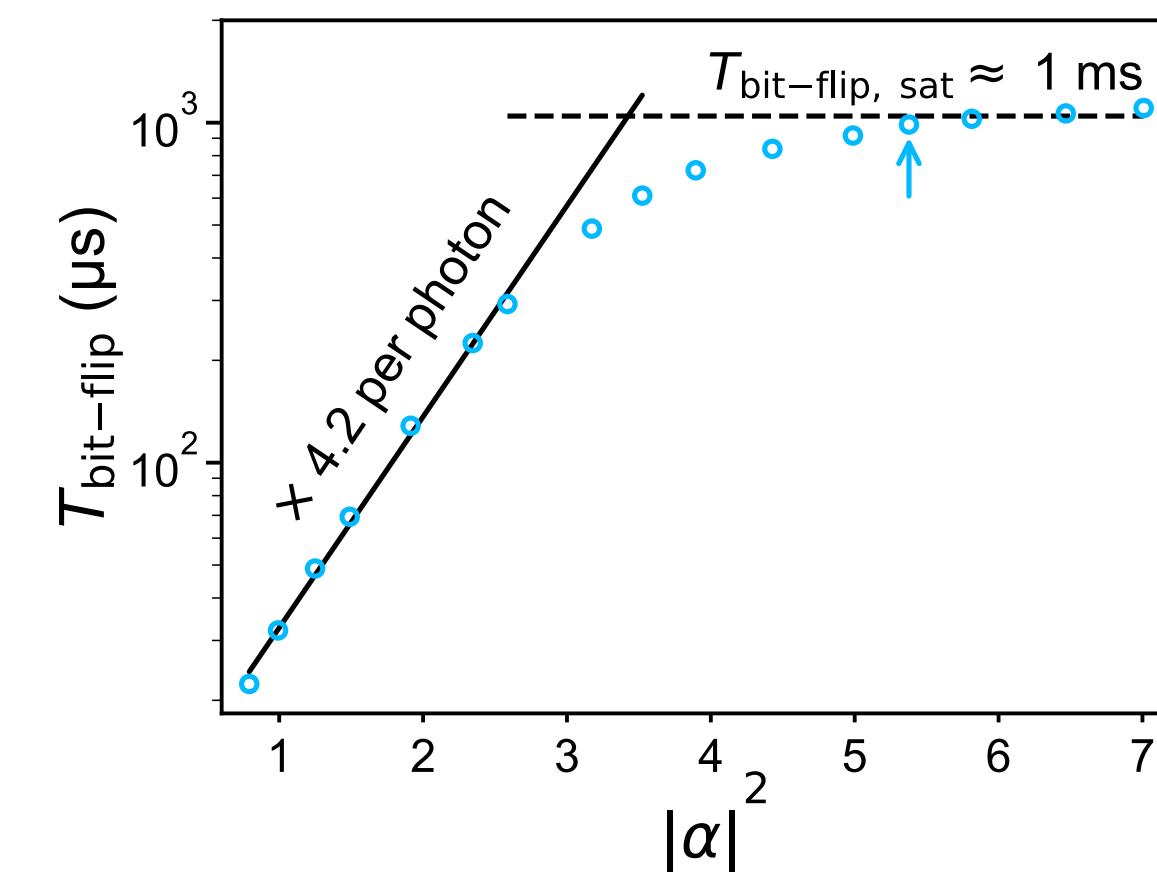
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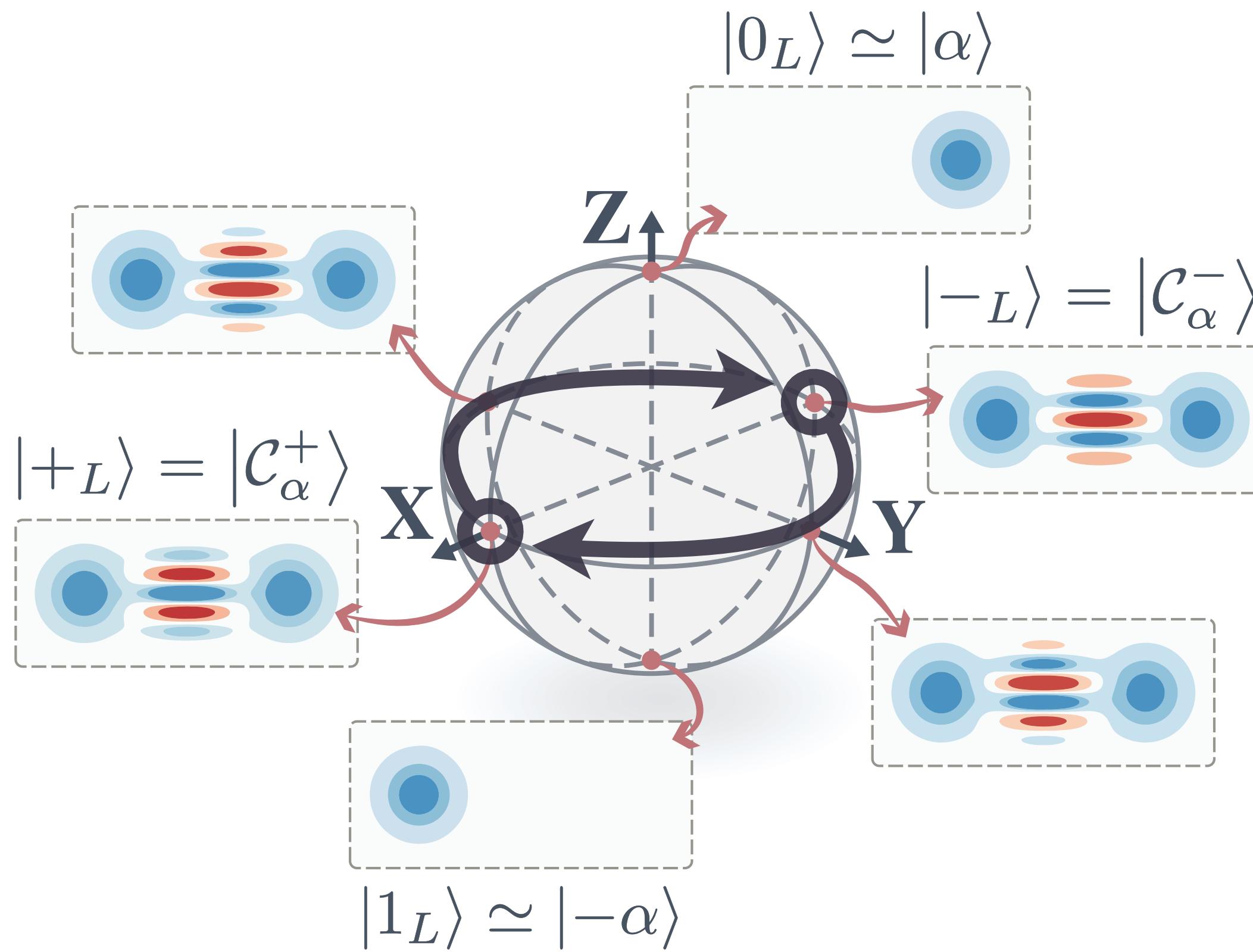
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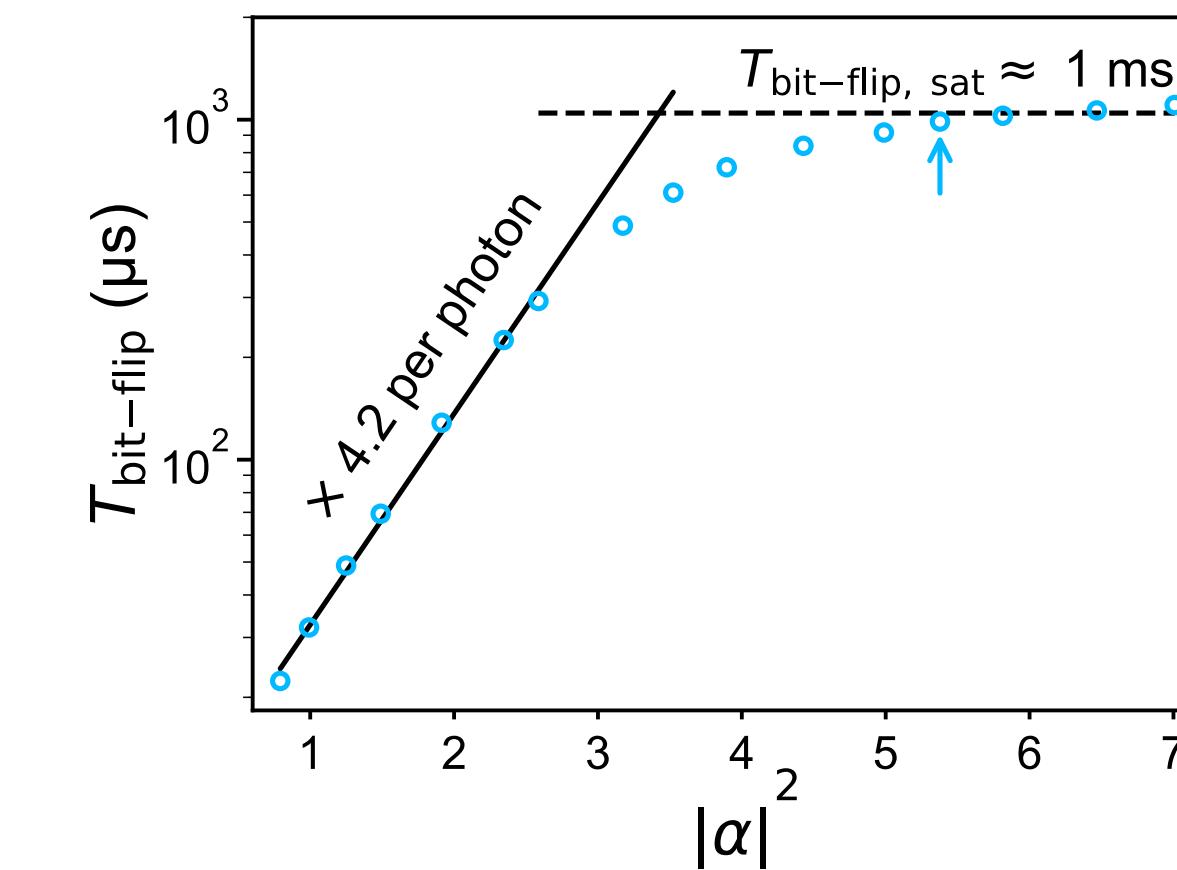
► Cat qubits are **exponentially** biased against bit-flip errors



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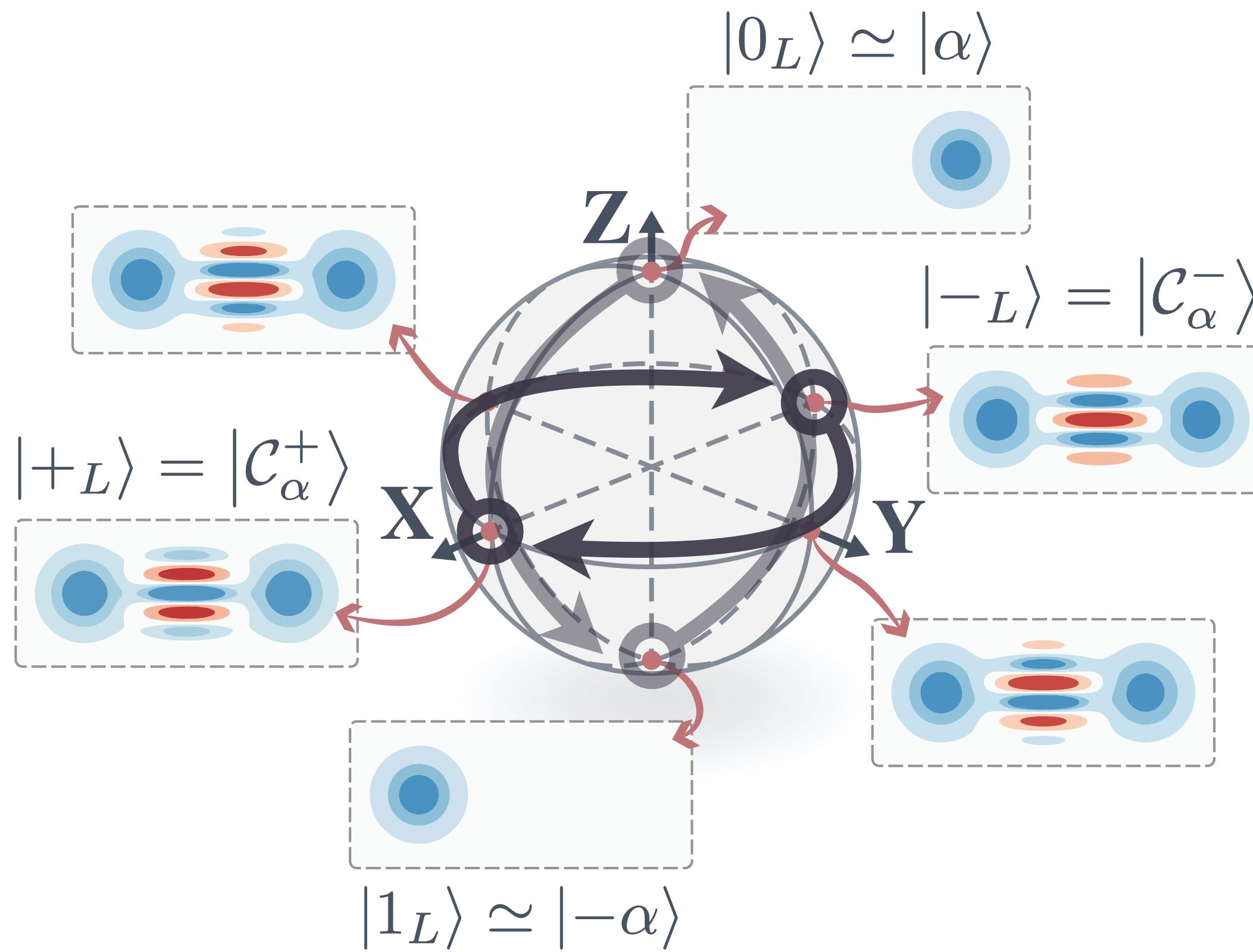
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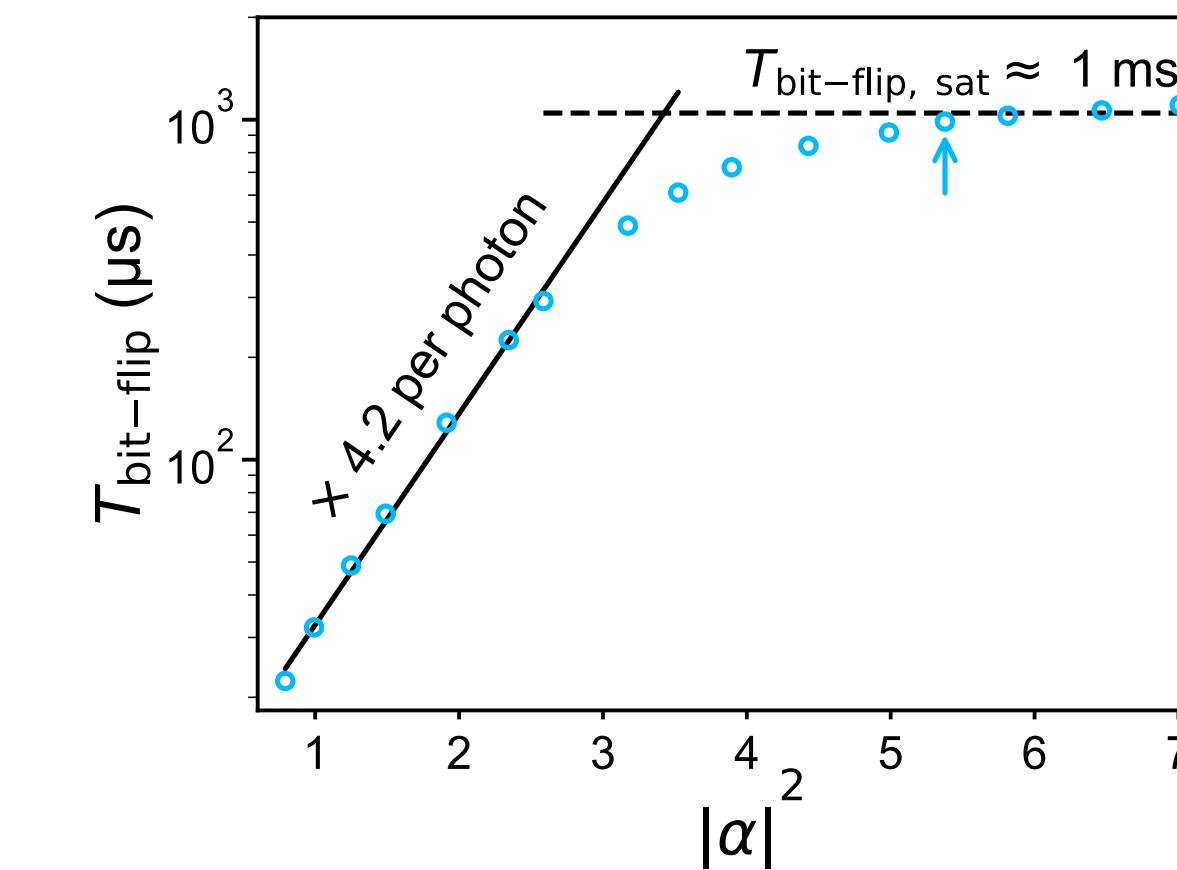
► A repetition code takes care of phase-flip errors



Cat qubits



- ▶ Cat qubits are **exponentially** biased against bit-flip errors



- ▶ A repetition code takes care of phase-flip errors



- ▶ Inner: cat qubits (bit-flips)
Outer: repetition code (phase-flips)

Protecting cat qubits

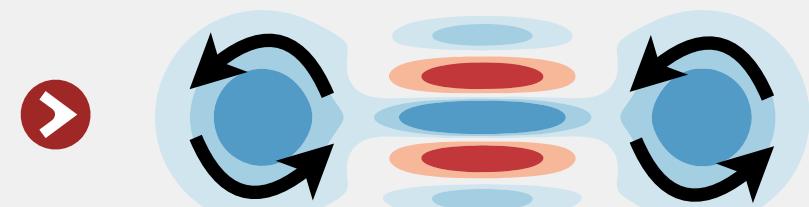
Kerr cat qubits

- Hamiltonian confinement

$$H = -K(a^{\dagger 2} - \alpha^{*2})(a^2 - \alpha^2)$$

since $(a^2 - \alpha^2)|\pm\alpha\rangle = 0$

- Kerr non-linearity
+ two-photon driving



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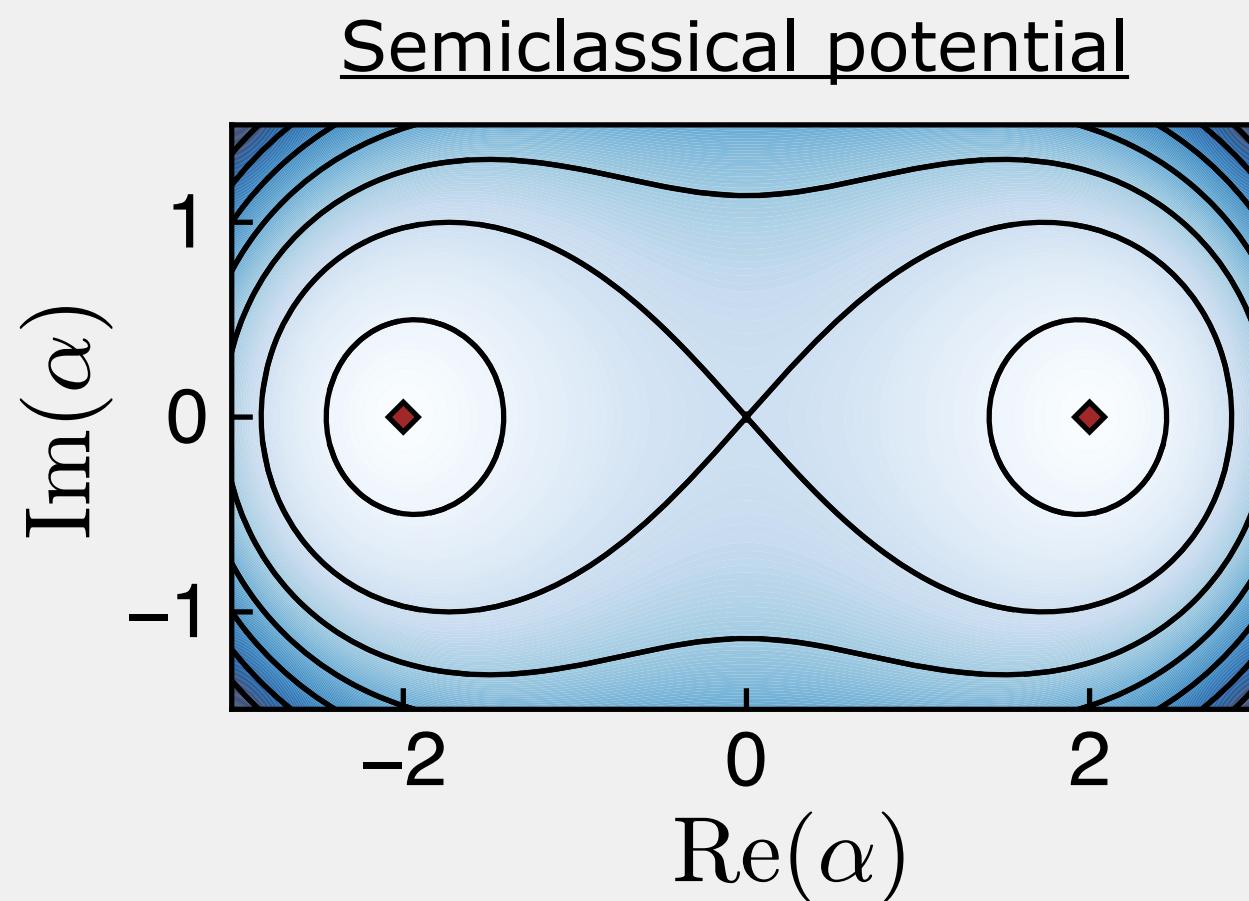
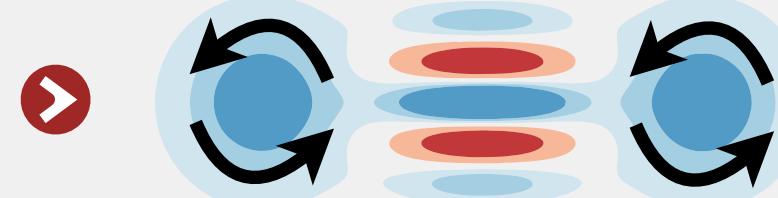
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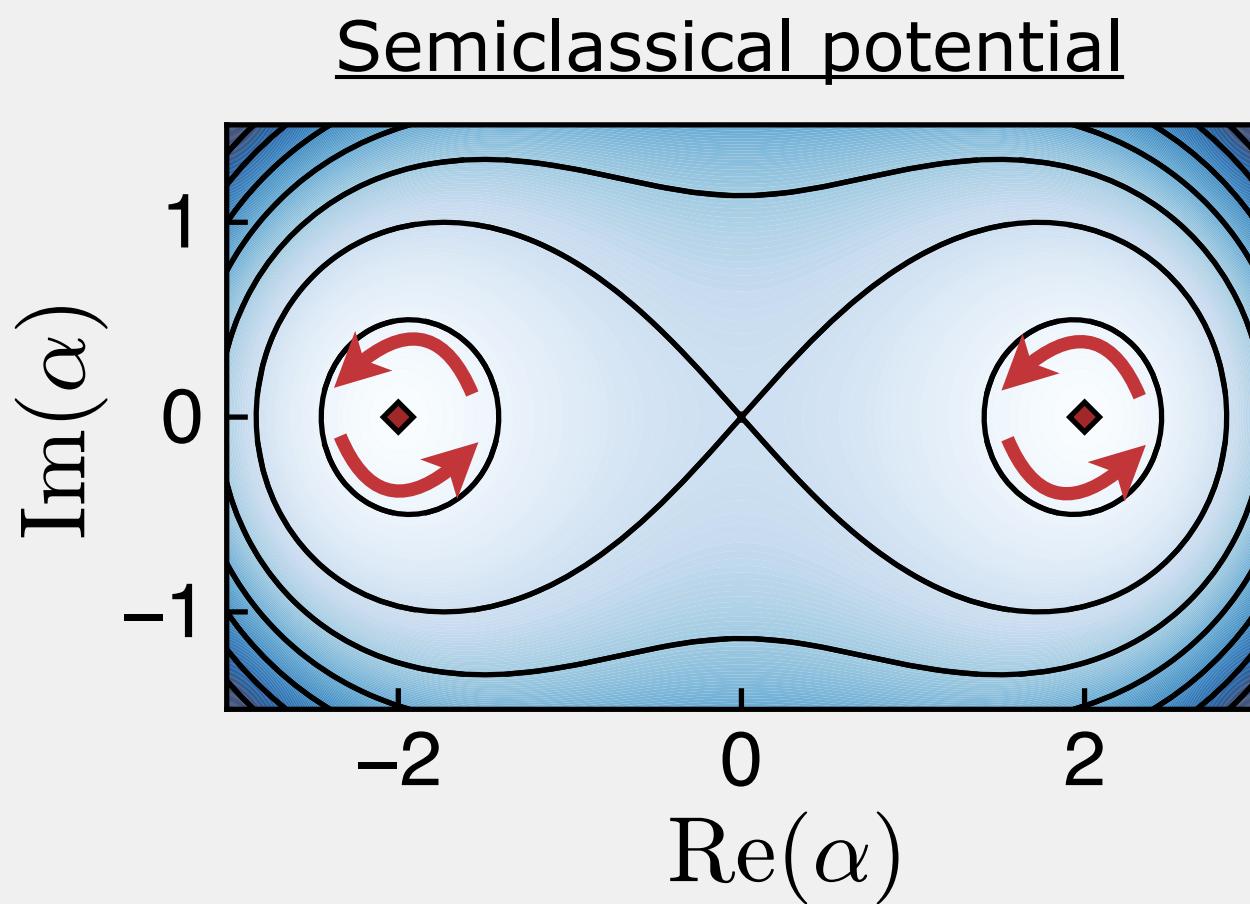
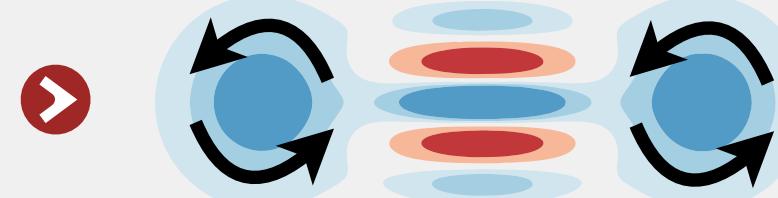
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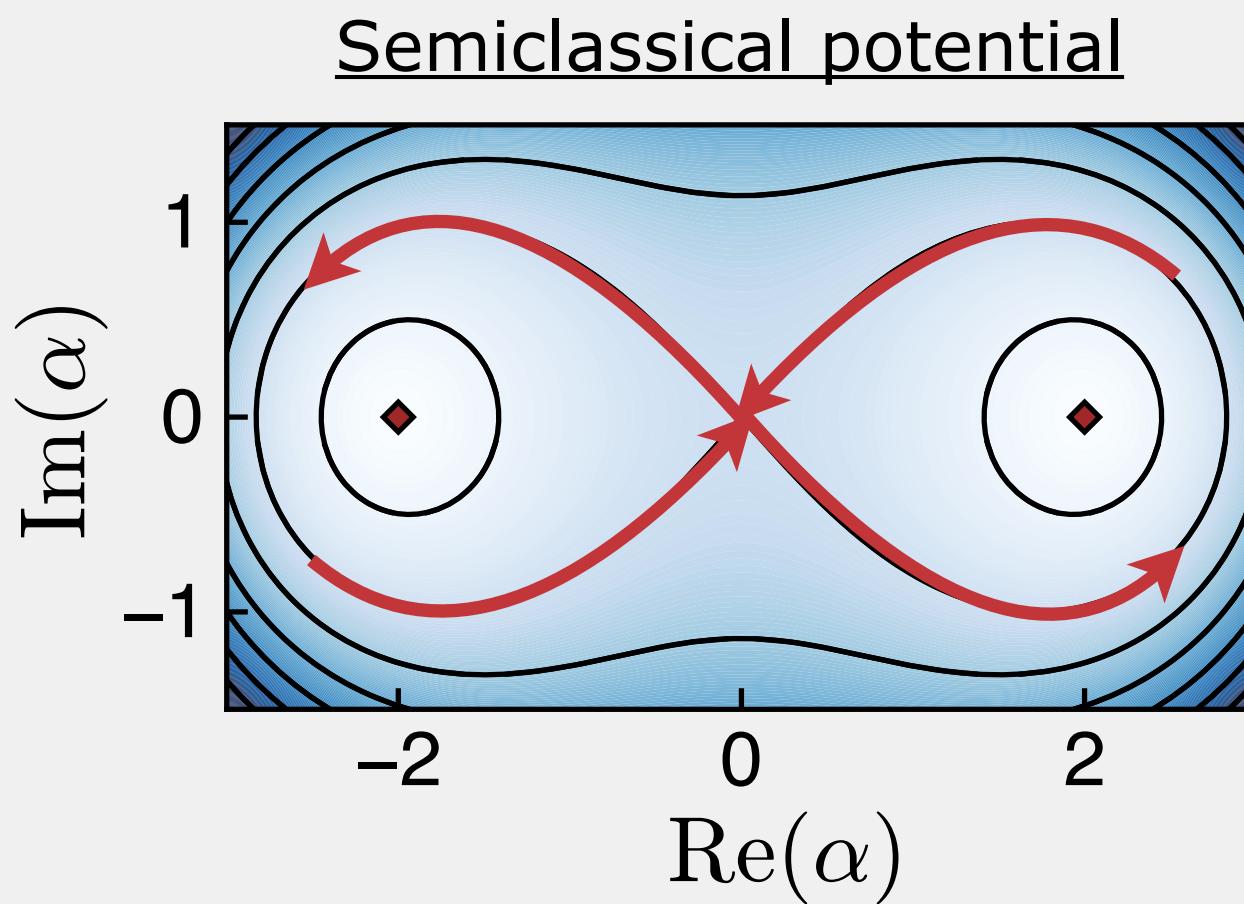
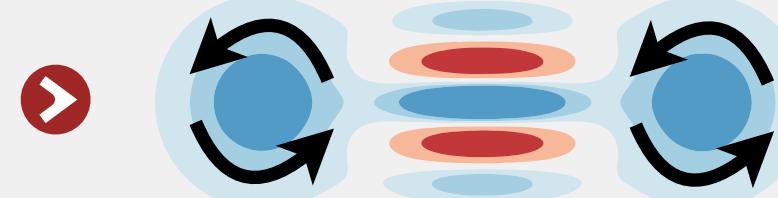
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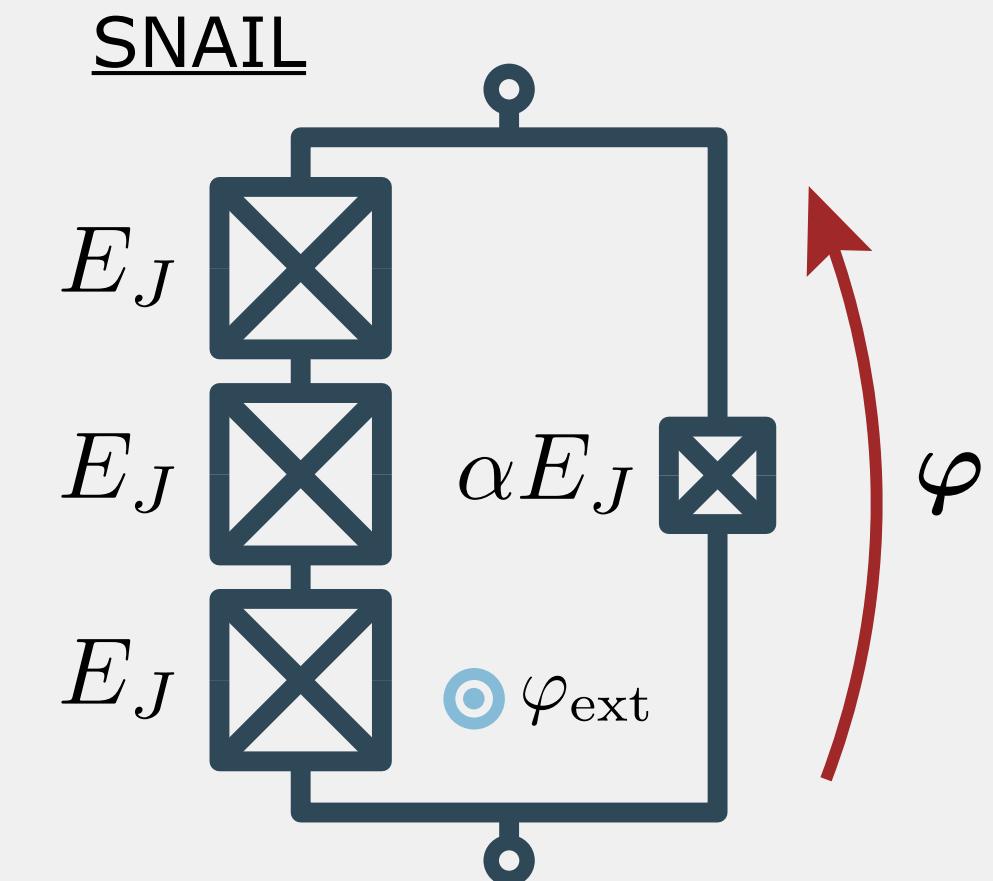
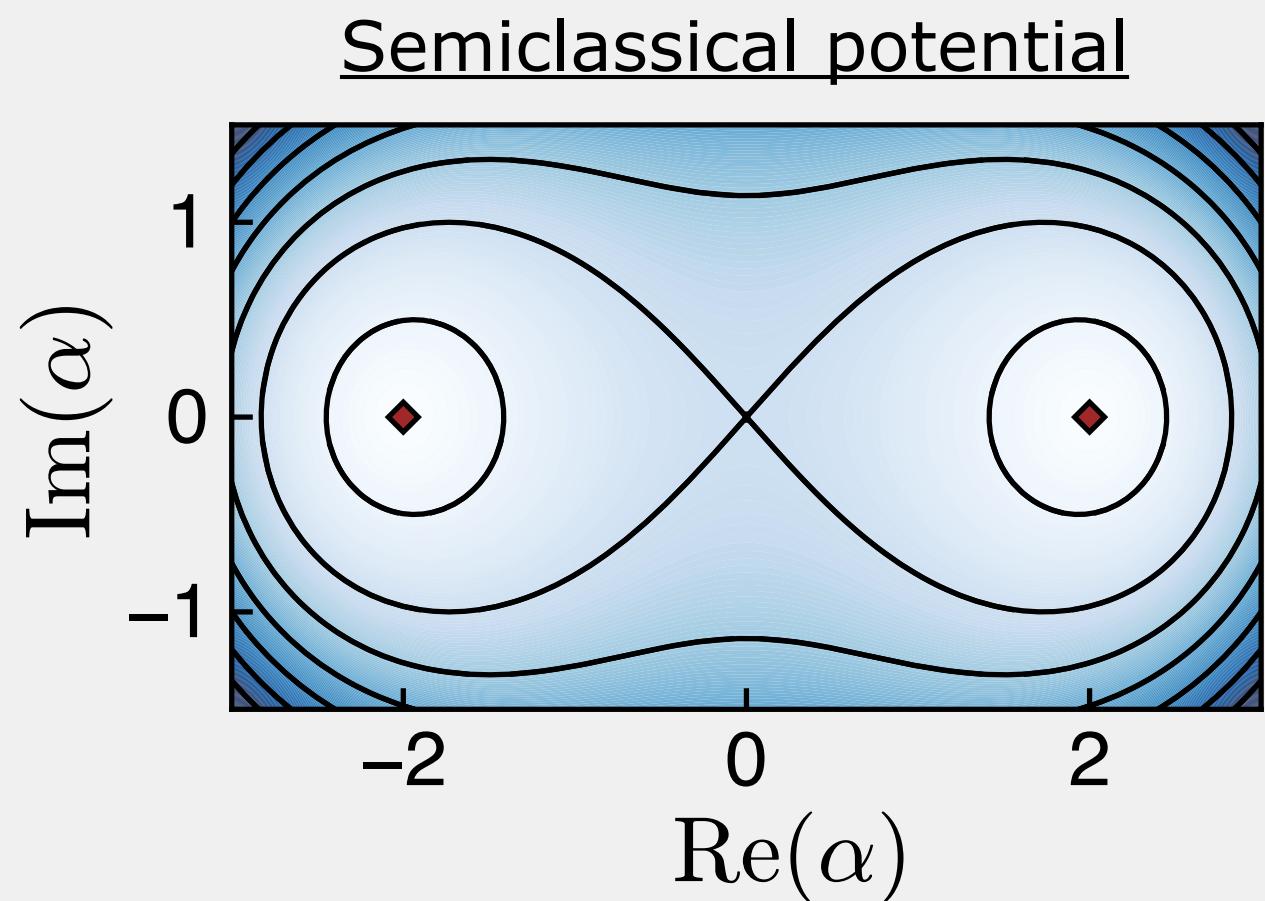
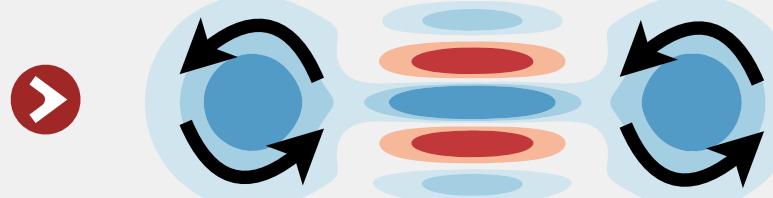
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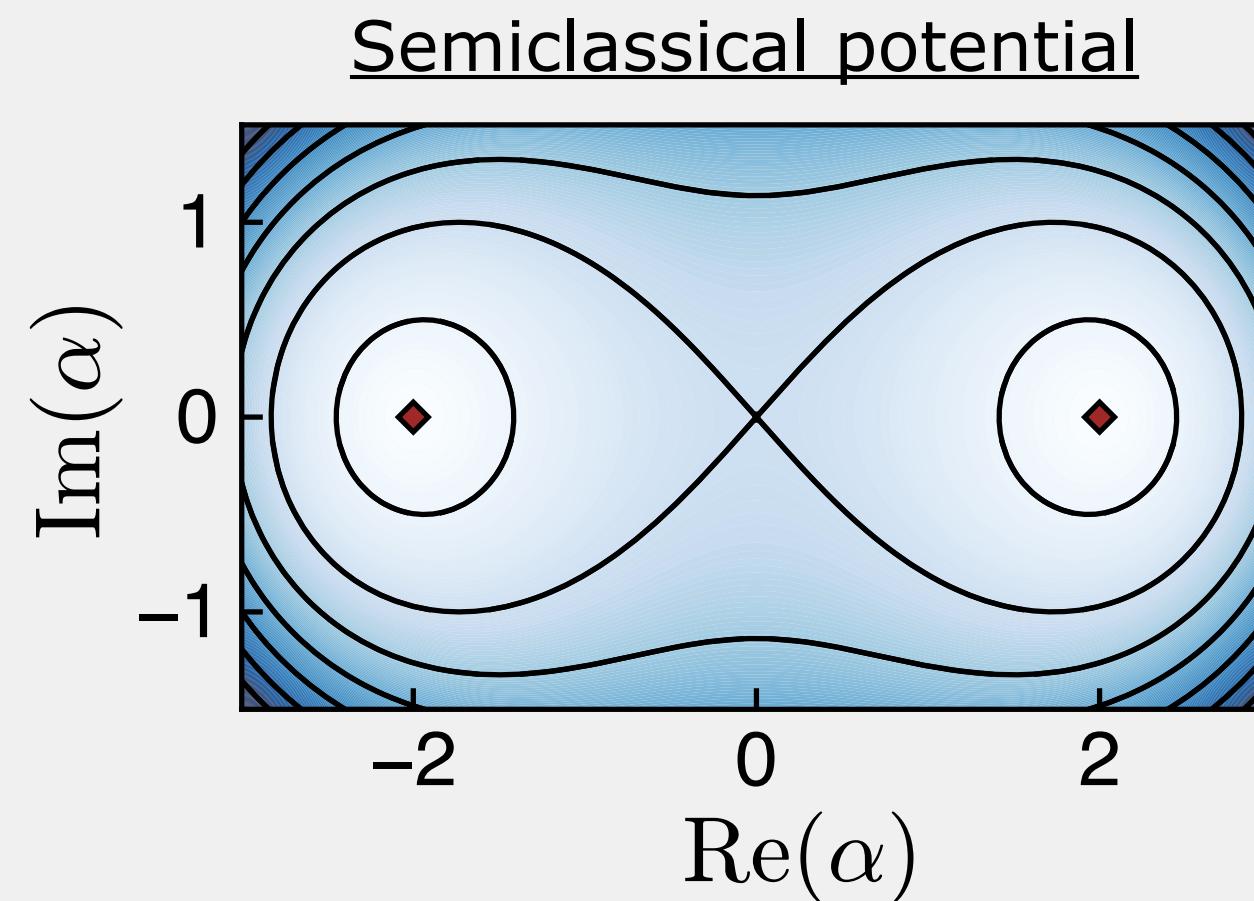
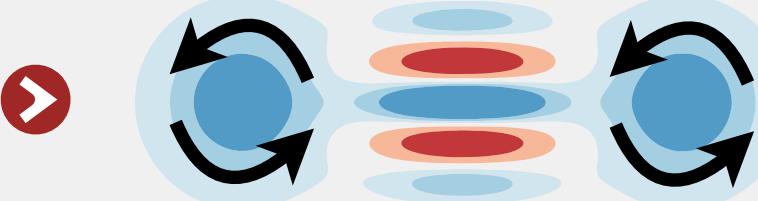
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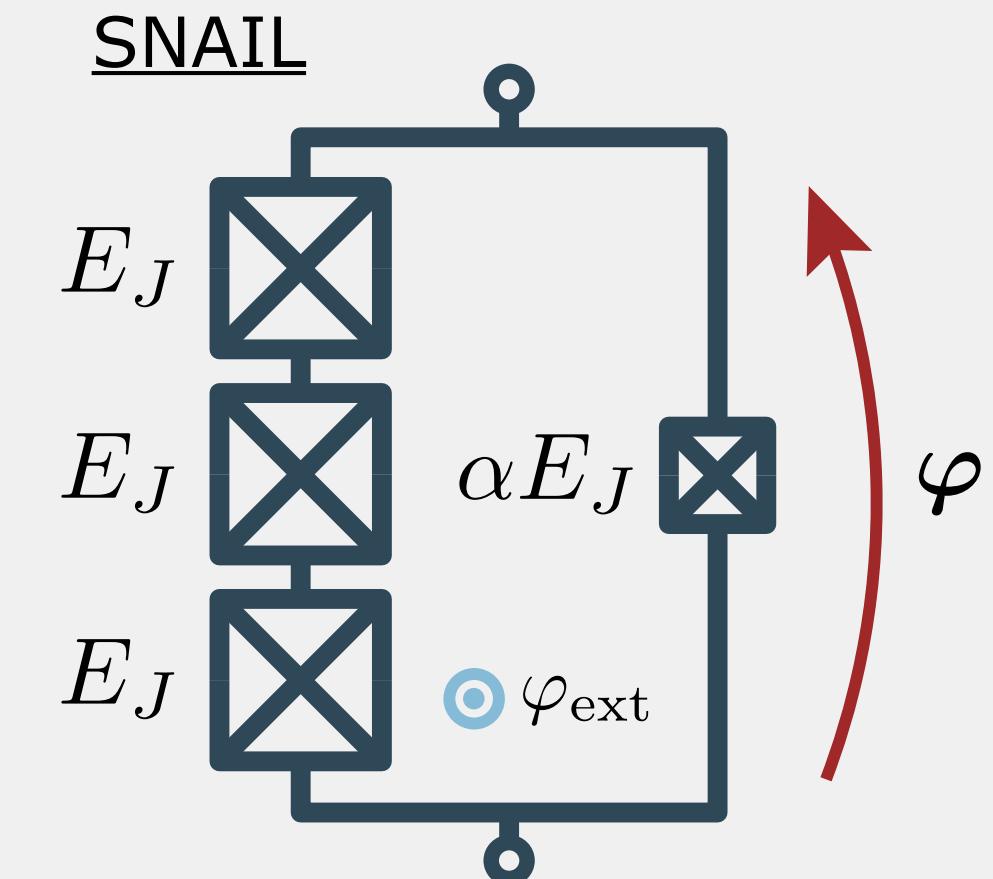
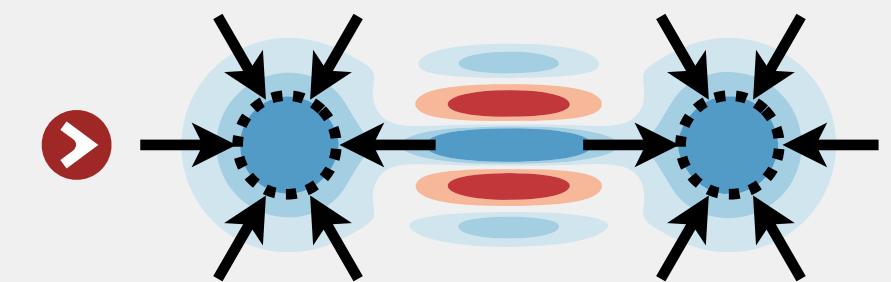


Dissipative cat qubits

- ▶ Dissipative stabilization

$$\kappa_2 \mathcal{D}[a^2 - \alpha^2]$$

- ▶ Two-photon dissipation + two-photon driving



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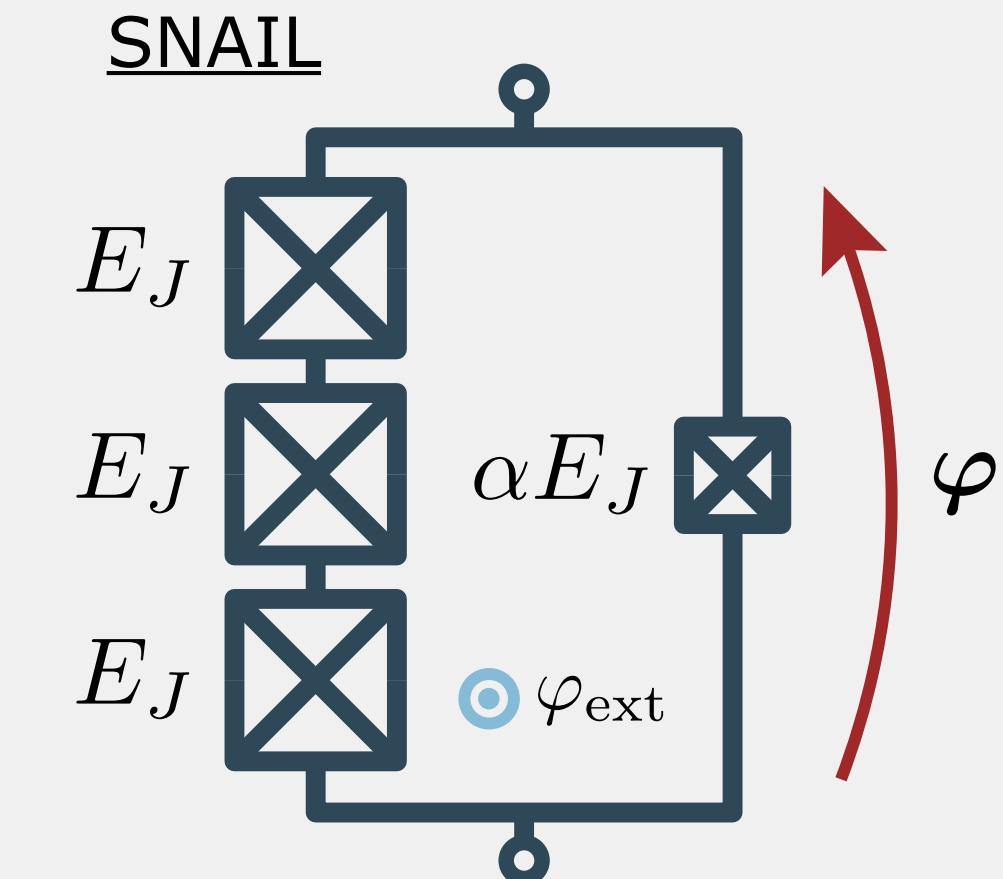
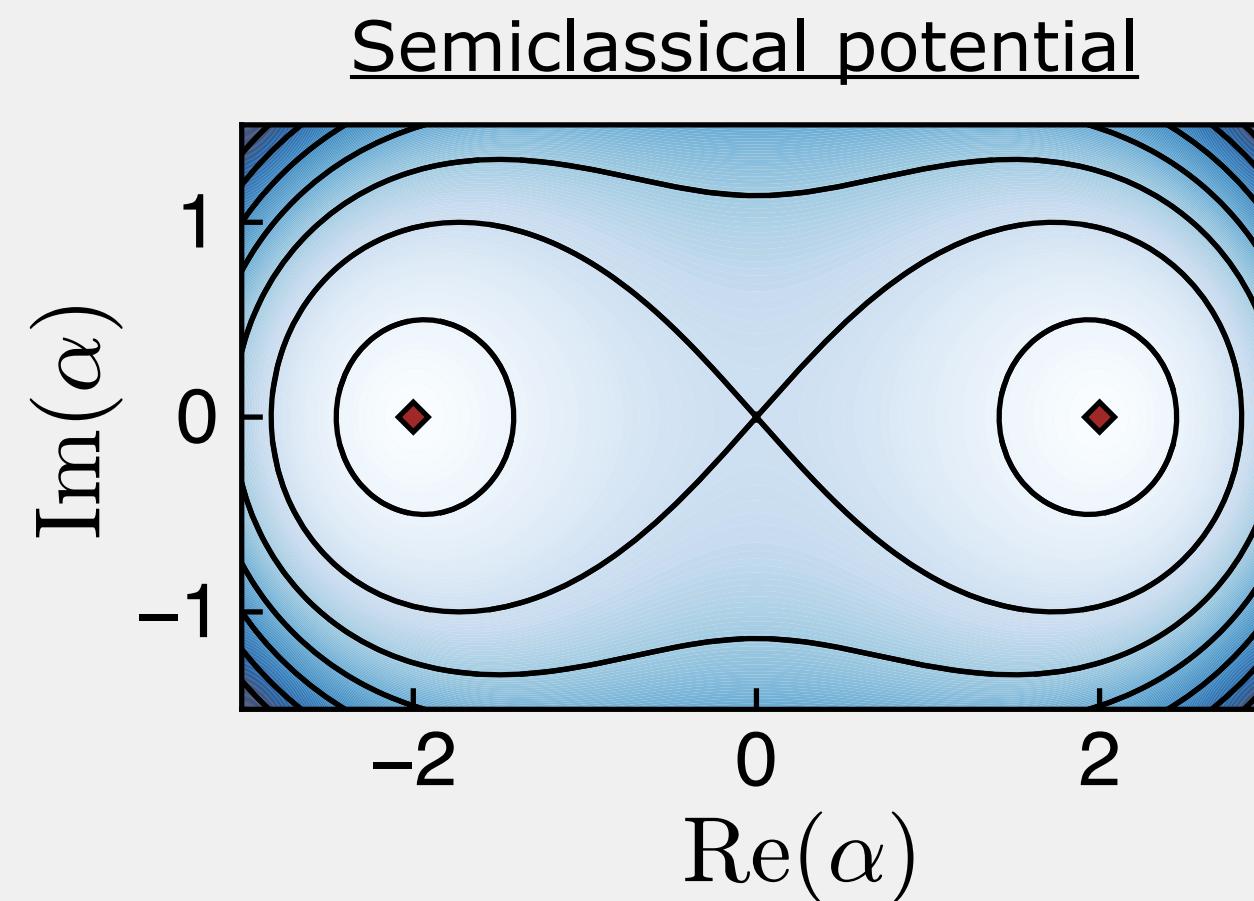
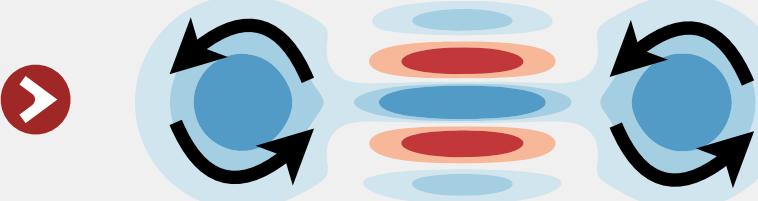
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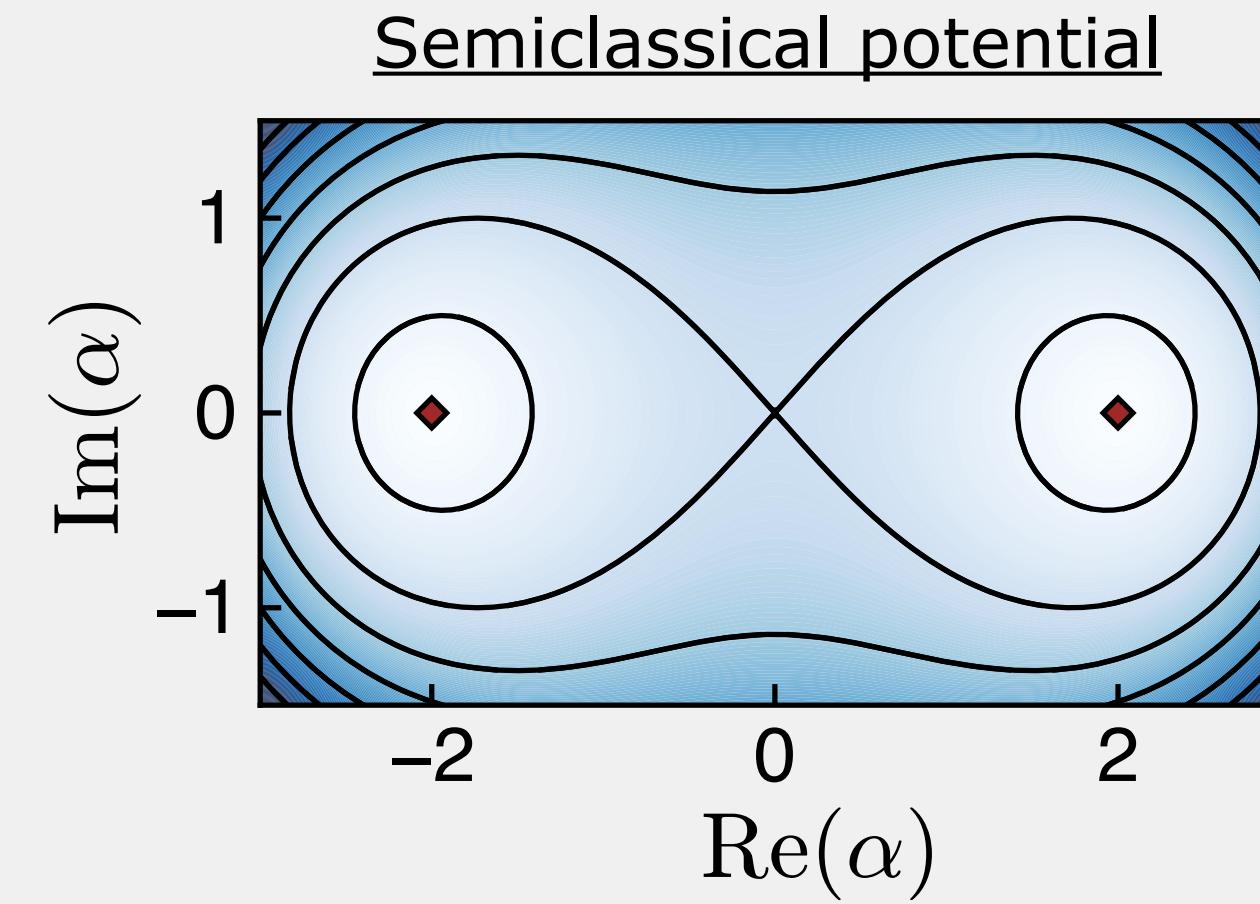
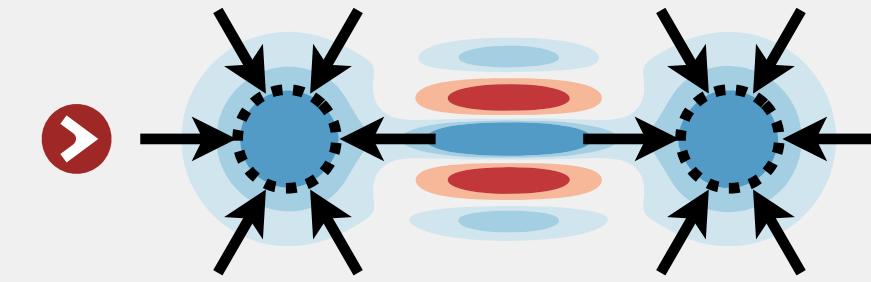


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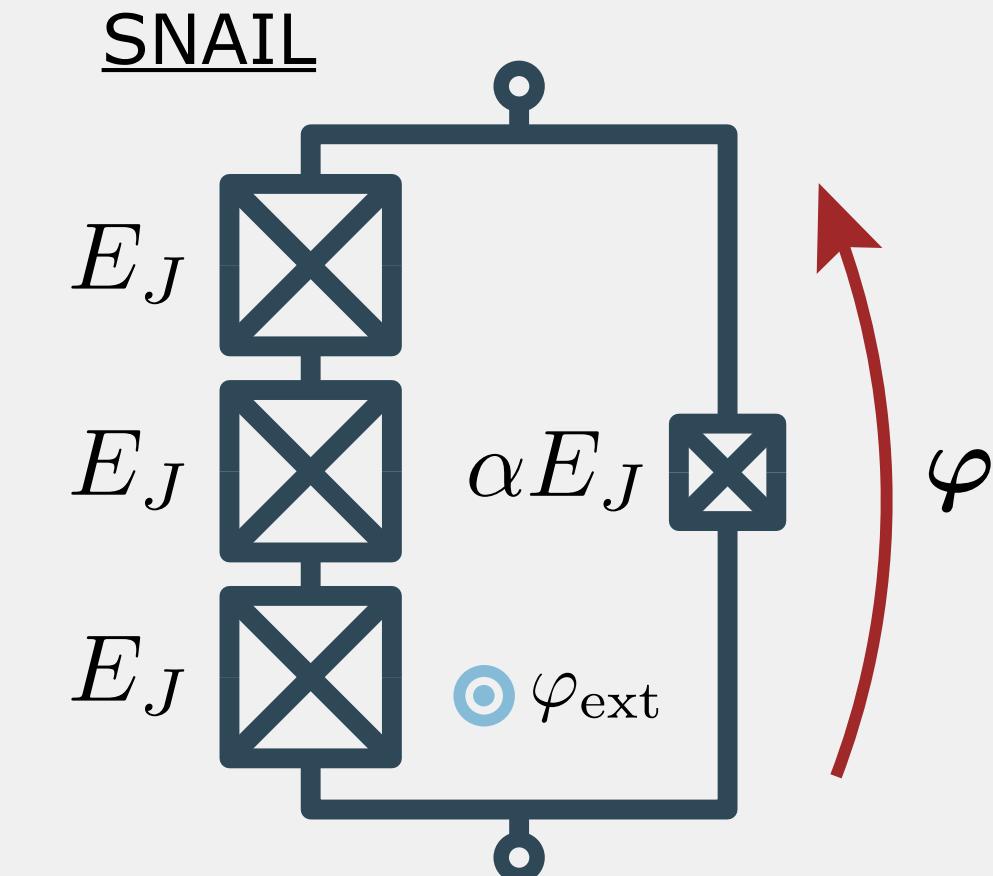
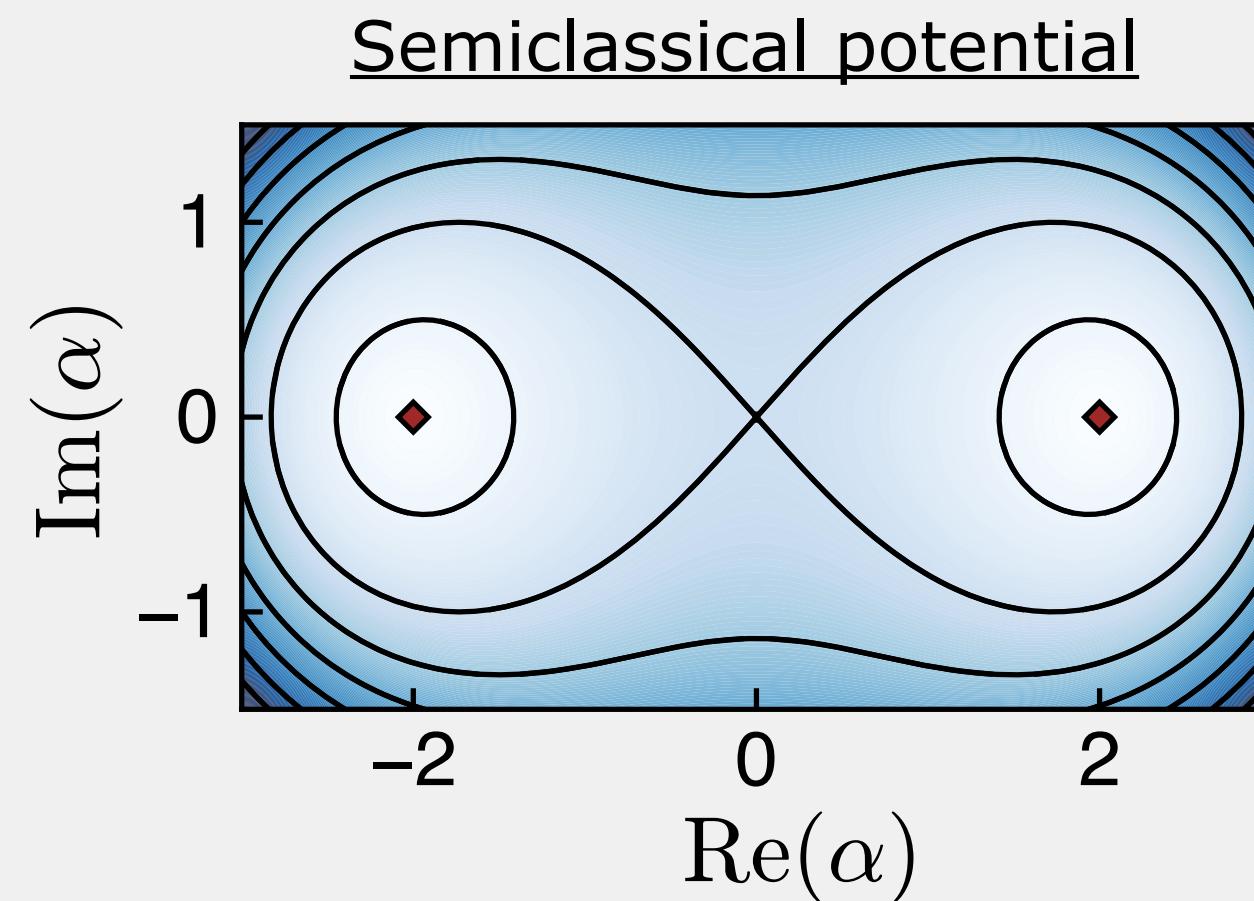
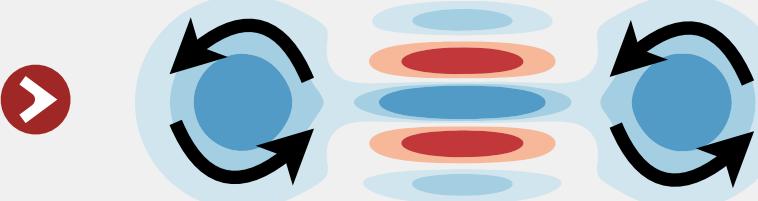
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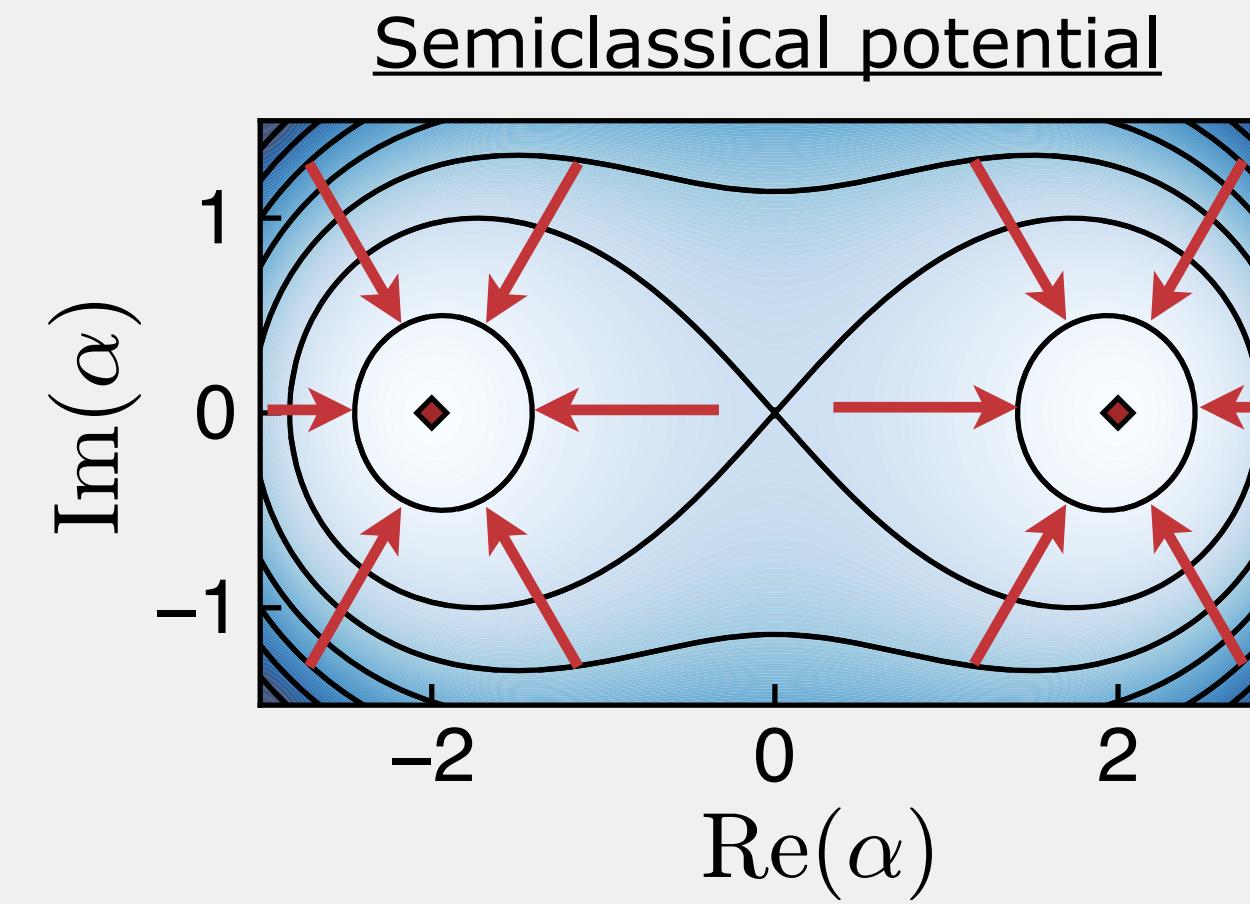
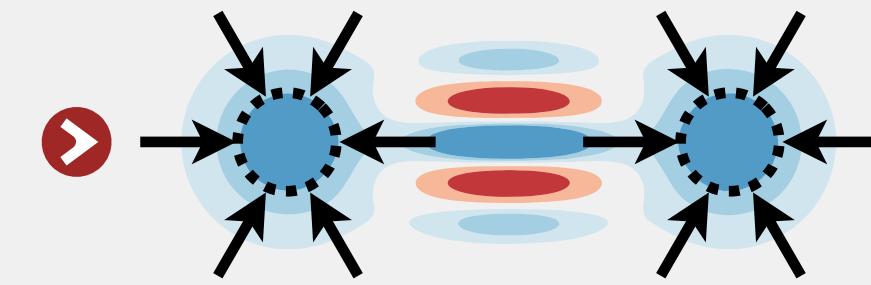


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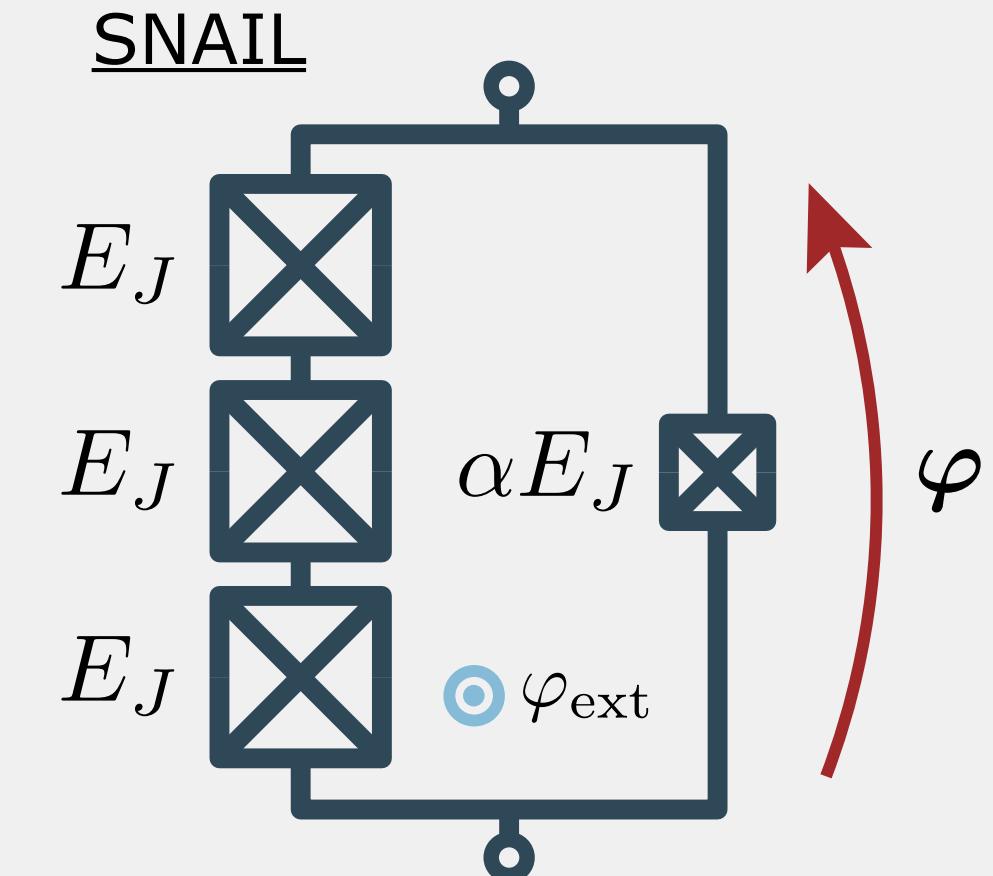
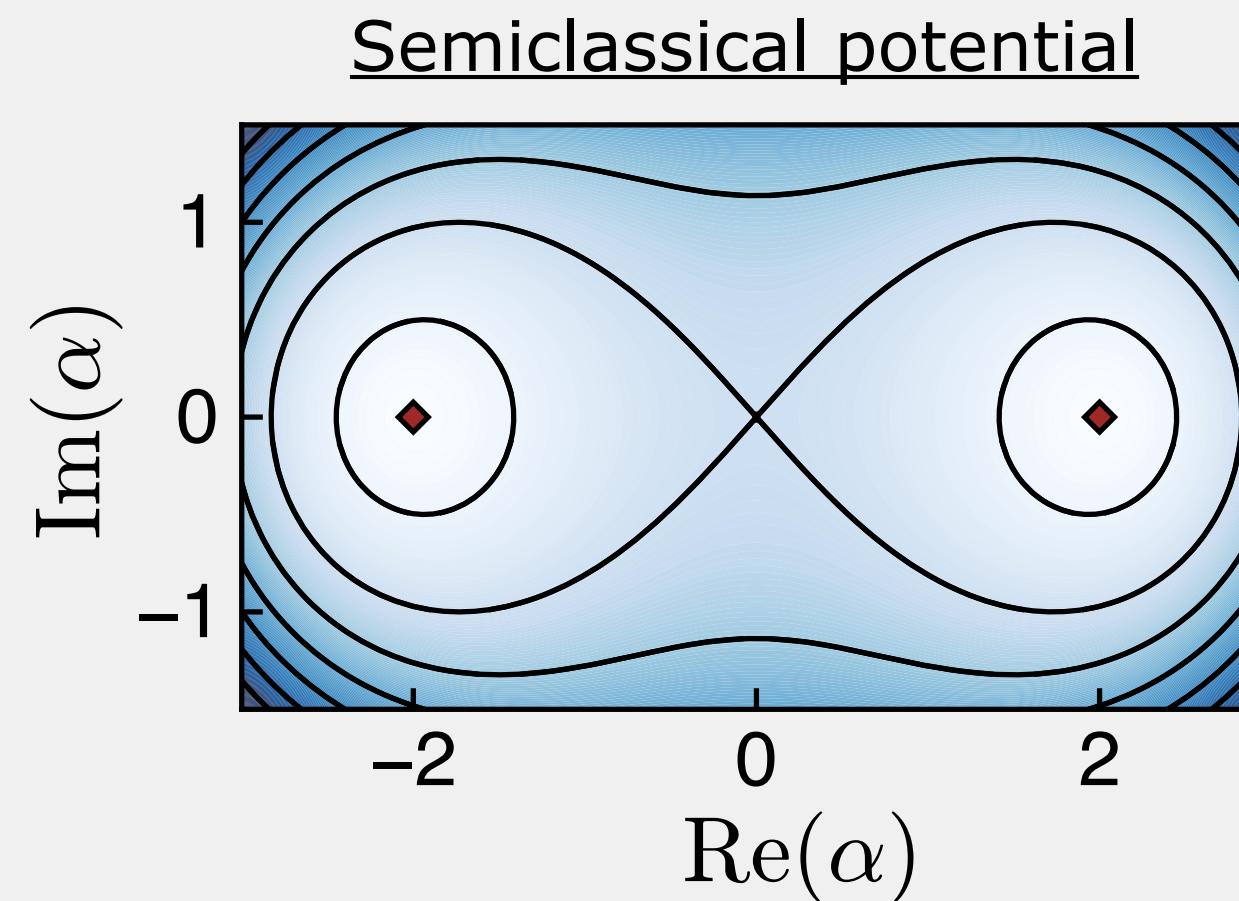
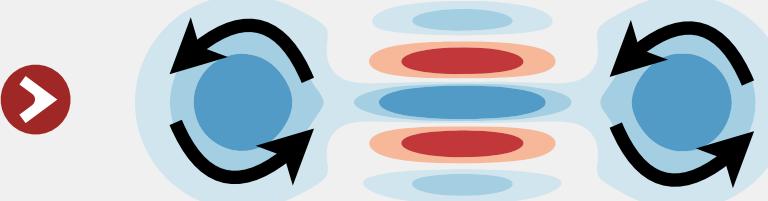
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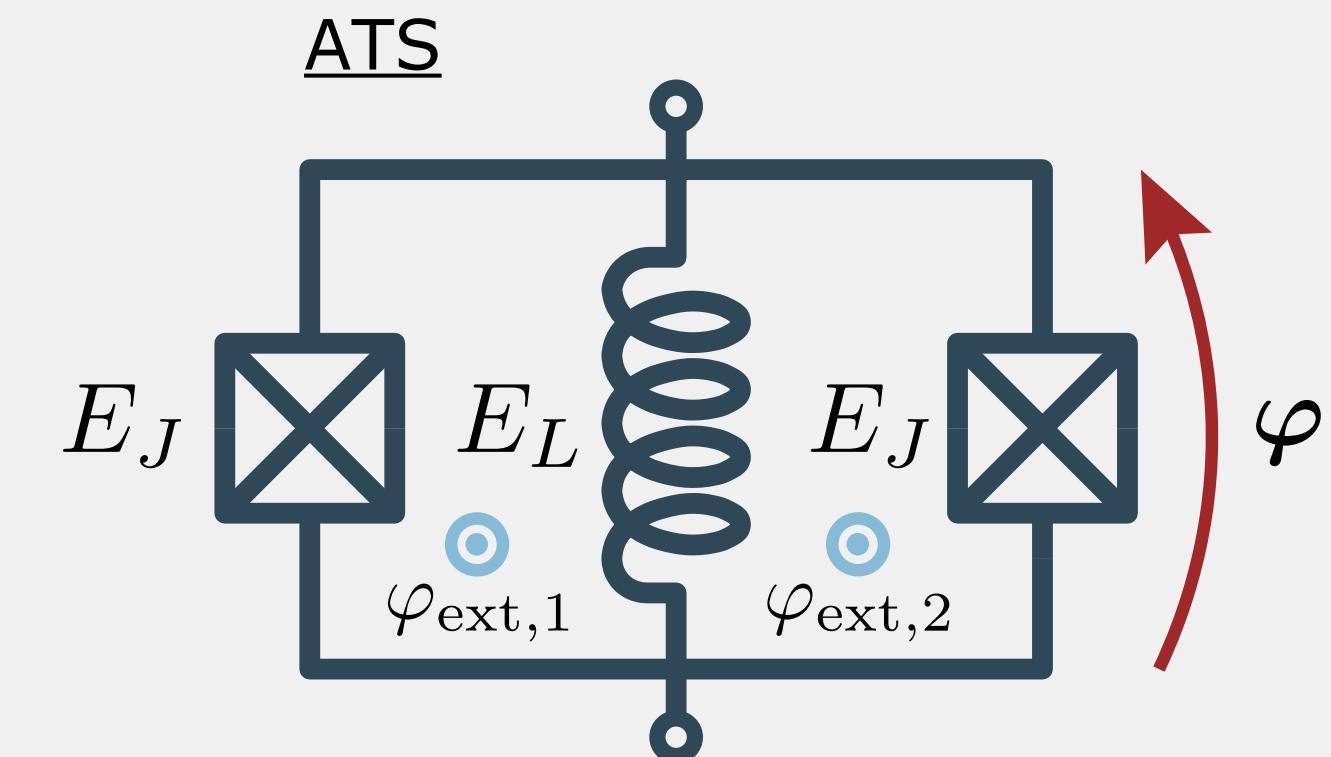
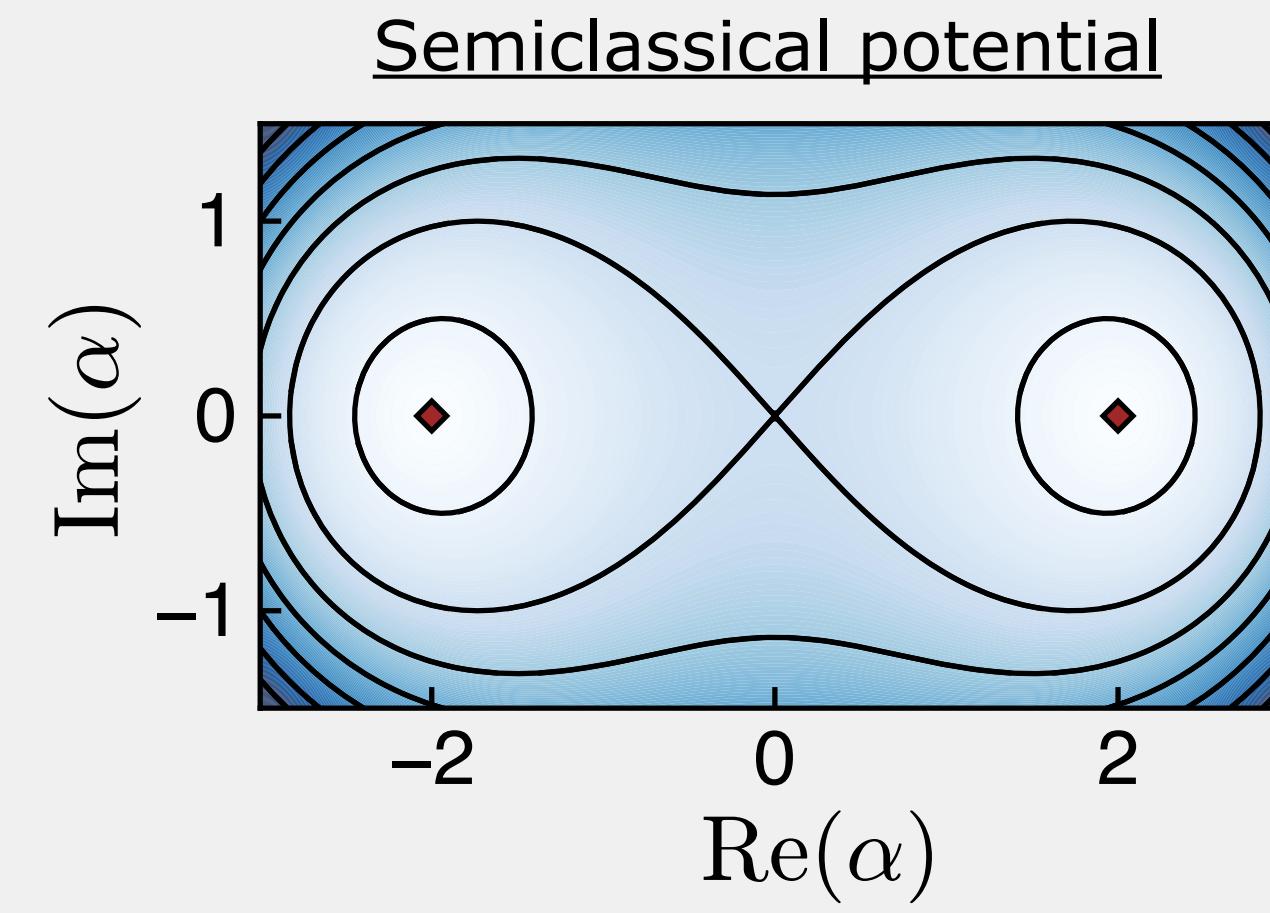
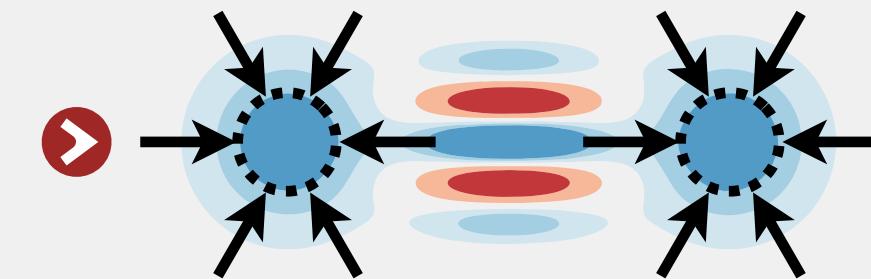


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Fault-tolerant universal QC

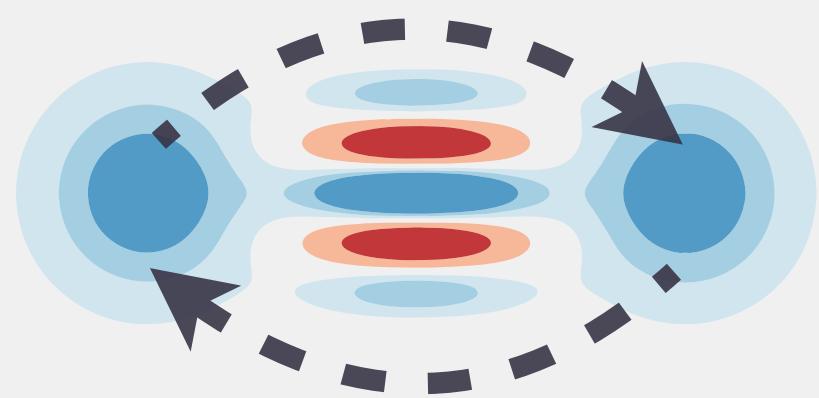
- Preparation (X)
 - Measurement (X)
 - Bias-preserving gates
- Guillaud et al. (2019)

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Pauli X



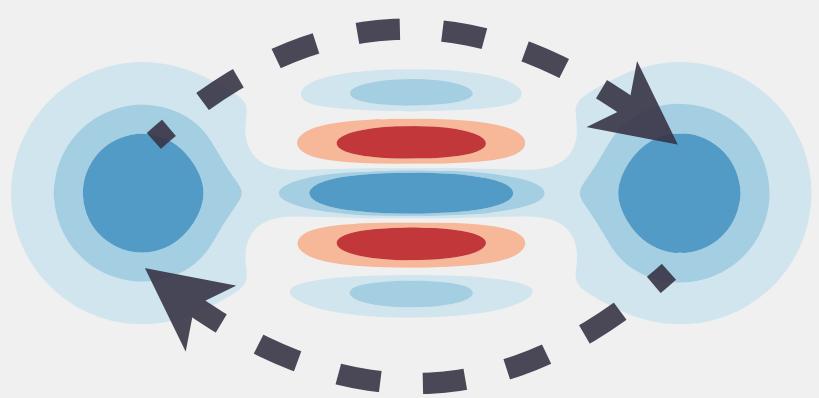
$$H = \Delta_X a^\dagger a$$

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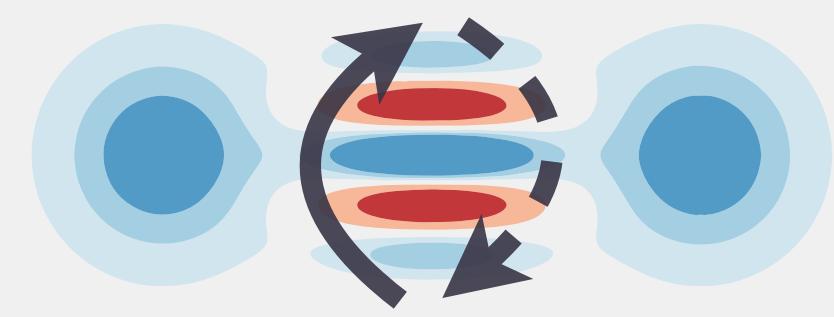
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Z rotation

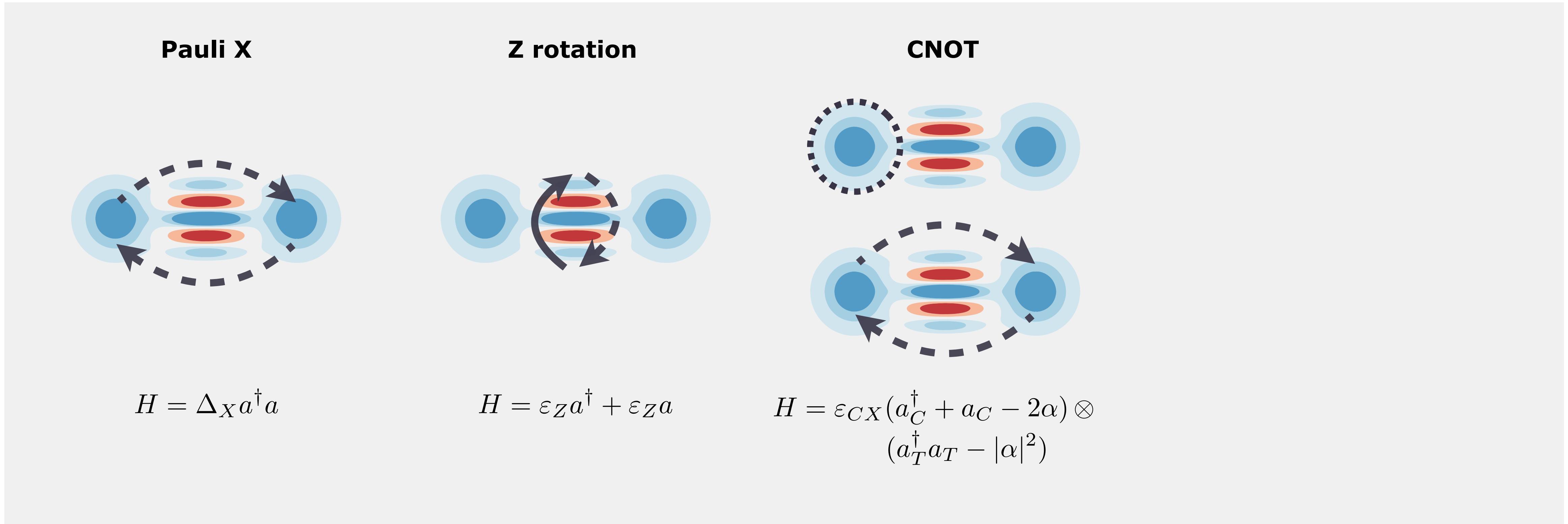


$$H = \varepsilon_Z a^\dagger + \varepsilon_Z a$$

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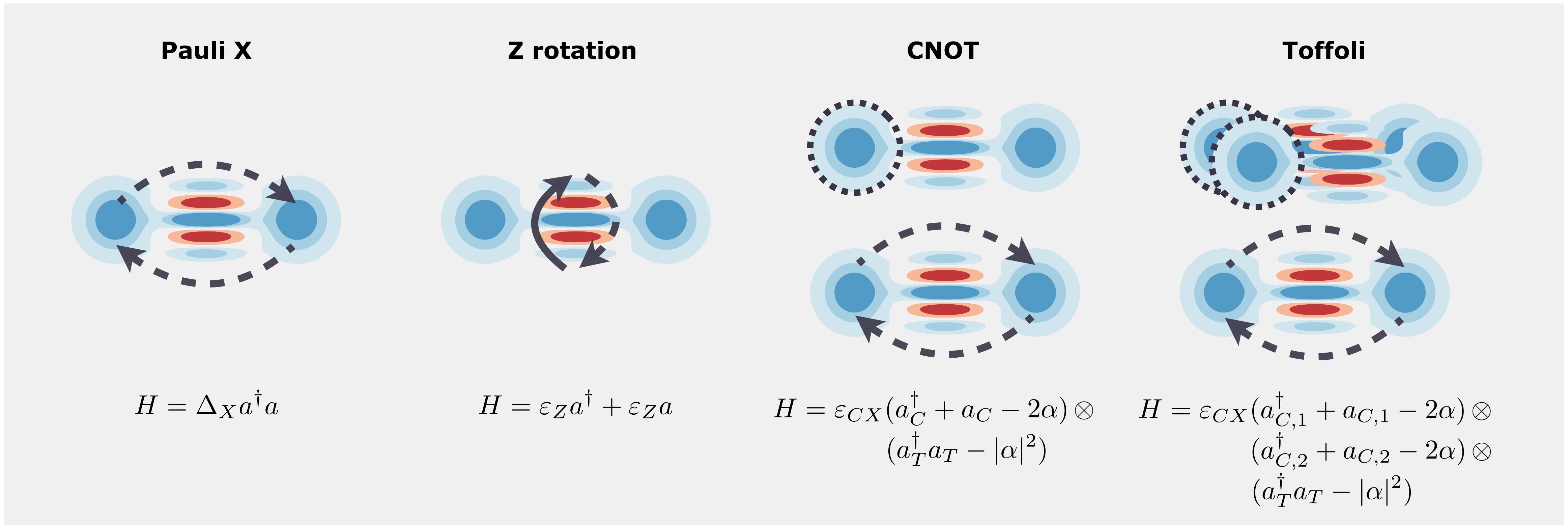
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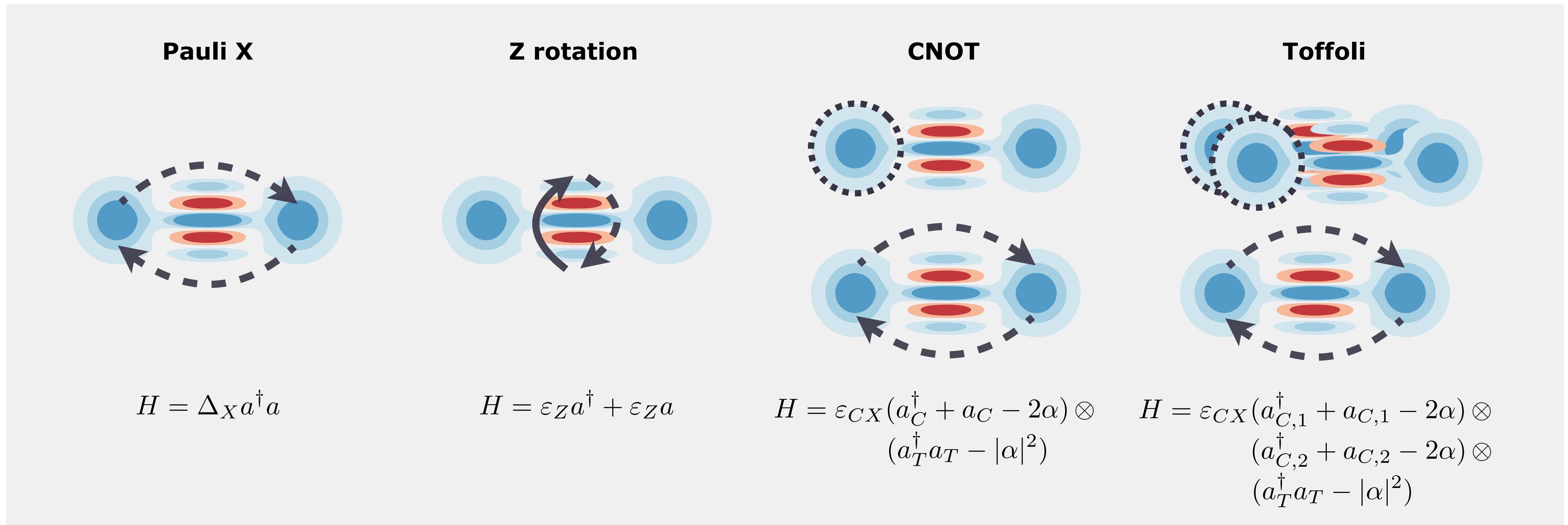
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- ▶ For Kerr cat qubits → $1 - F \propto \exp(-\gamma T_{\text{gate}})$ (adiabatic theorem)
- ▶ For dissipative cat qubits → $1 - F \propto 1/T_{\text{gate}}$ (Zeno effect)

Summary of PhD contributions

Work on cat qubits

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- D. Ruiz, RG, J. Guillaud, M. Mirrahimi, *Two-photon driven Kerr quantum oscillator with multiple spectral degeneracies*, Phys. Rev. A (2022)
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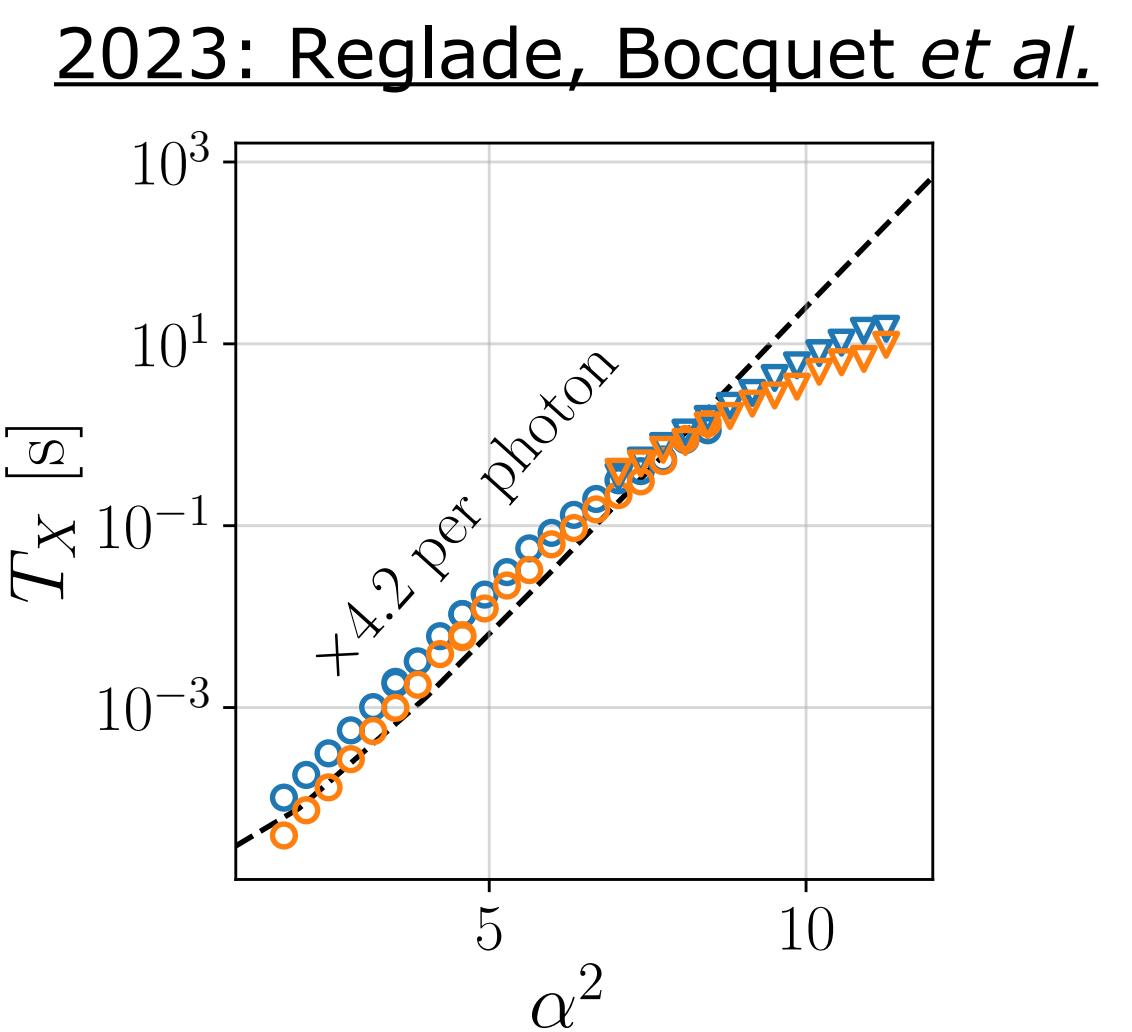
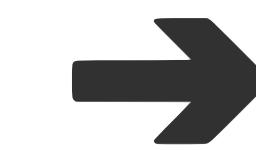
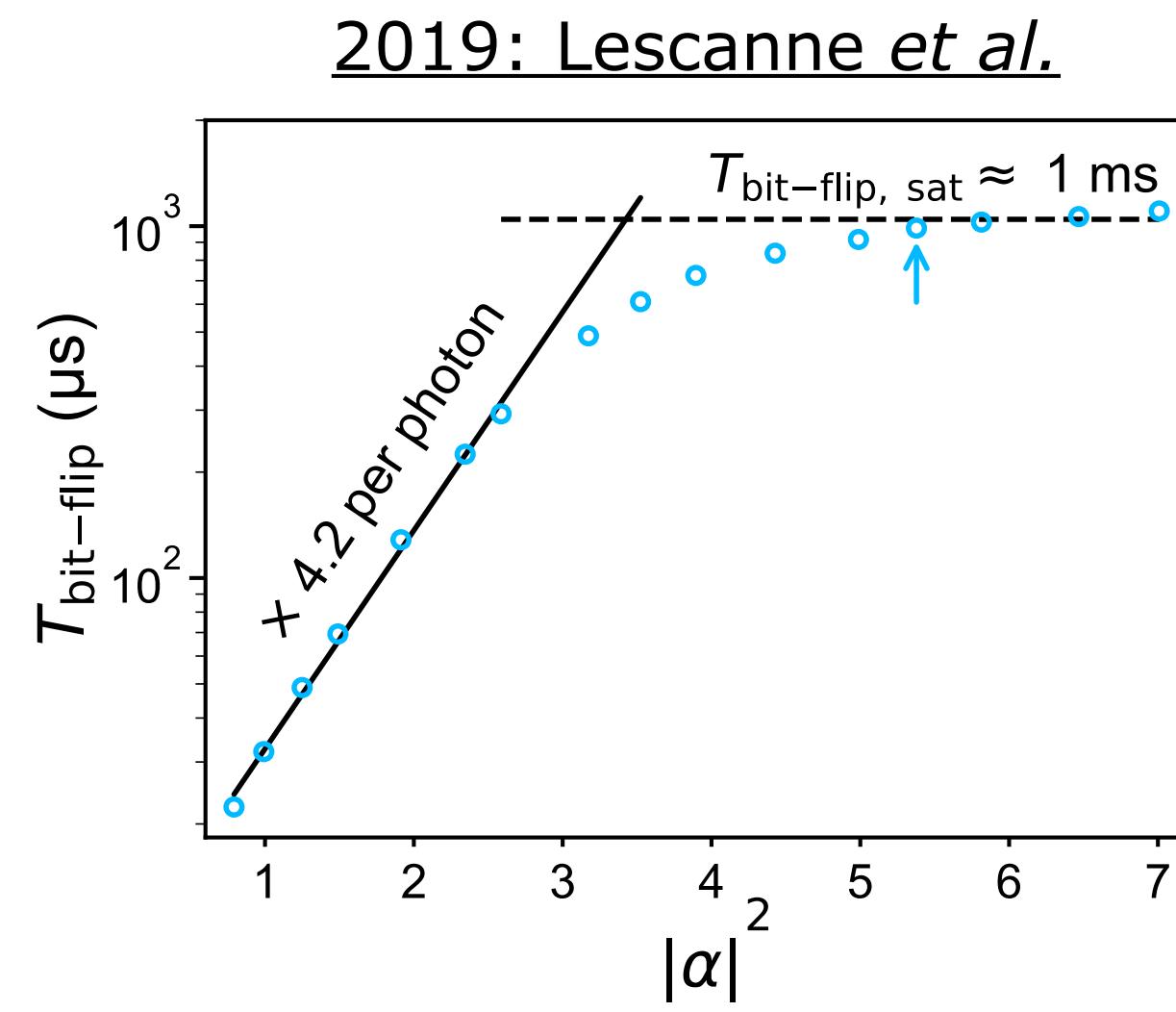
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Bit-flip lifetime of cat qubits

Dissipative cat qubits

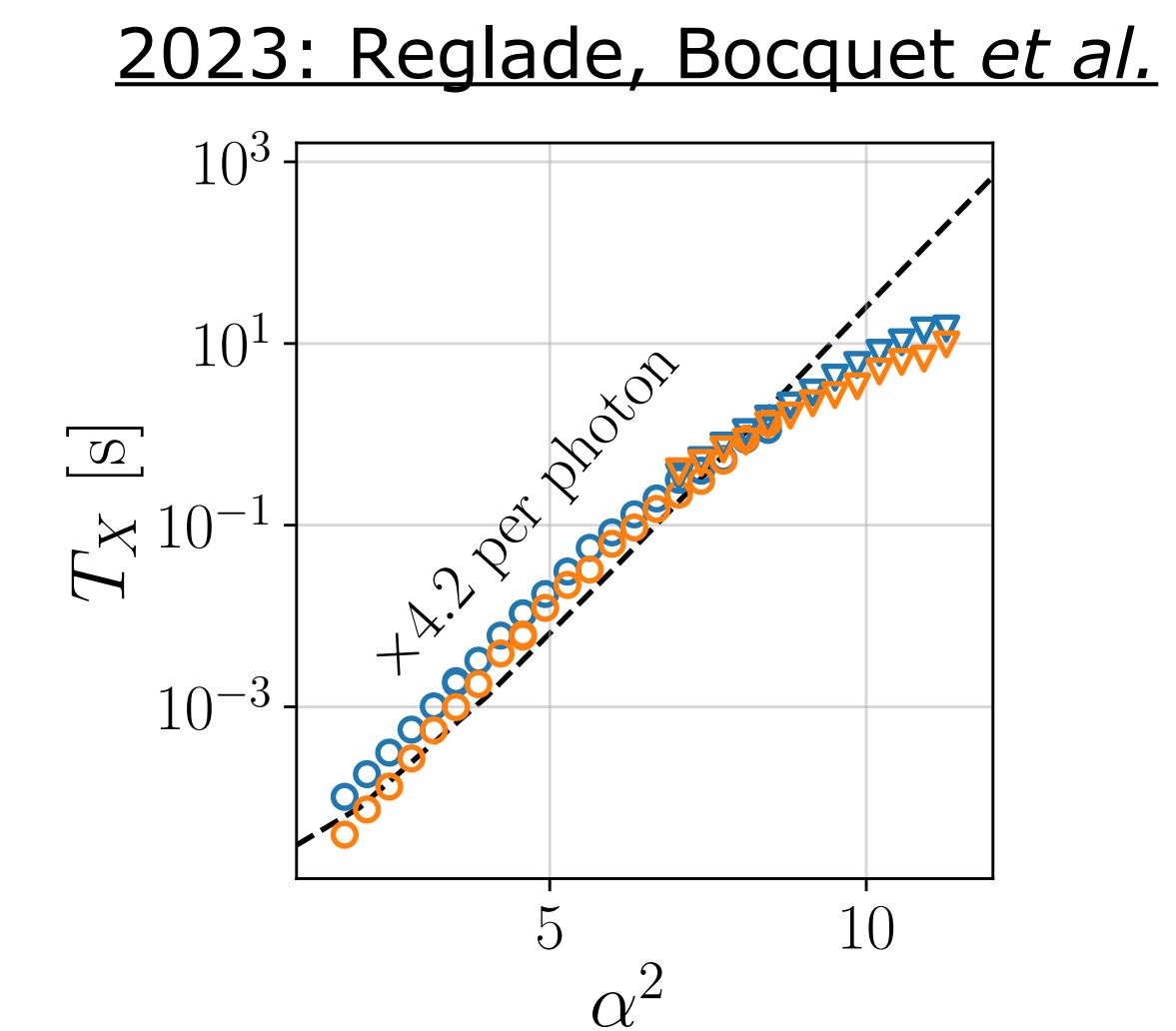
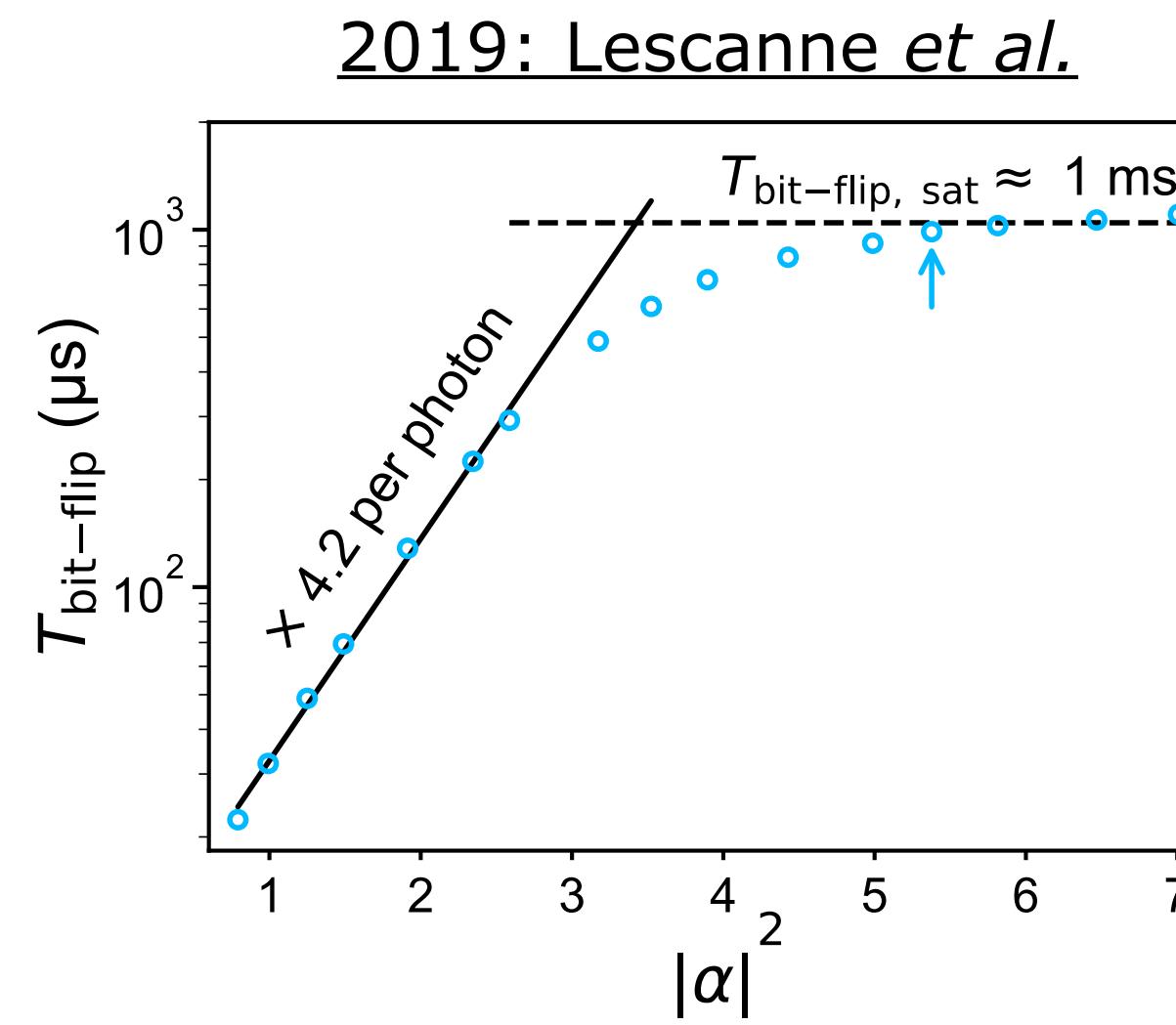
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- >10s in 2023
- Exponential → Saturation



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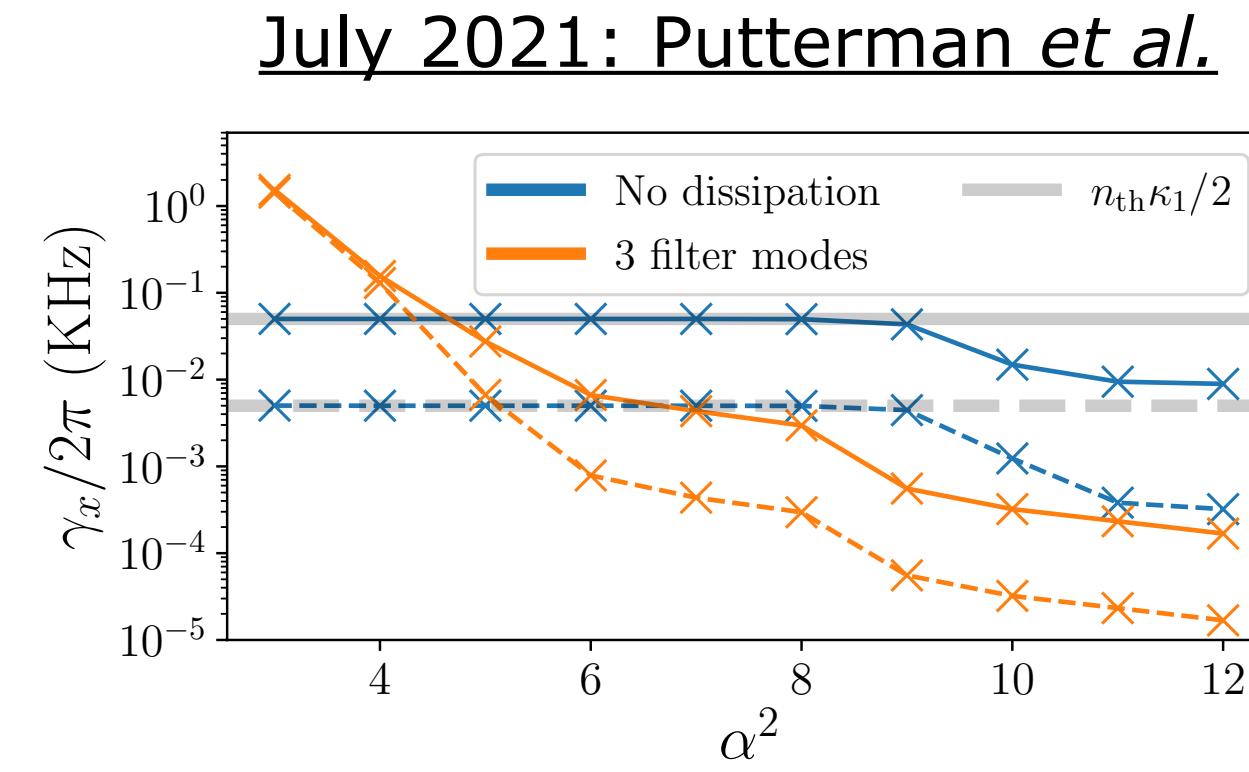
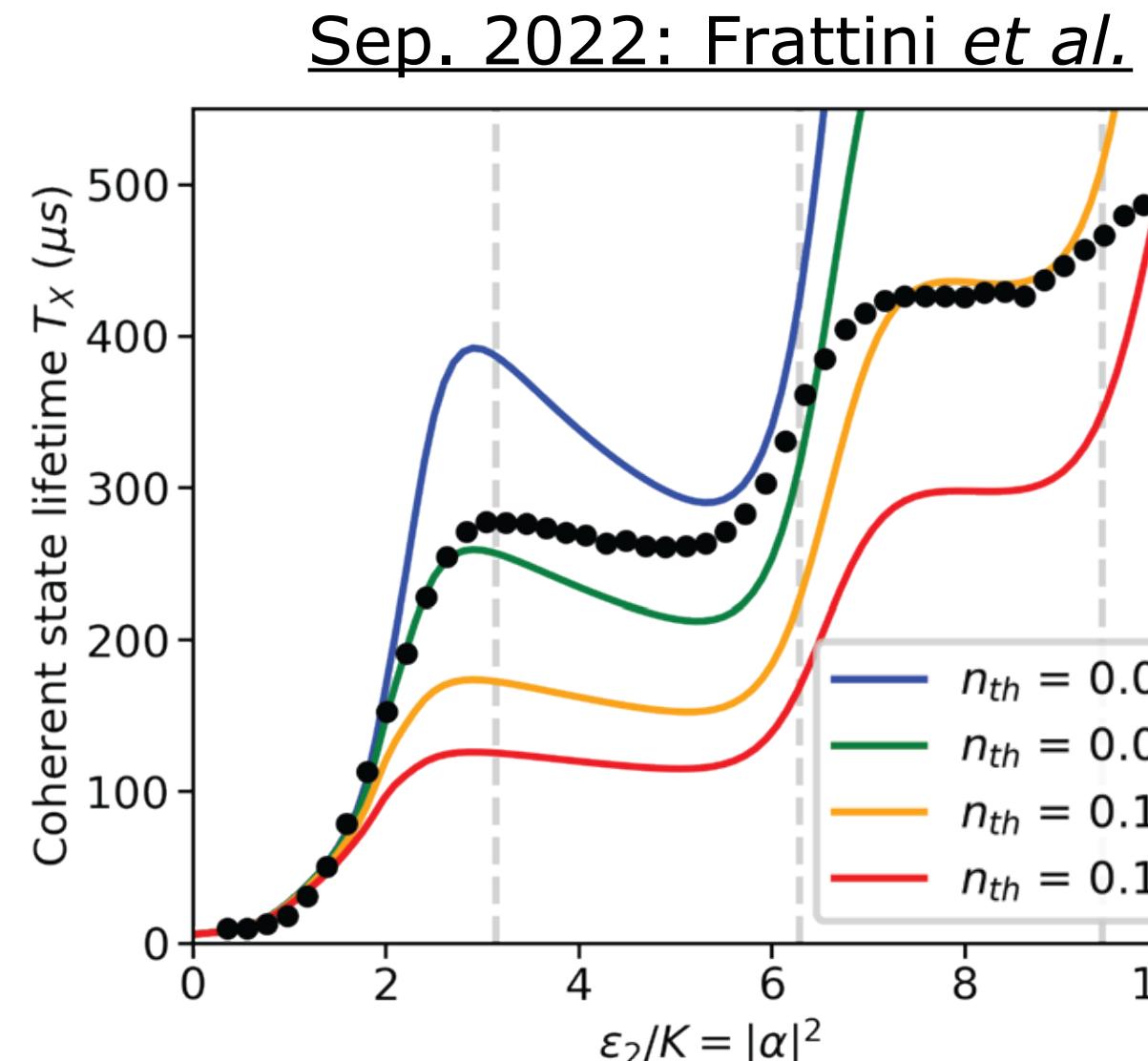
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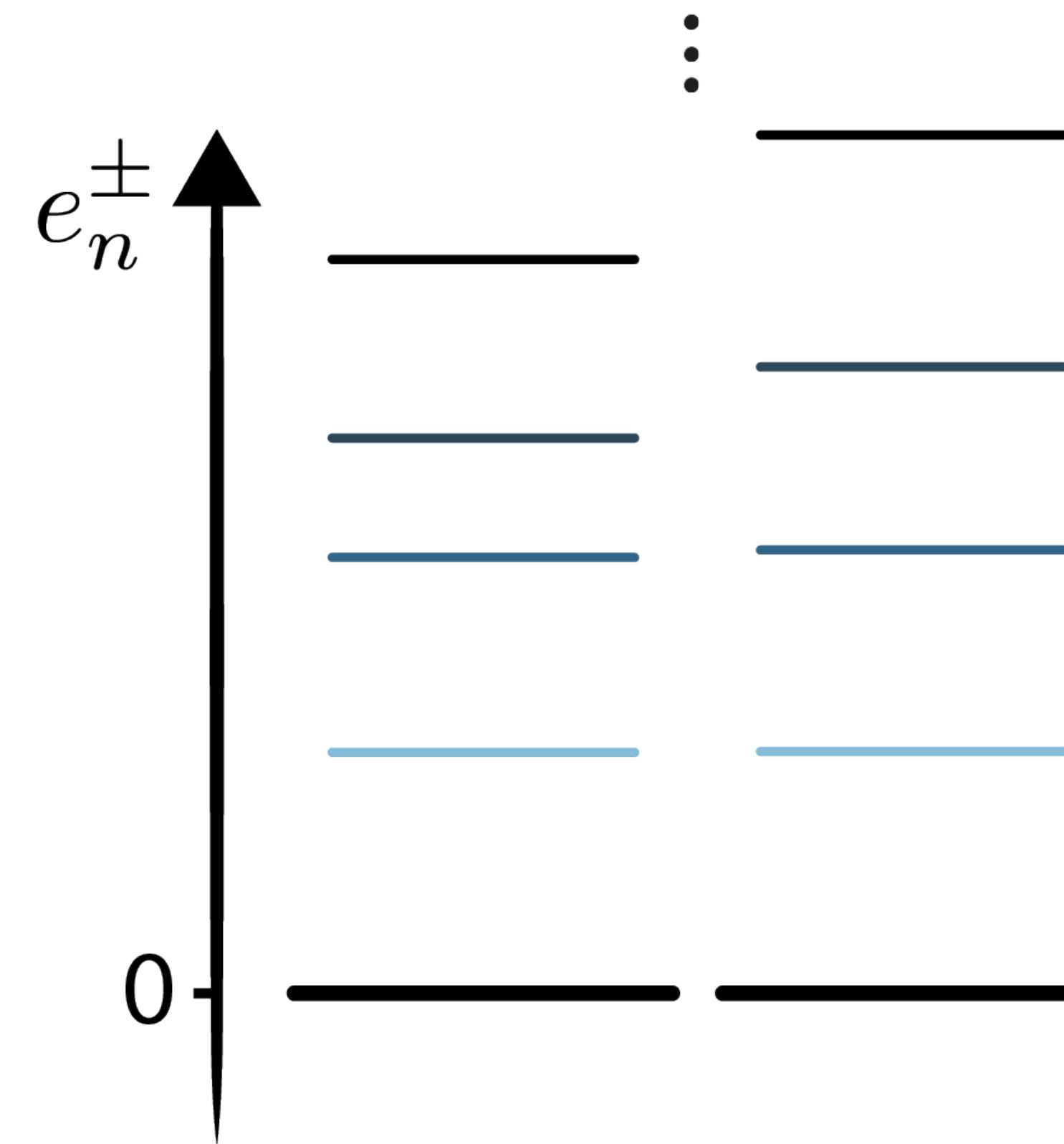
Kerr cat qubits

- ~500us in 2022
- Exponential → Plateaus
- Saturation predicted in 2021 by Puttermann et al.
- Why plateaus?



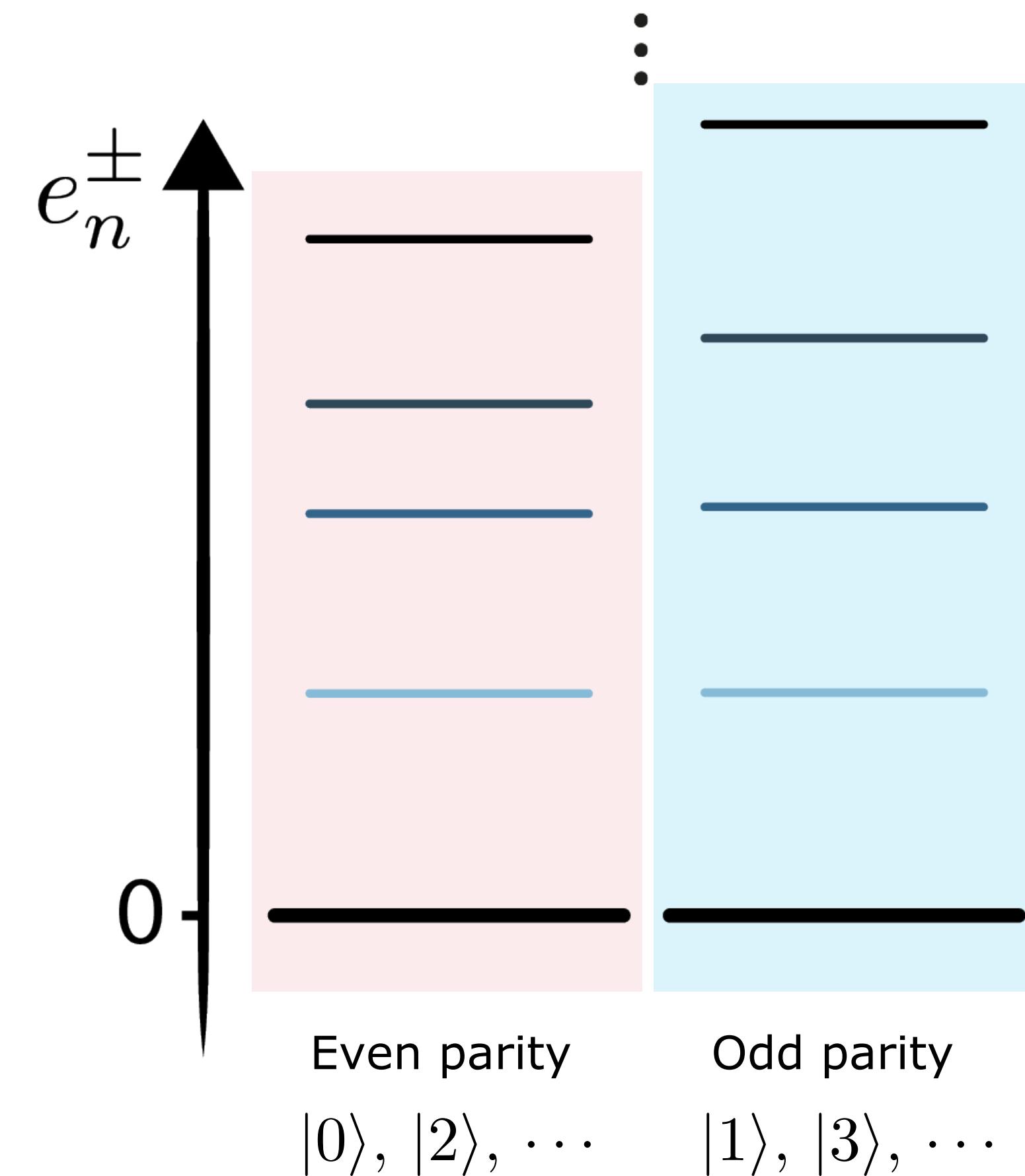
The spectrum of Kerr cat qubits

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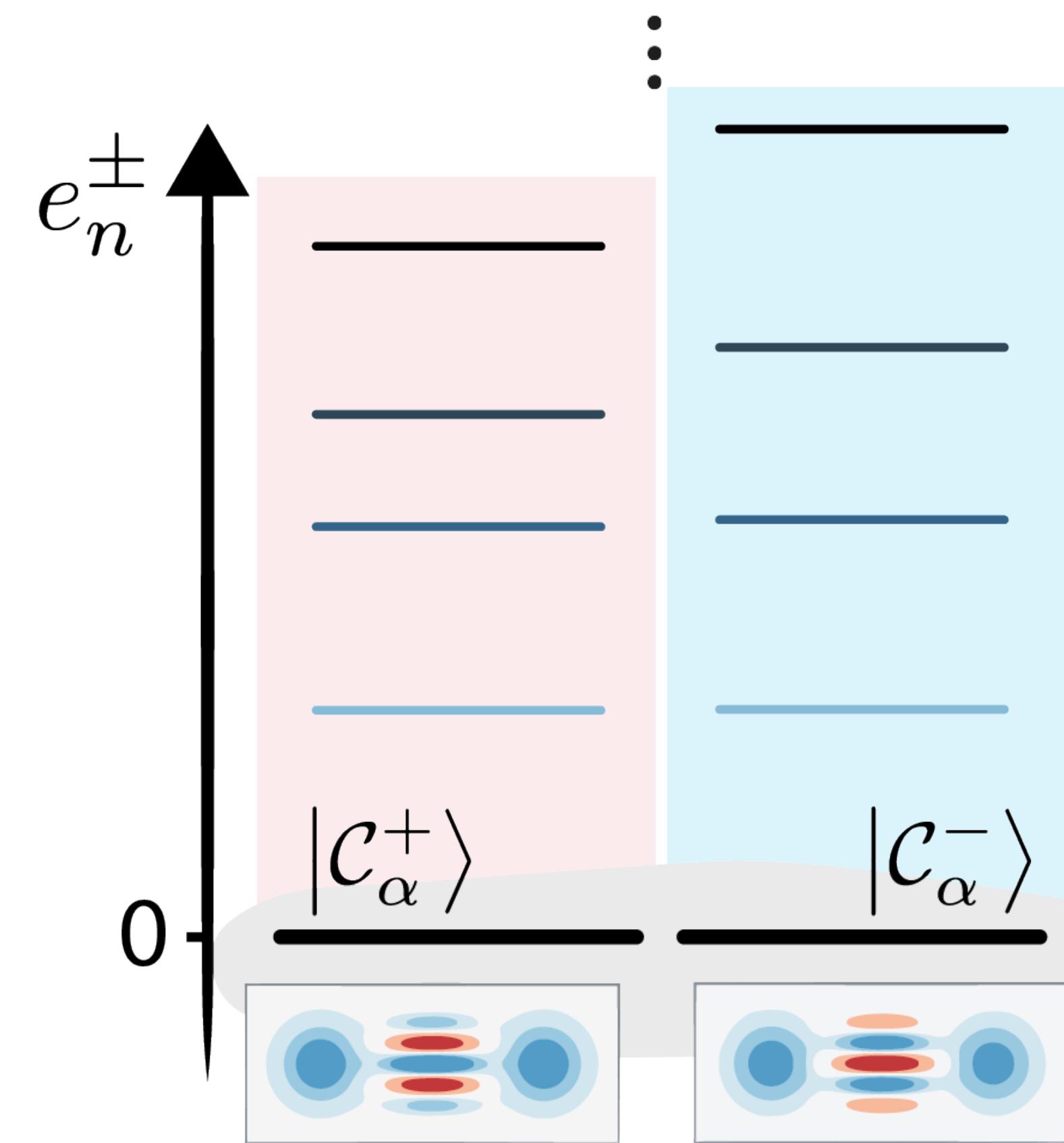
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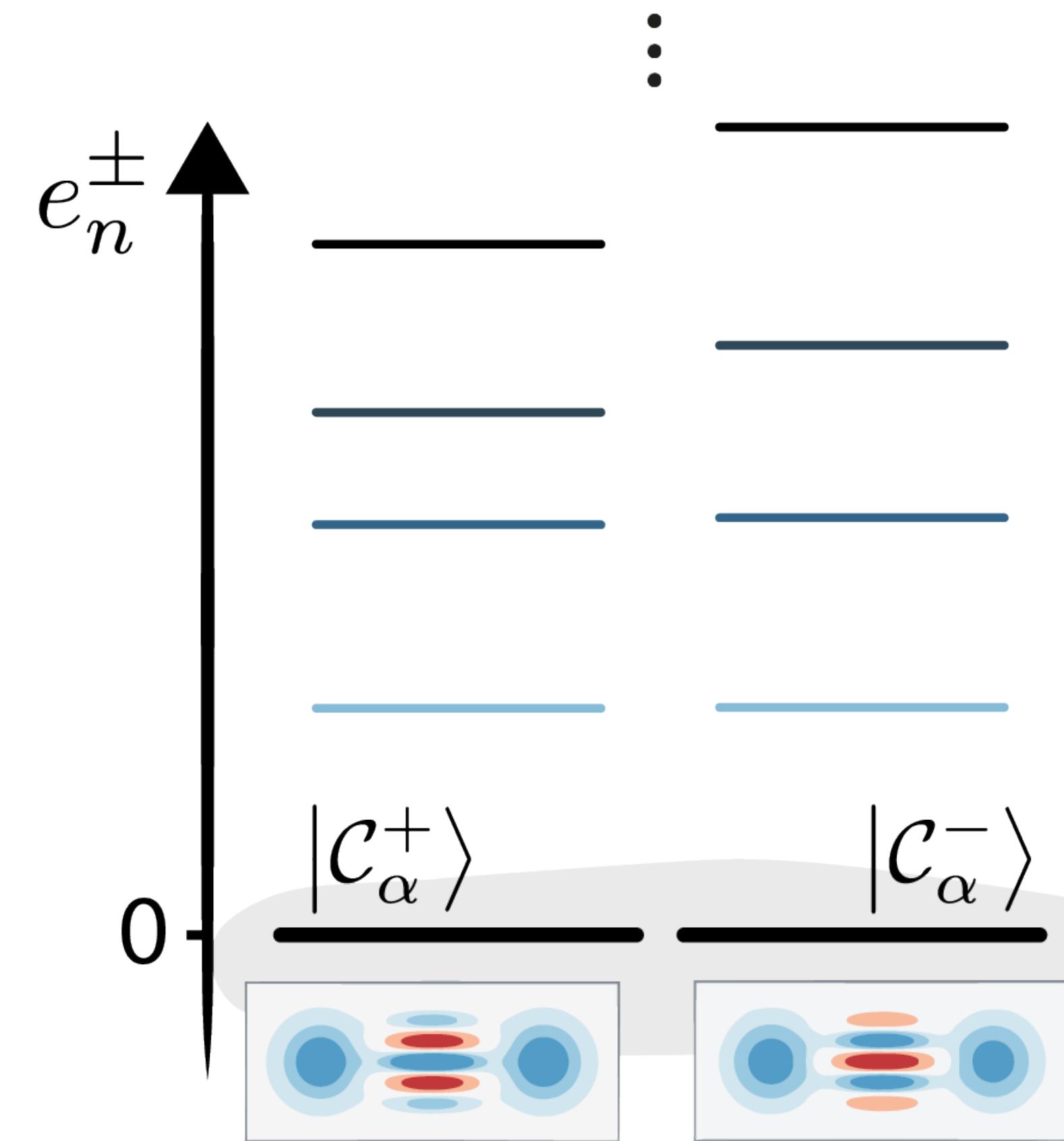
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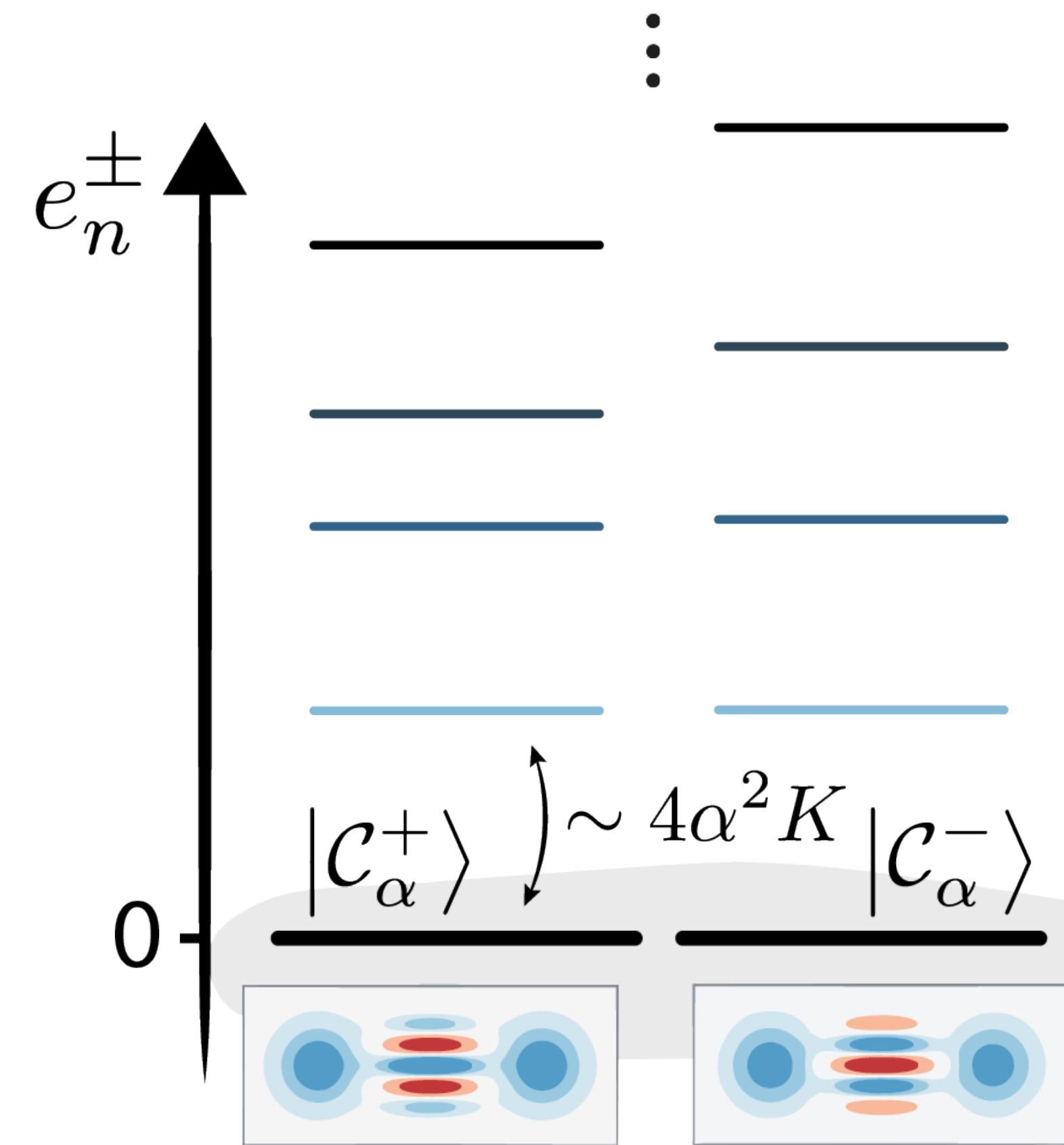
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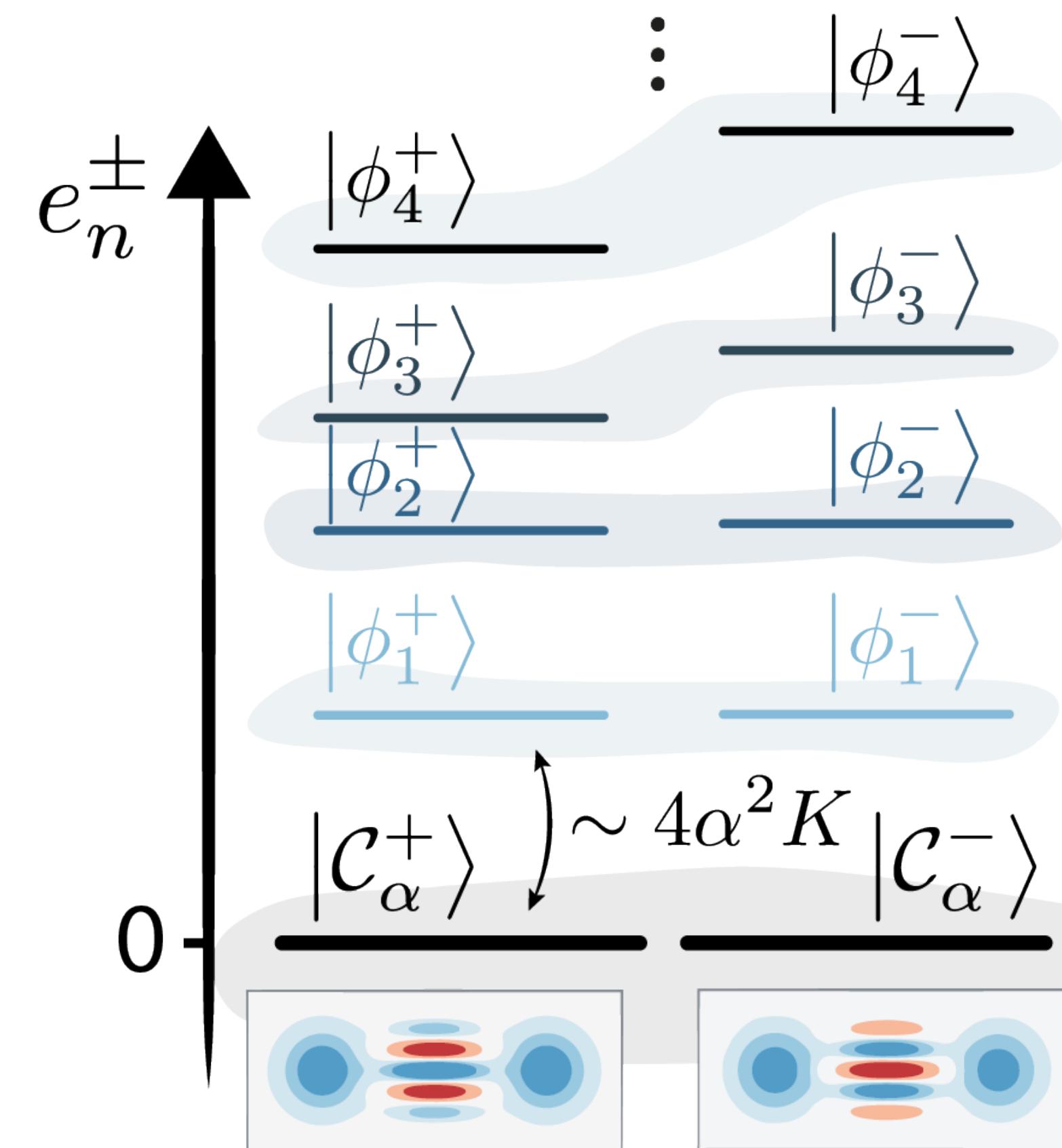
⌚ Gapped spectrum



The spectrum of Kerr cat qubits

$$H = -K(a^{\dagger 2} - \alpha^{*2})(a^2 - \alpha^2)$$

► Gapped spectrum



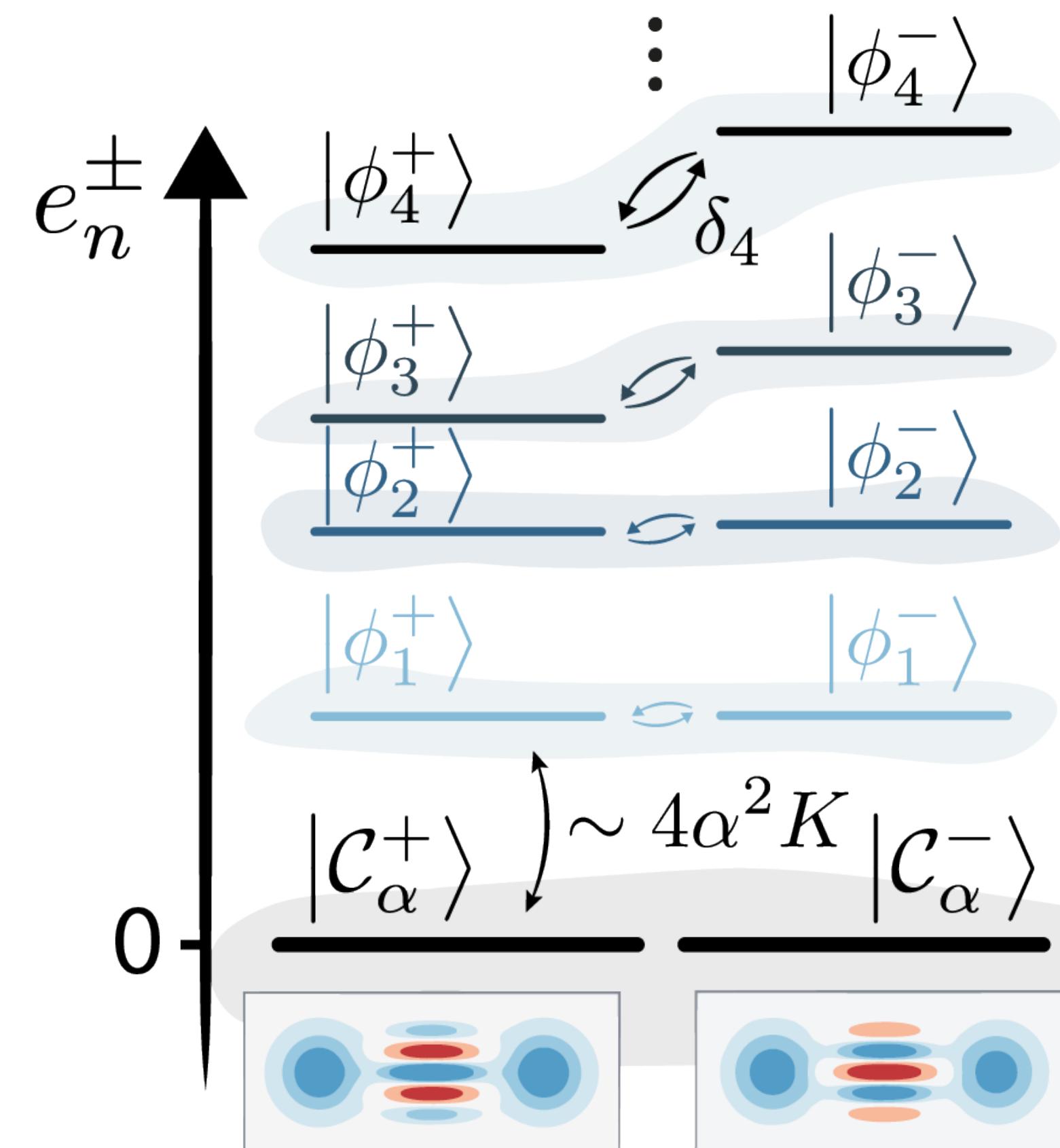
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- Gapped spectrum
- Standard Tunneling Model

phase basis

$$\begin{pmatrix} e_n + \delta_n/2 & 0 \\ 0 & e_n - \delta_n/2 \end{pmatrix}$$

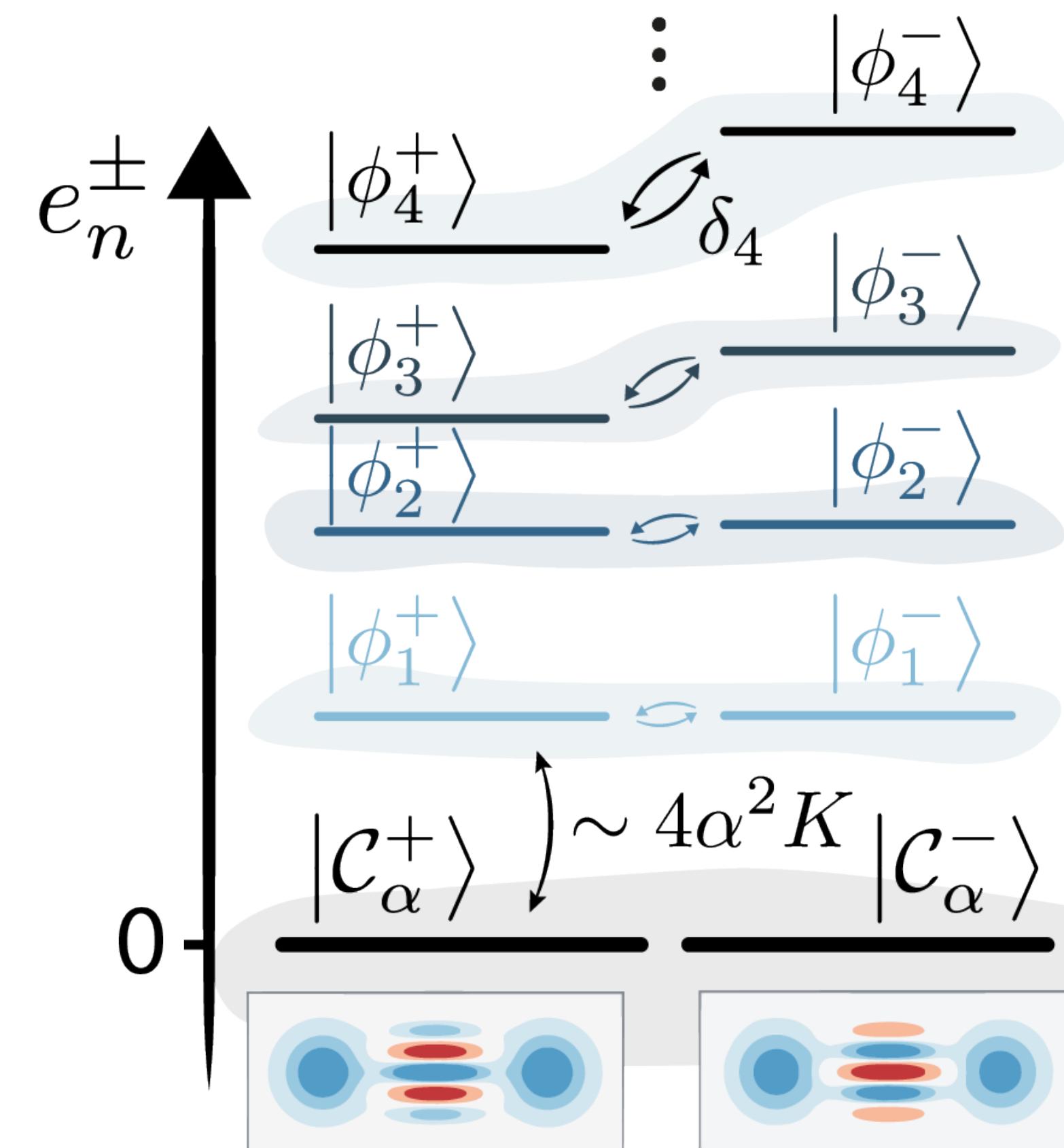


The spectrum of Kerr cat qubits

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$$\begin{array}{c} \text{bit basis} \\ \left(\begin{array}{cc} e_n & \delta_n/2 \\ \delta_n/2 & e_n \end{array} \right) \end{array} \xleftarrow{} \begin{array}{c} \text{phase basis} \\ \left(\begin{array}{cc} e_n + \delta_n/2 & 0 \\ 0 & e_n - \delta_n/2 \end{array} \right) \end{array}$$

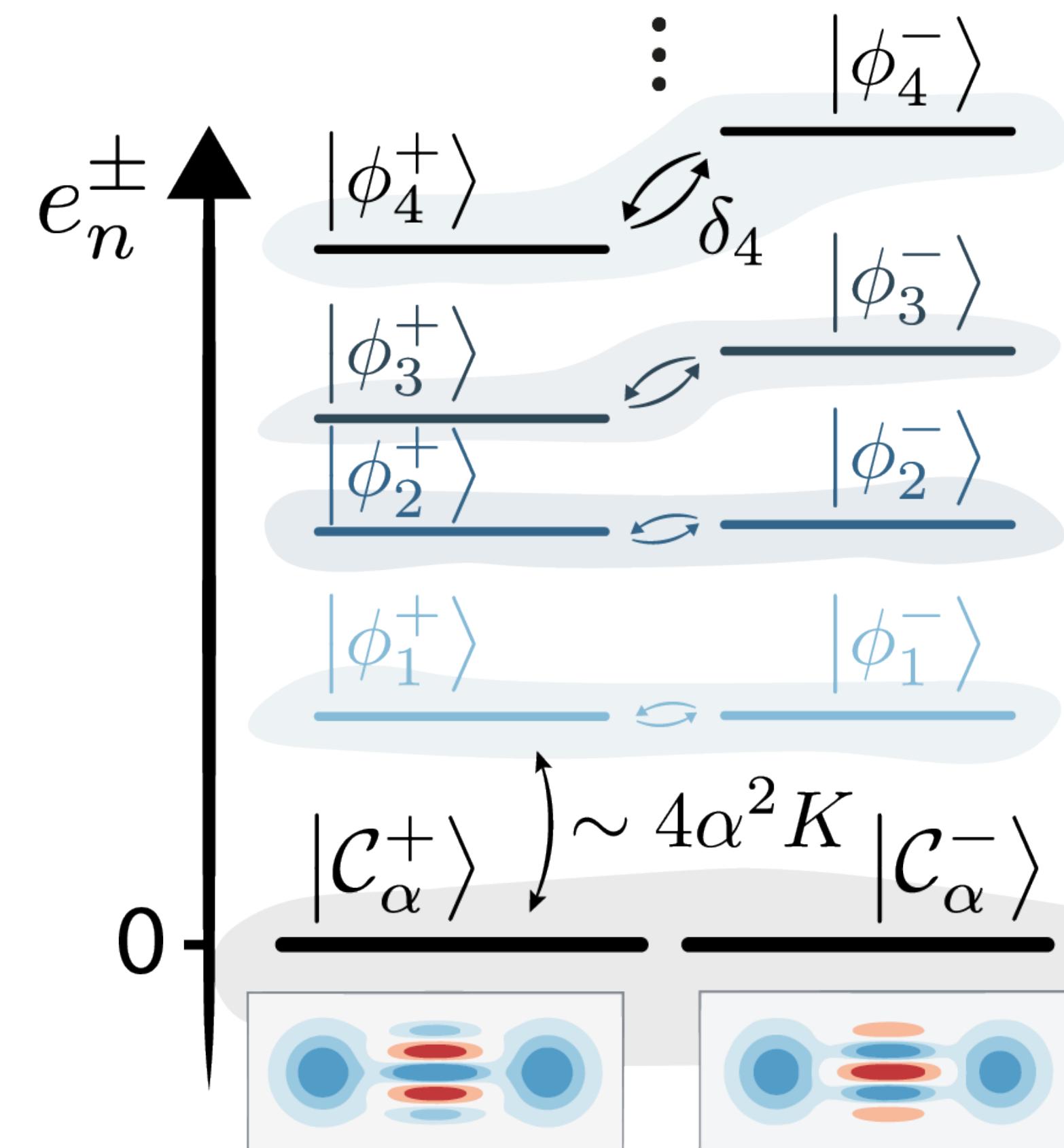


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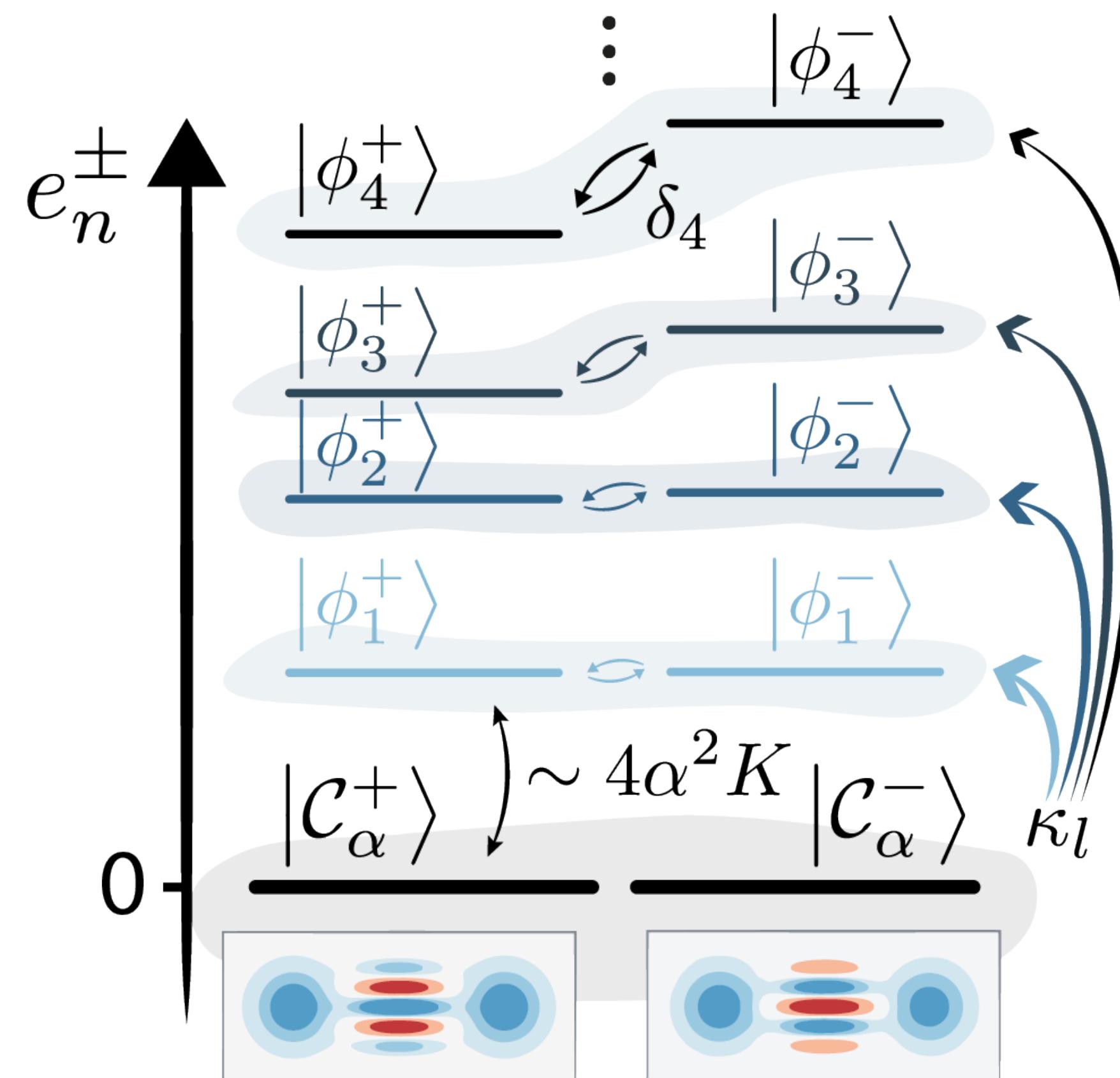
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- Incoherent leakage
 - ➔ Thermal photons
 - ➔ Pure dephasing



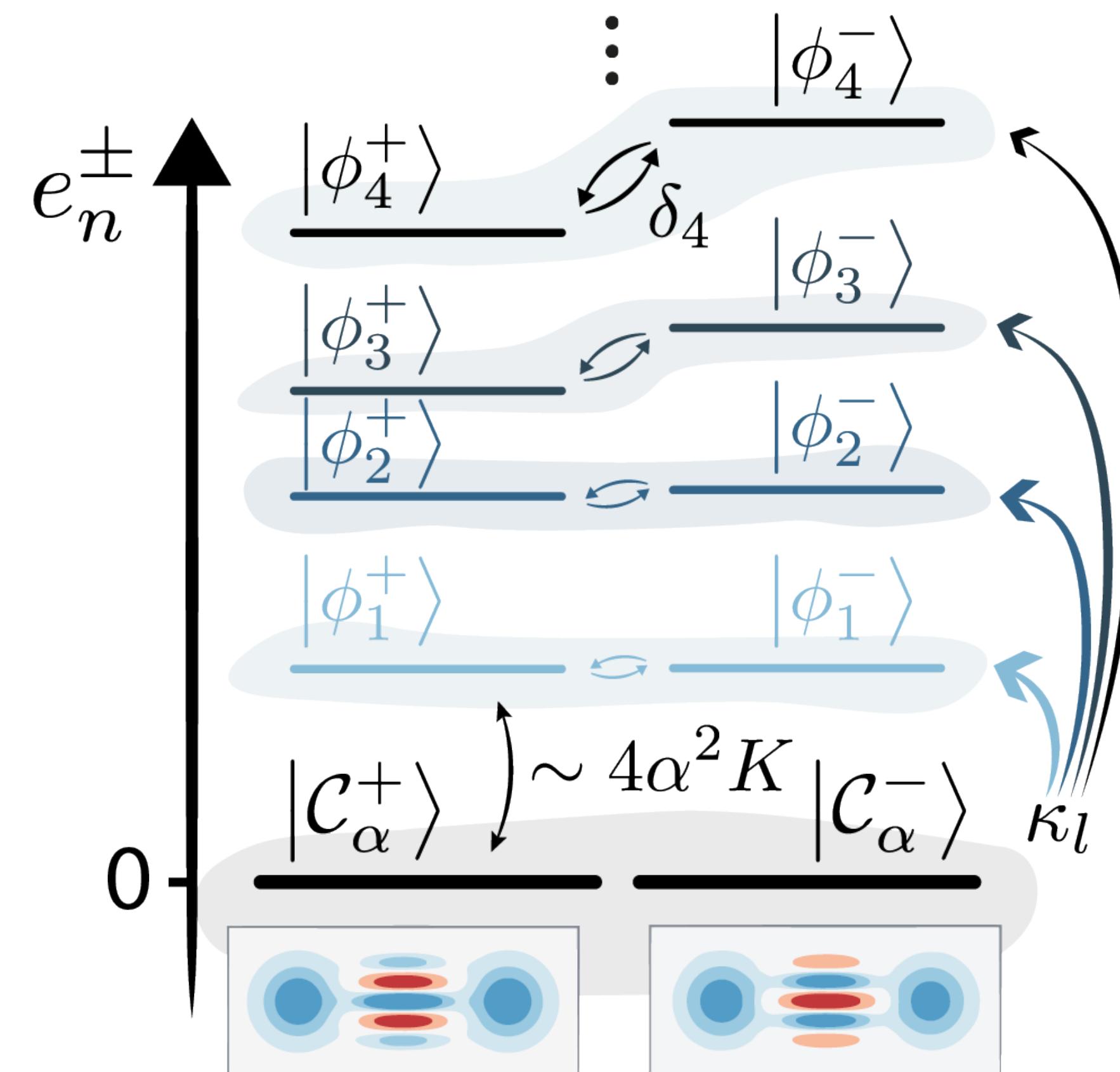
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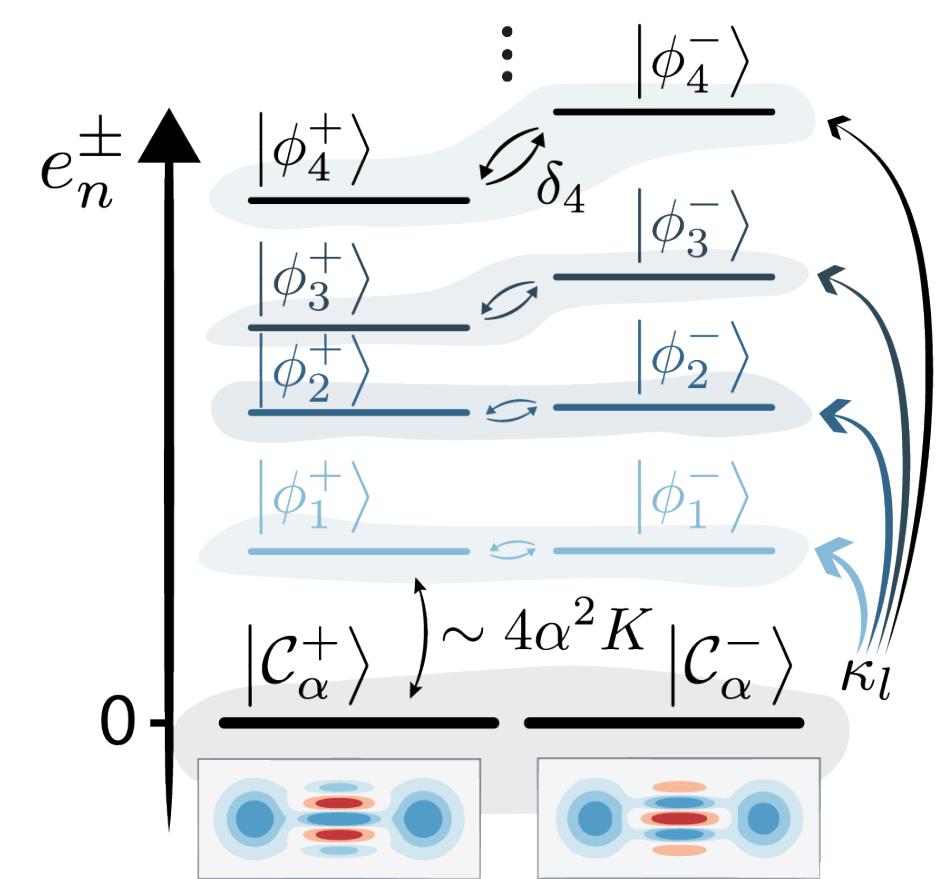
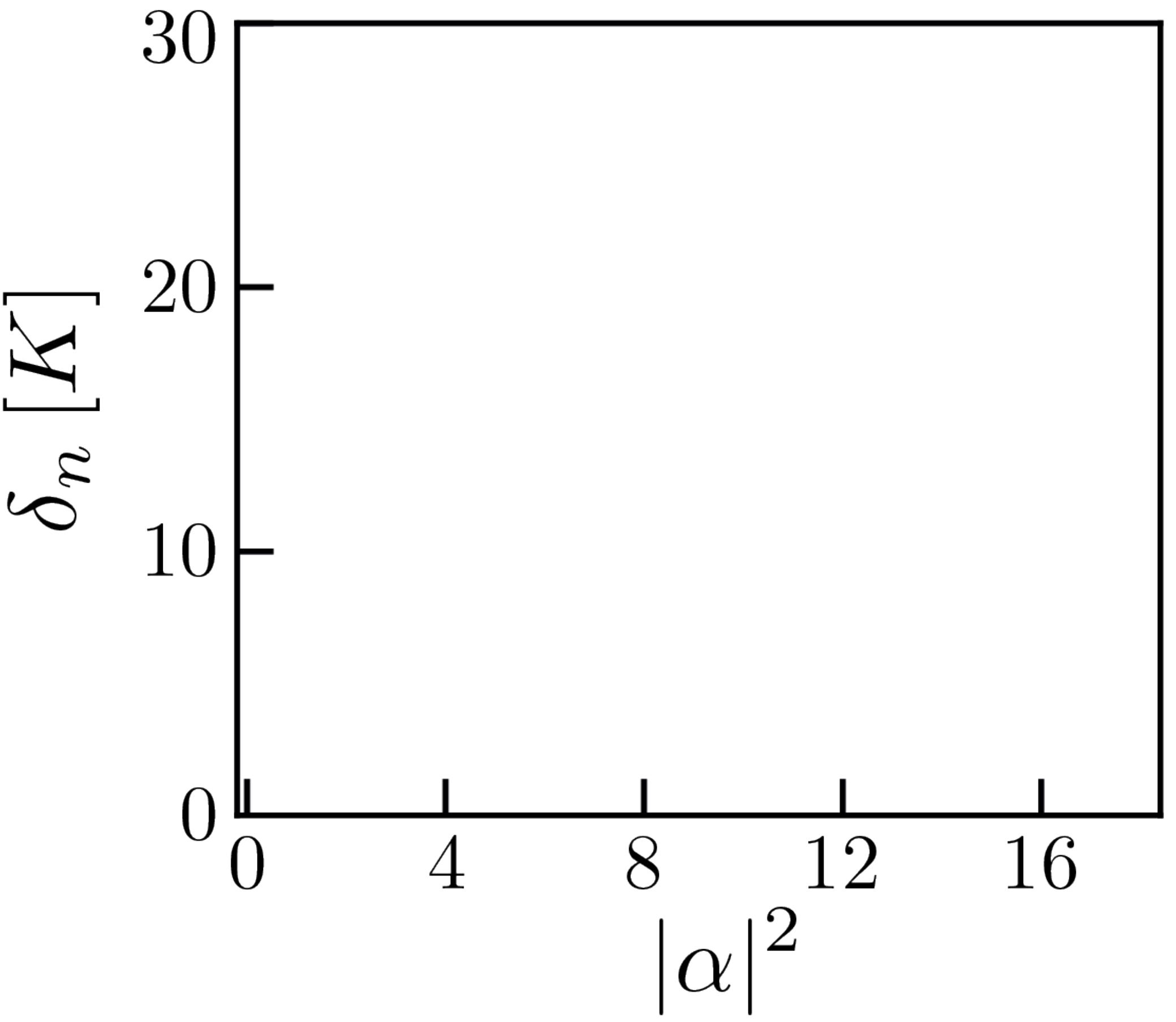
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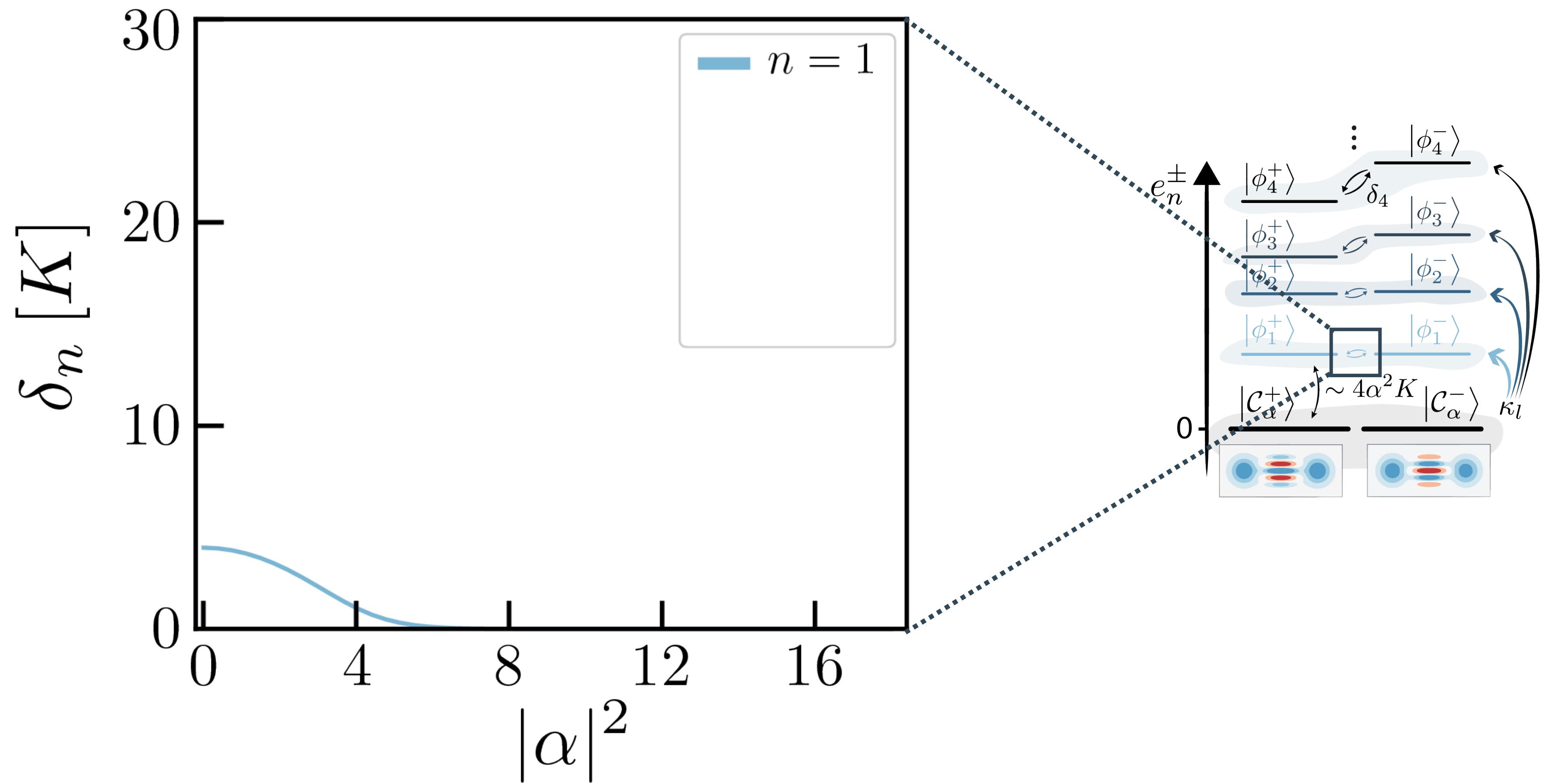


Bit-flip induced by incoherent leakage

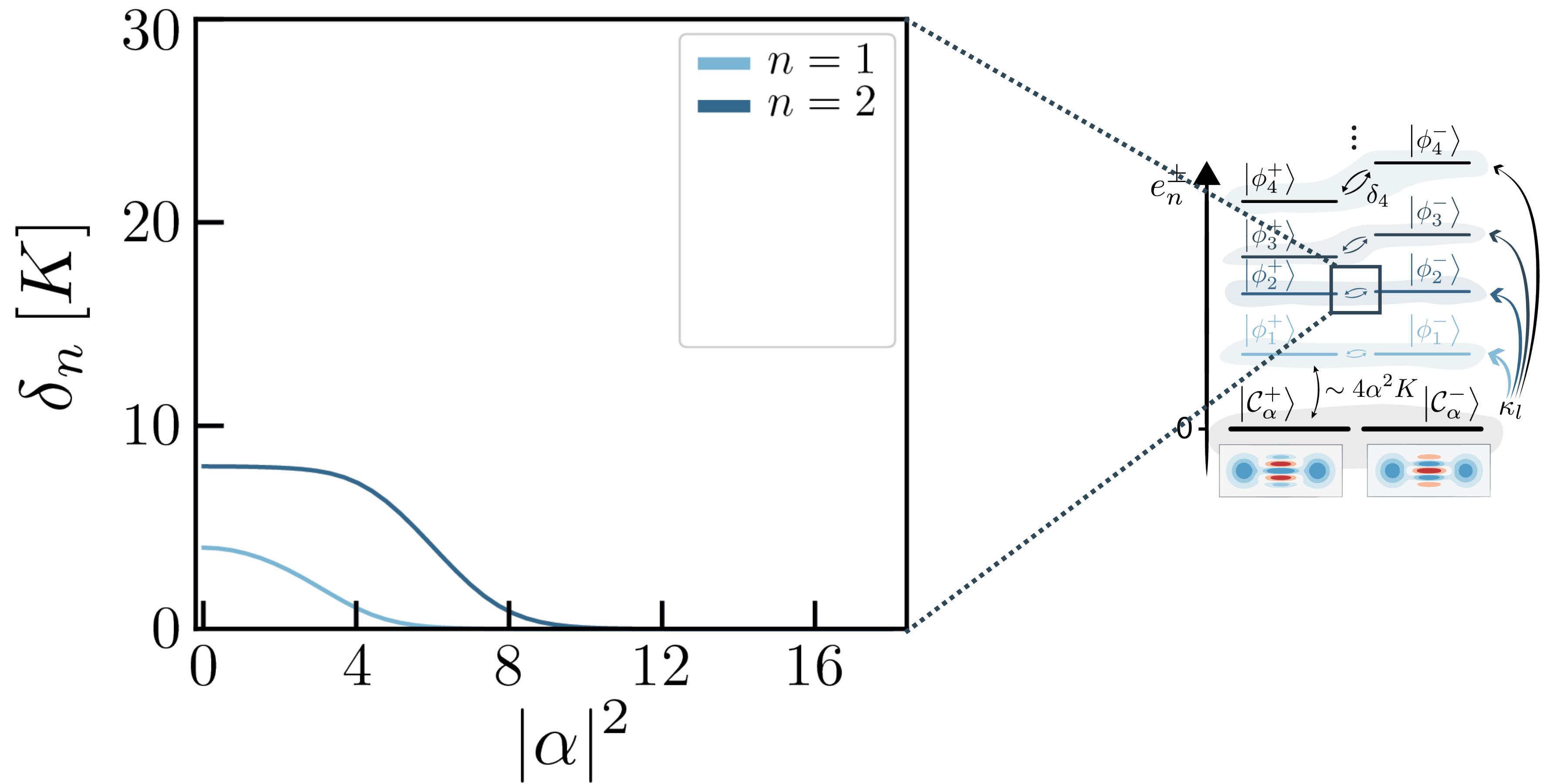
Bit-flip plateaus



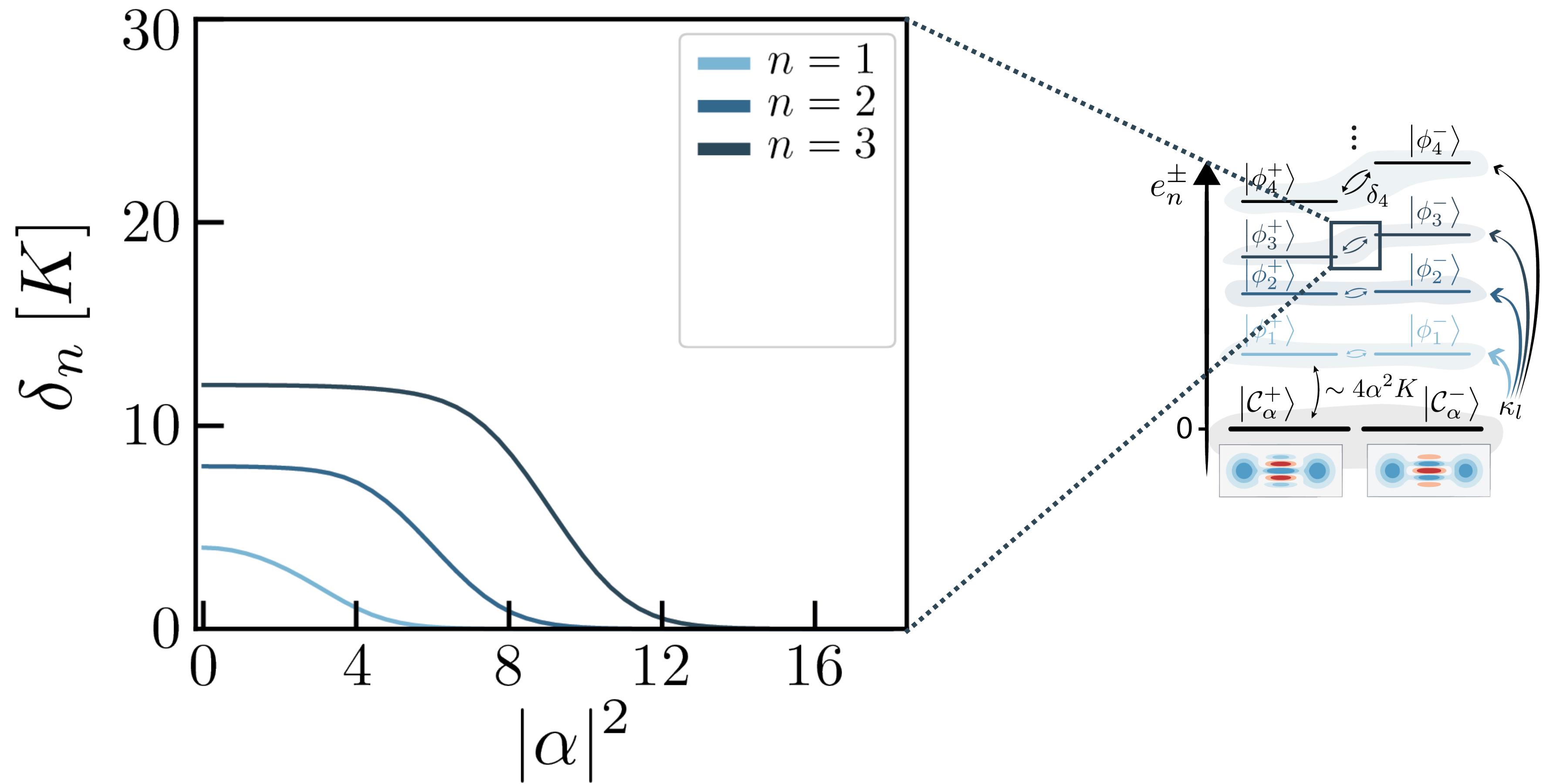
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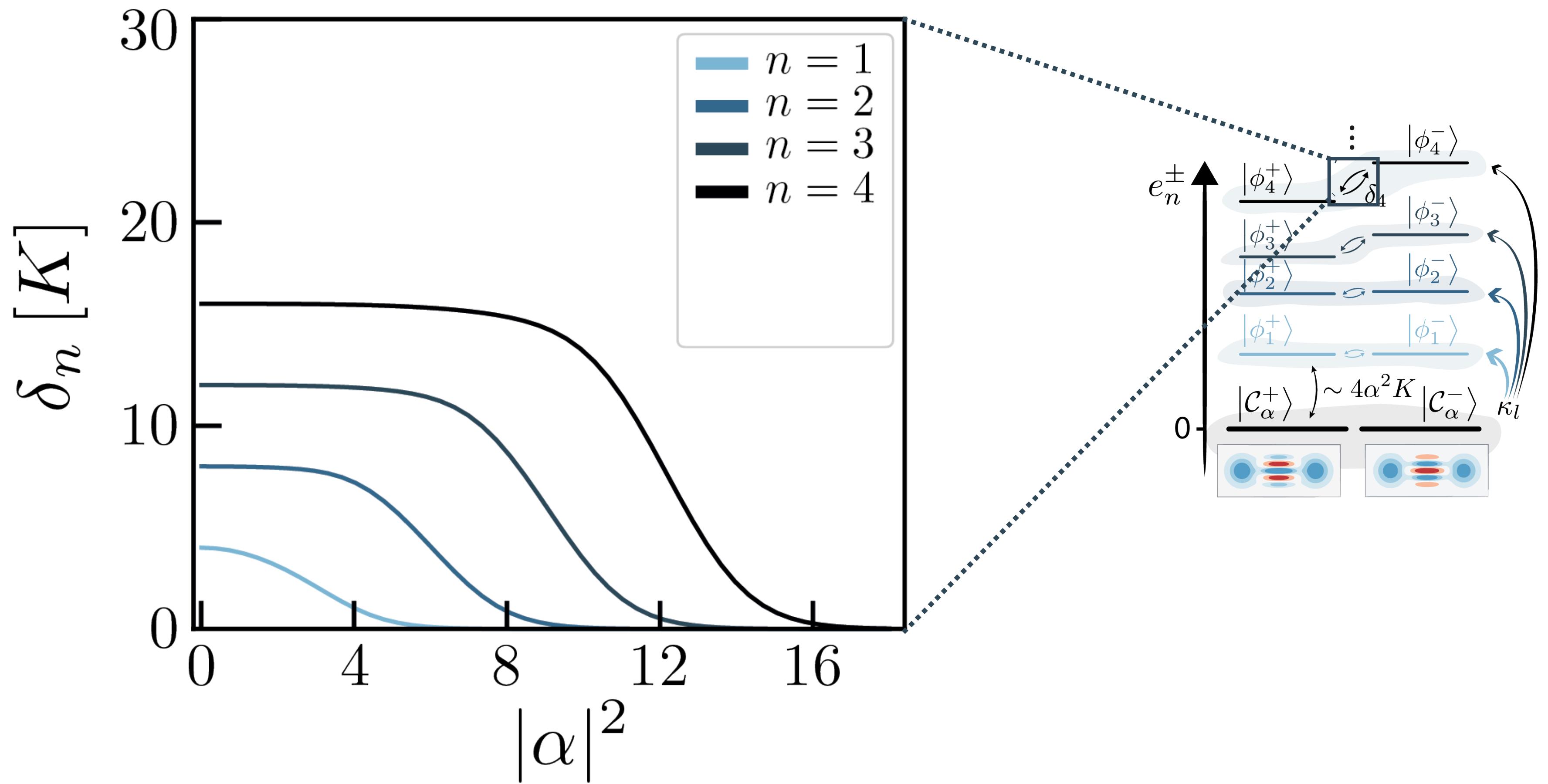
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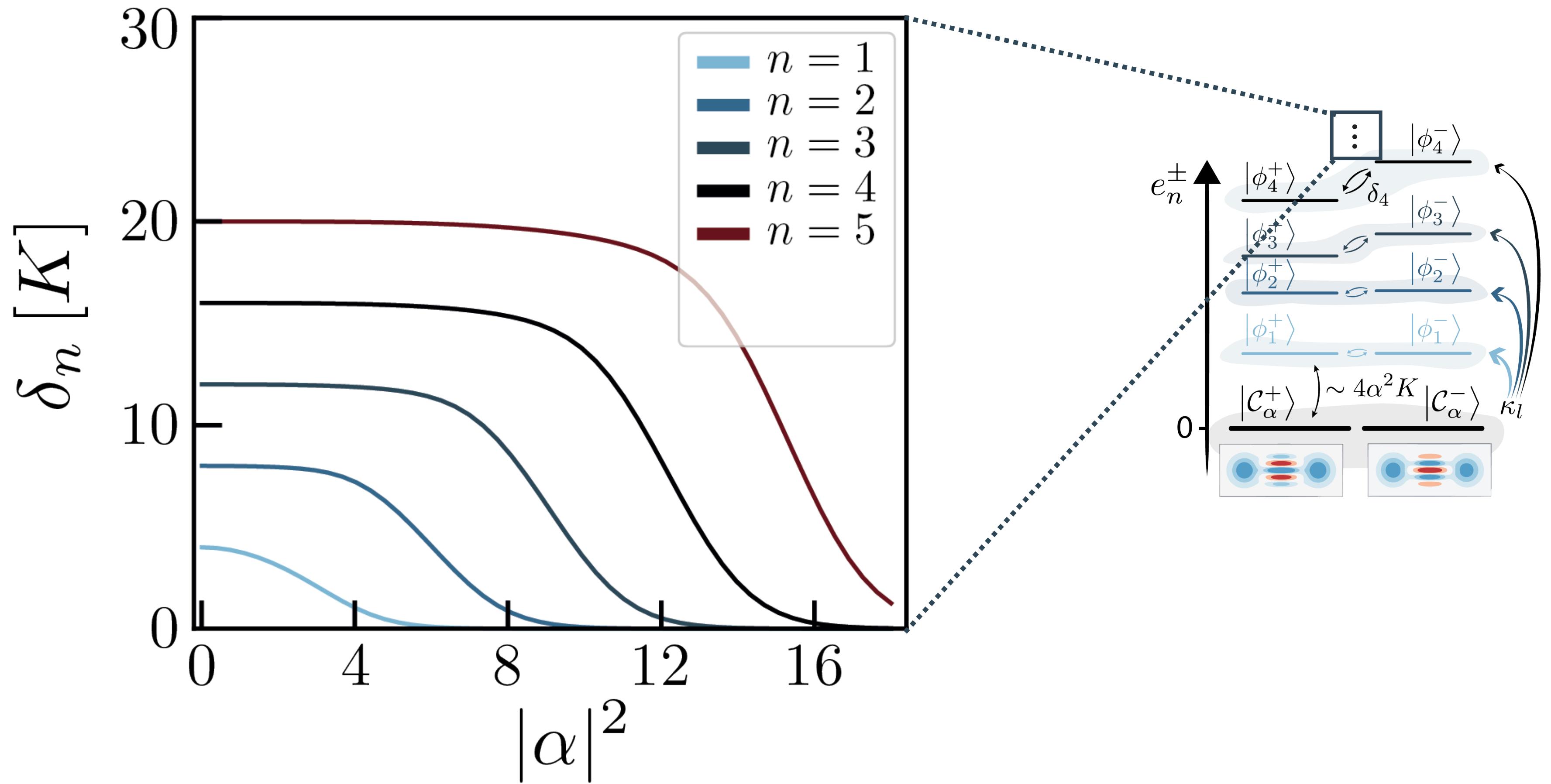
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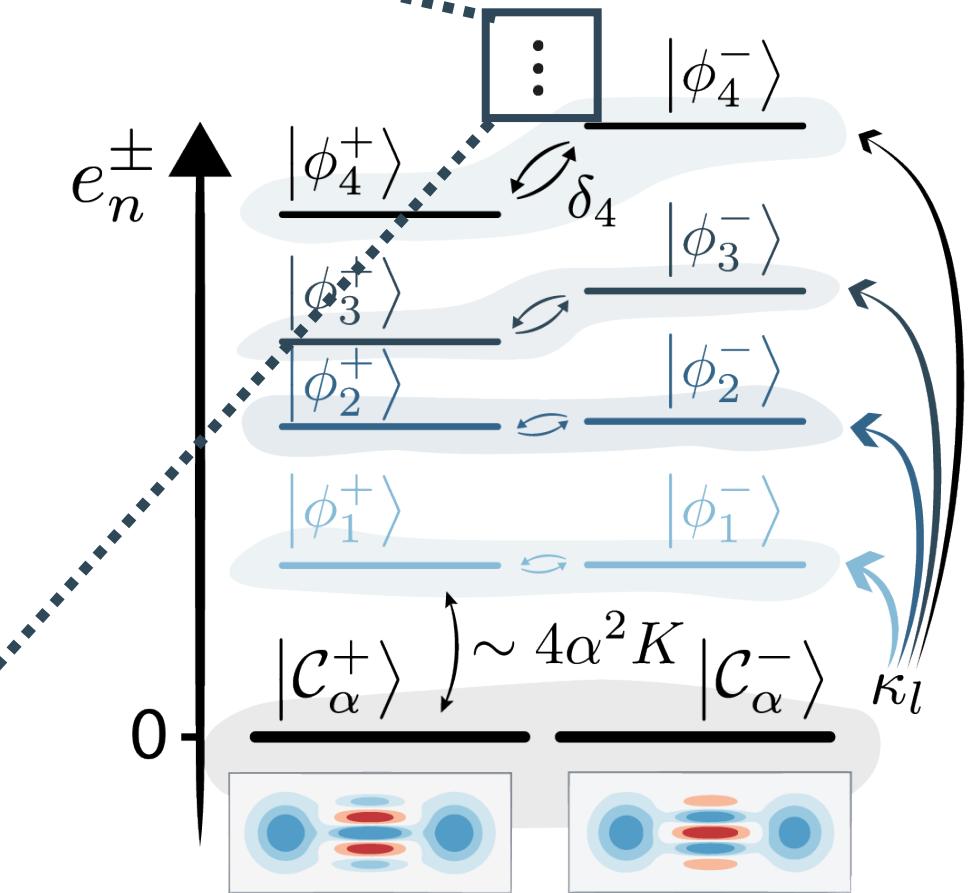
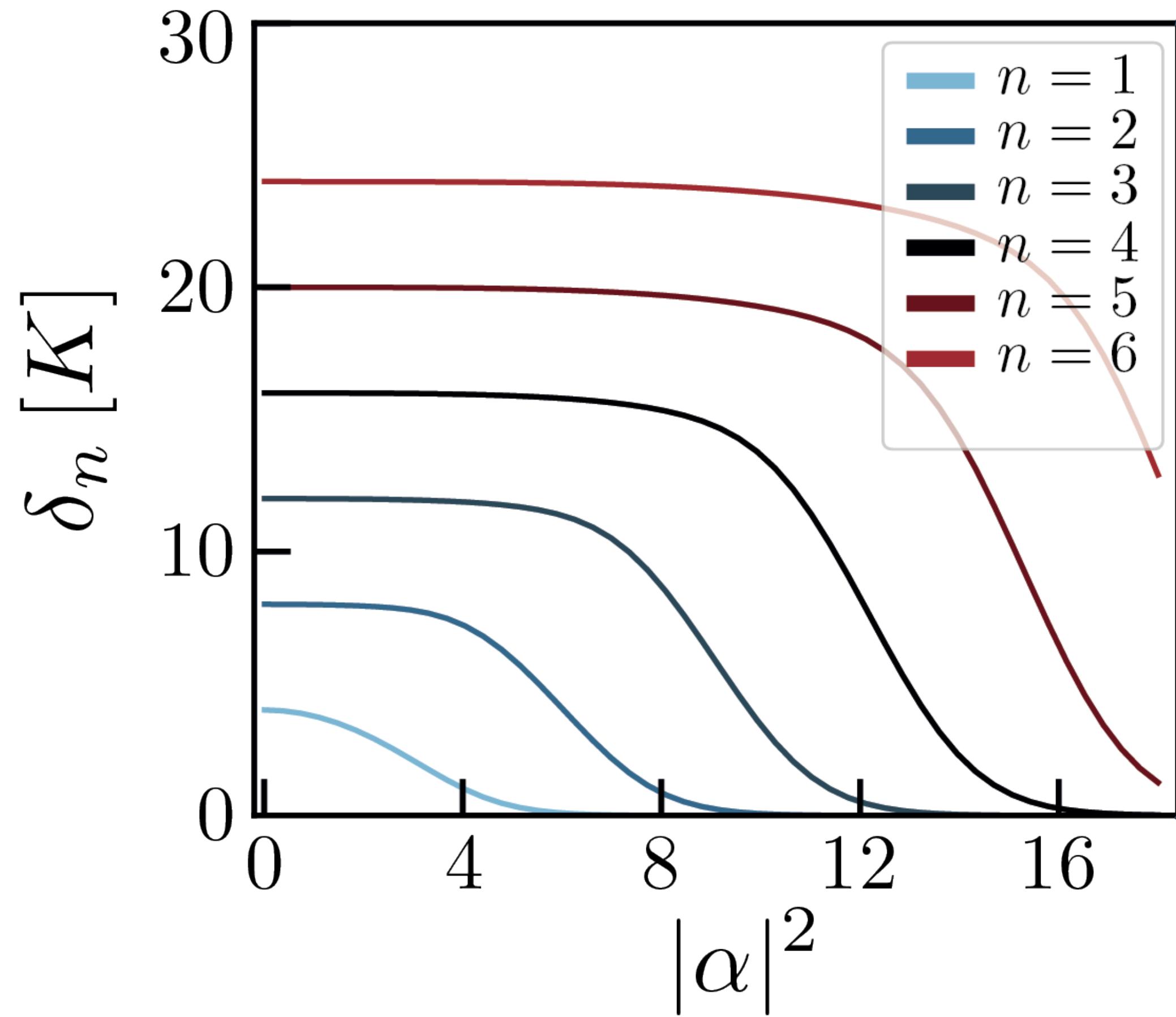
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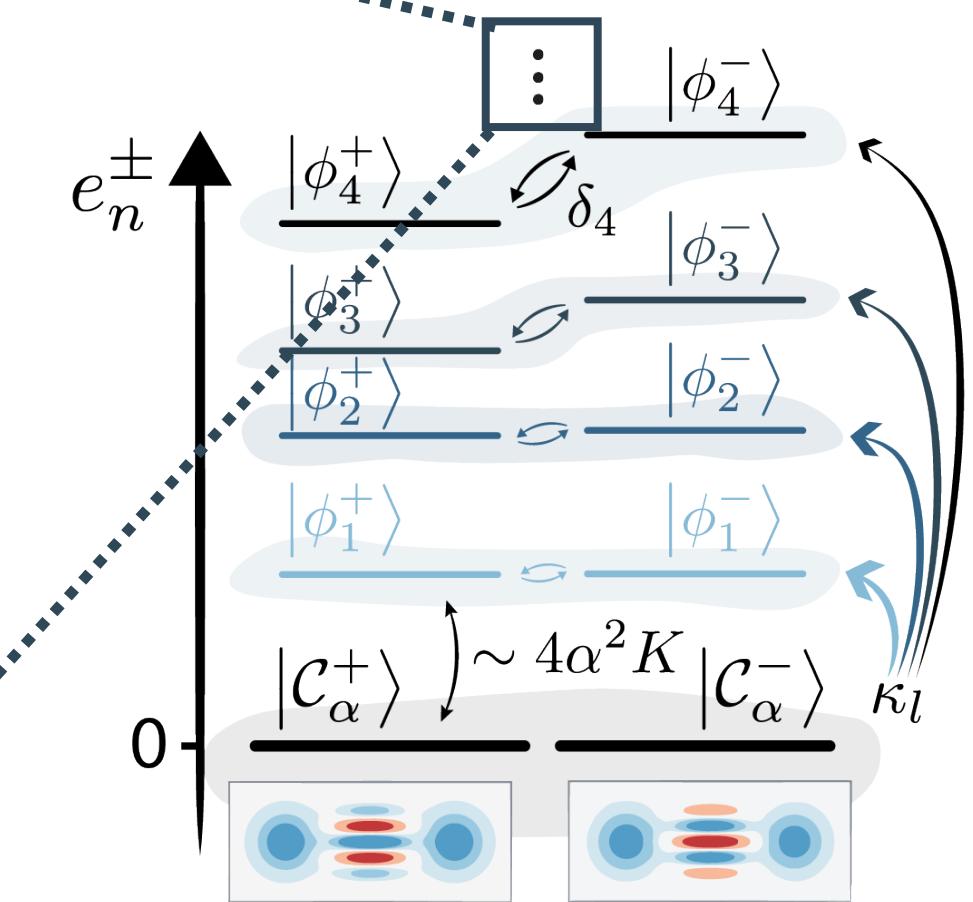
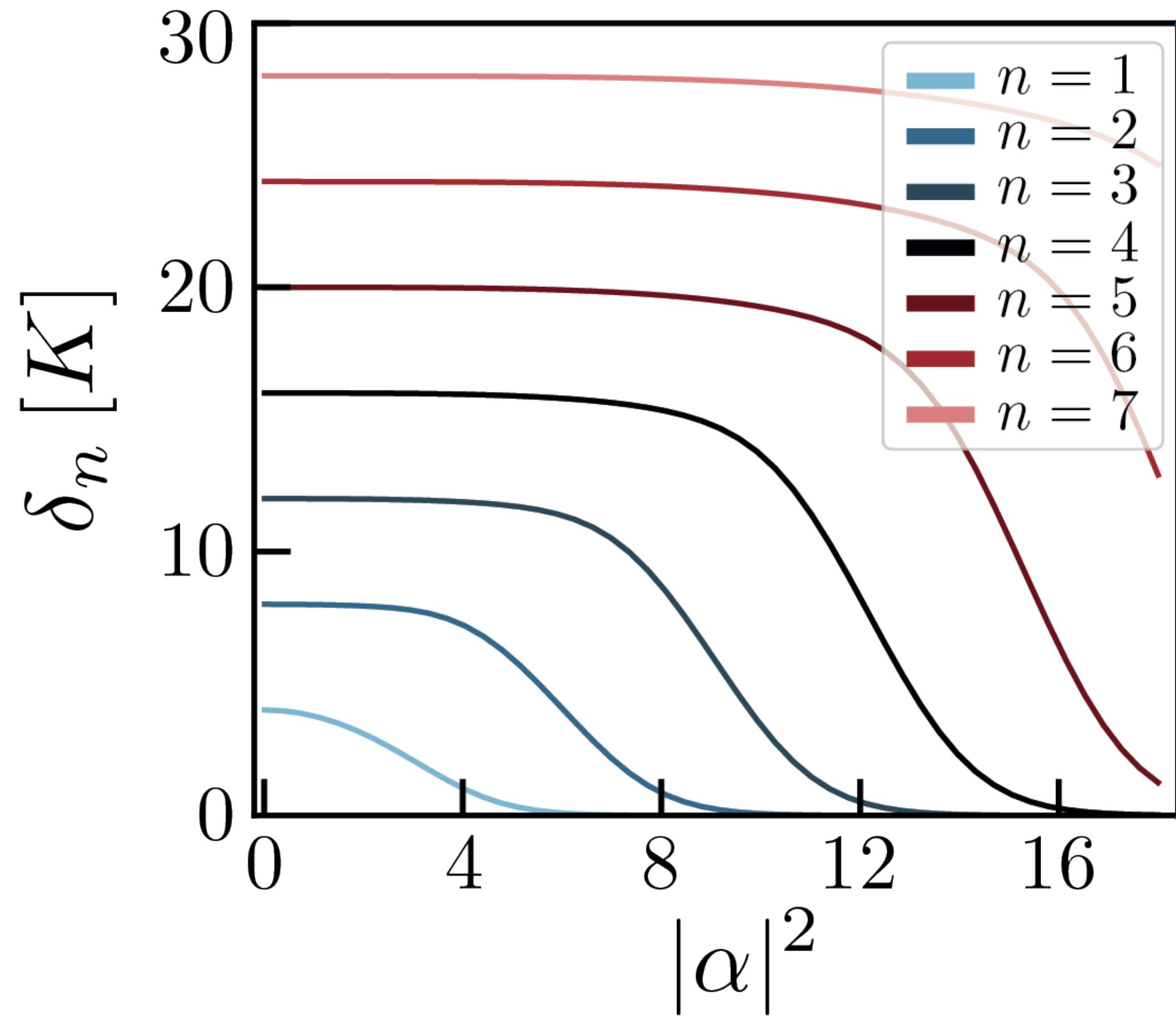
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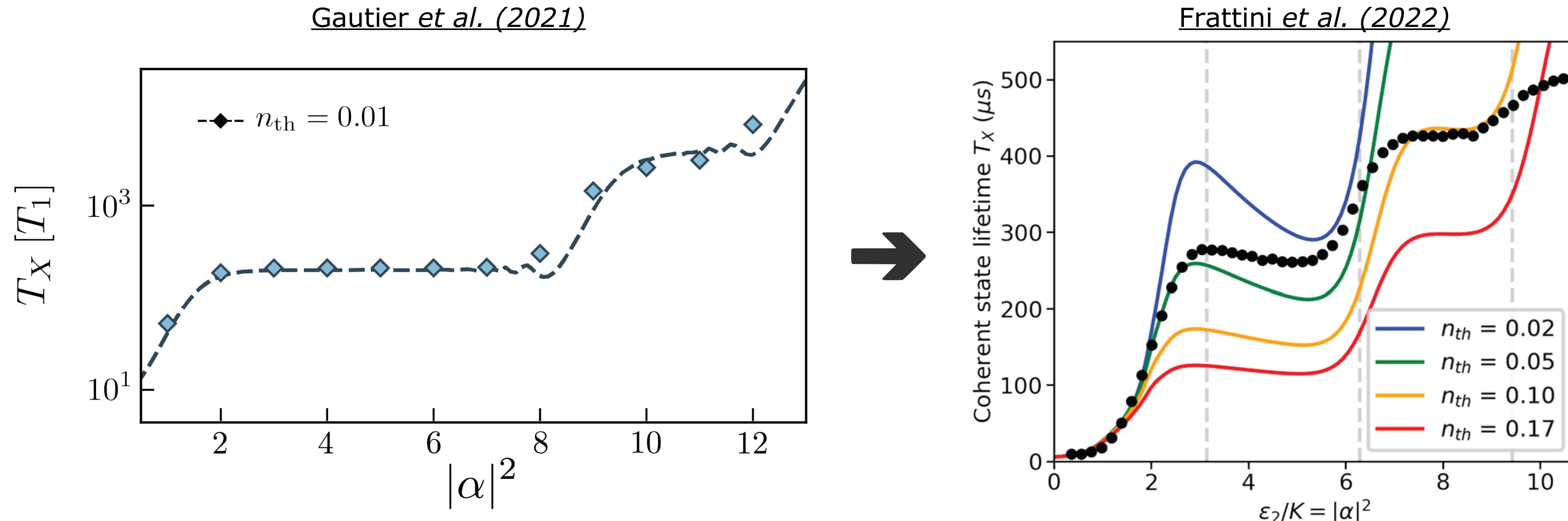
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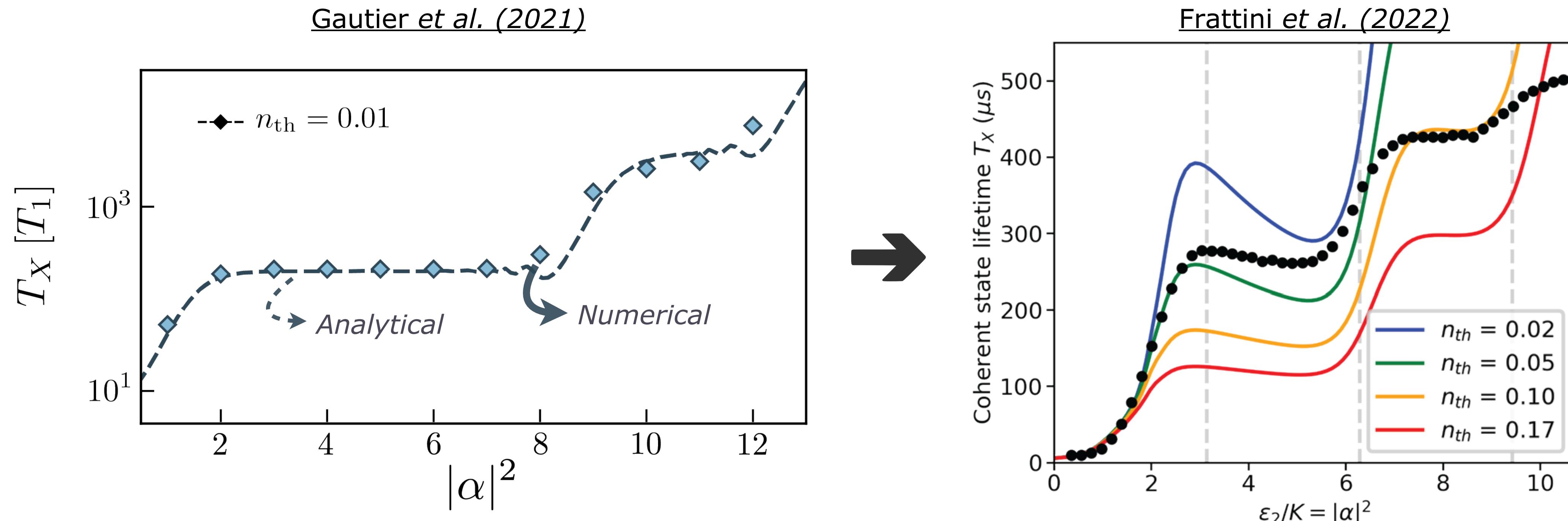
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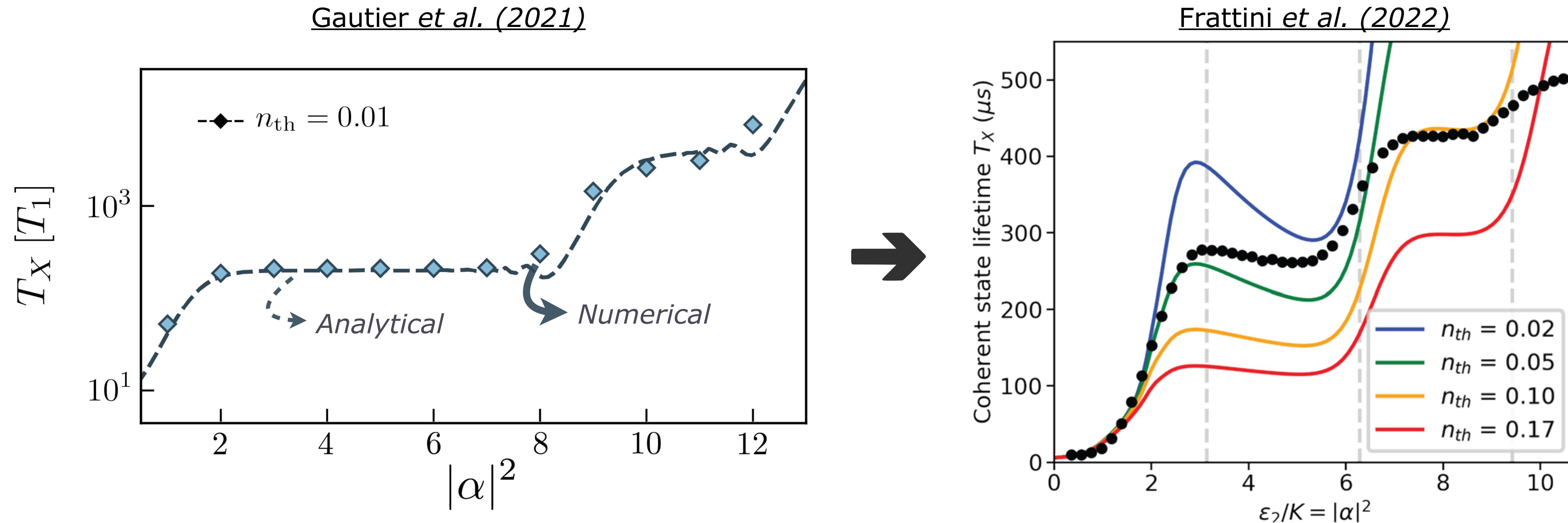
Bit-flip plateaus



Bit-flip plateaus



Bit-flip plateaus



How to retrieve an exponentially biased Kerr cat qubit?

Retrieving the exponential

Combining with two-photon dissipation

$$\frac{d\rho}{dt} = -i[H_{\text{Kerr}}, \rho] + \kappa_2 \mathcal{D}[a^2 - \alpha^2]\rho$$

Retrieving the exponential

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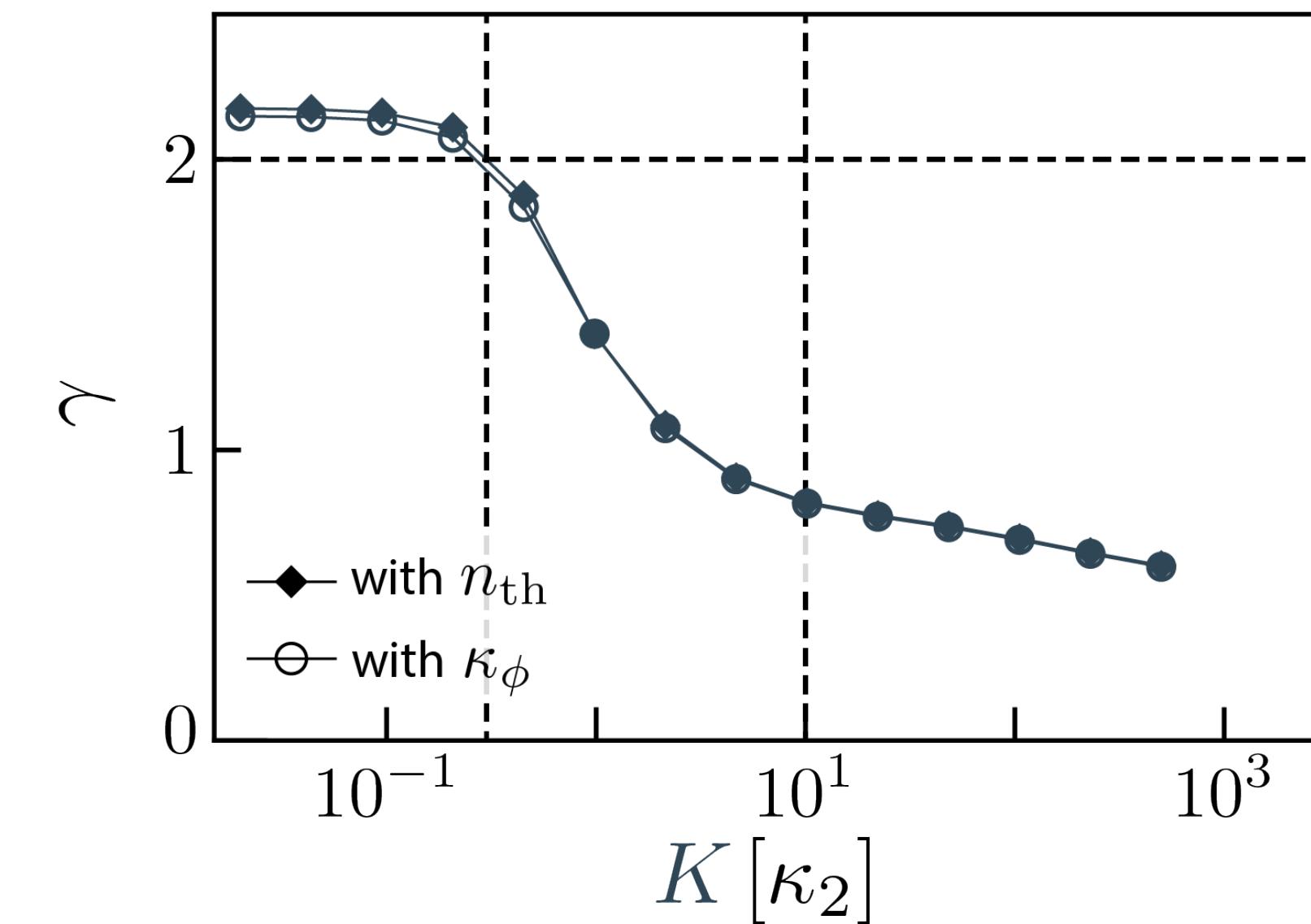
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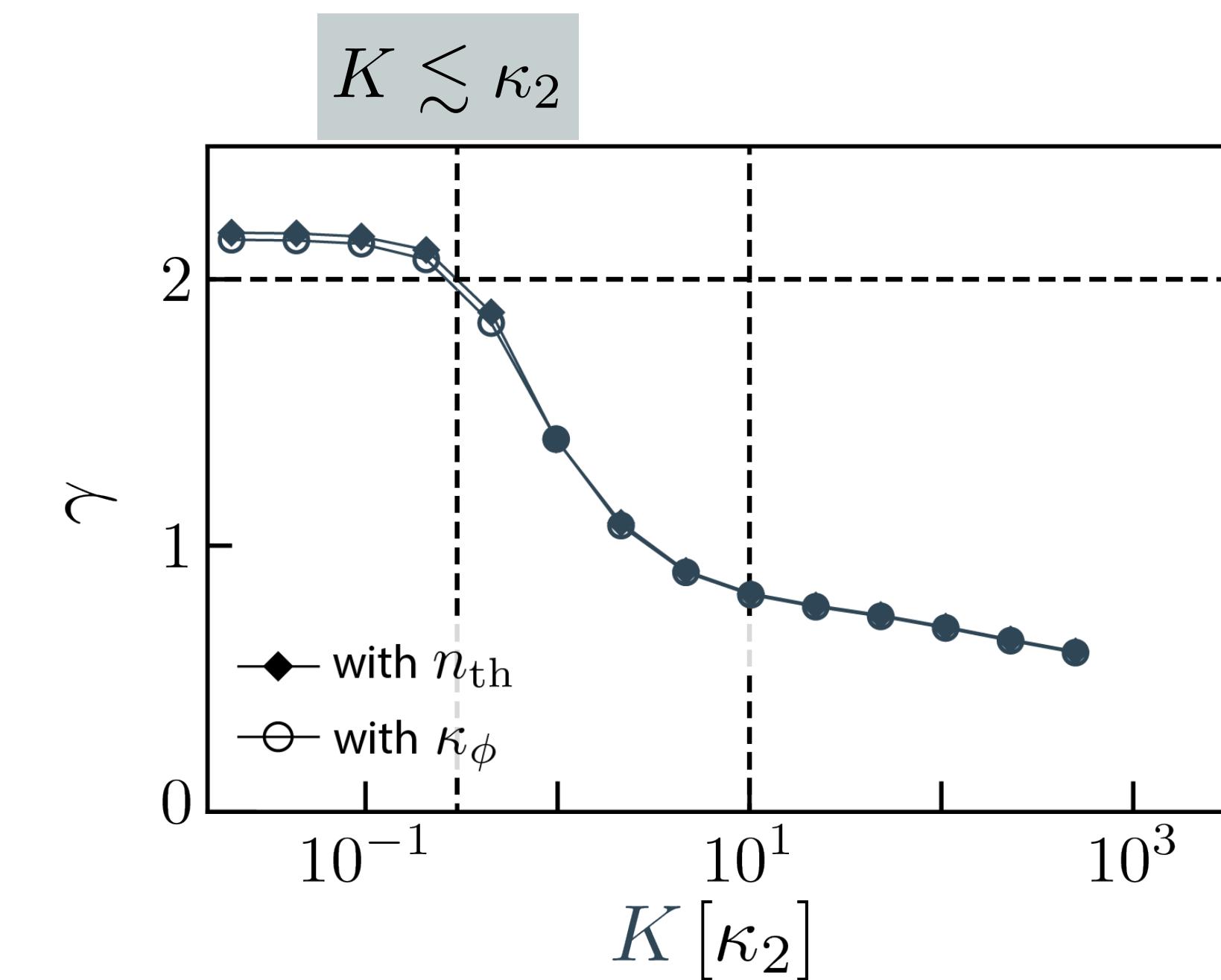


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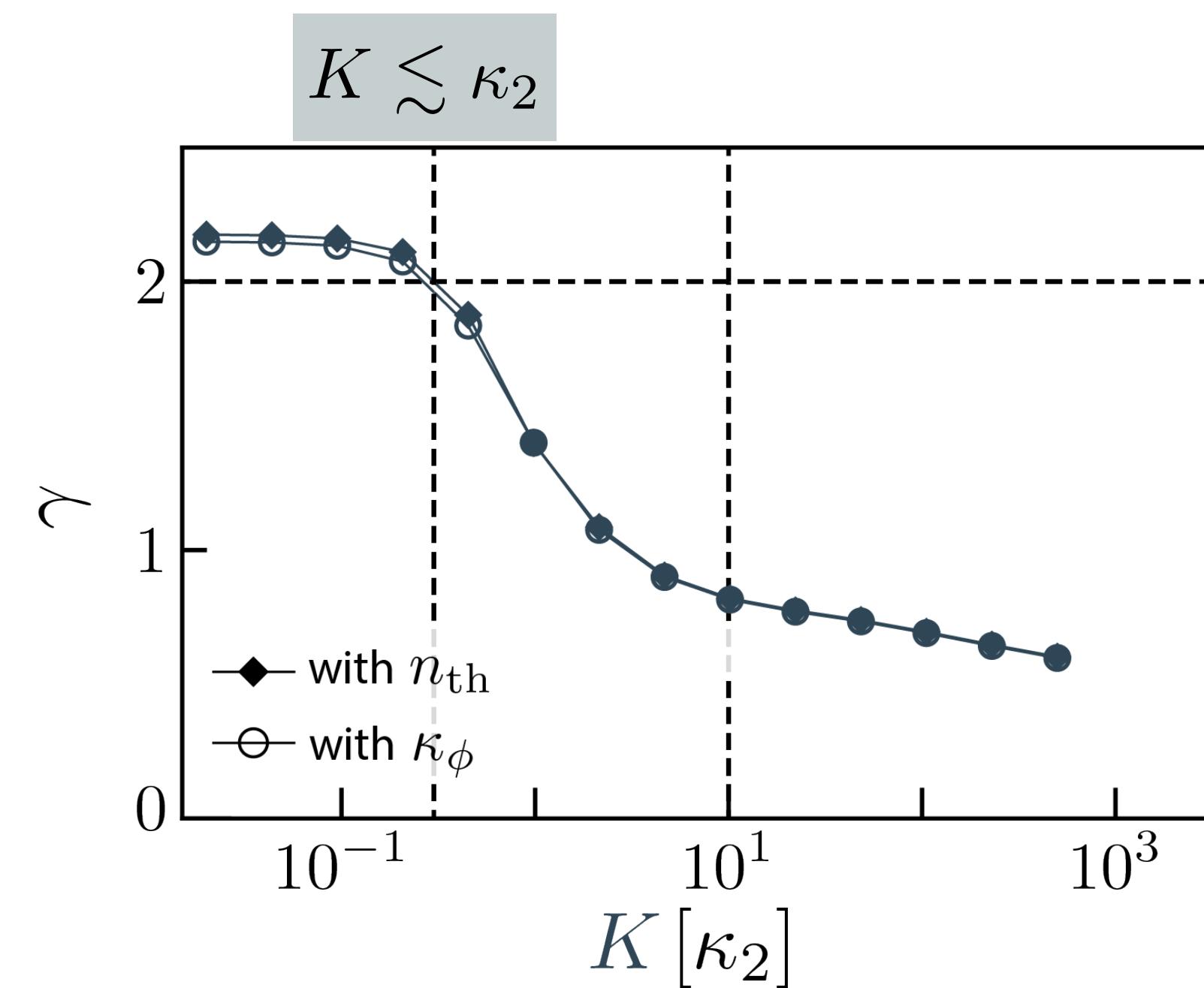
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The two-photon exchange Hamiltonian

$$\frac{d\rho}{dt} = -i[H_{\text{TPE}}, \rho] + \kappa_2 \mathcal{D}[a^2 - \alpha^2]\rho$$

with $H_{\text{TPE}} = g_2(a^2 - \alpha^2)\sigma_+ + g_2^*(a^{\dagger 2} - \alpha^{*2})\sigma_-$

- Jaynes-Cummings-like interaction between a two-photon memory and a qubit



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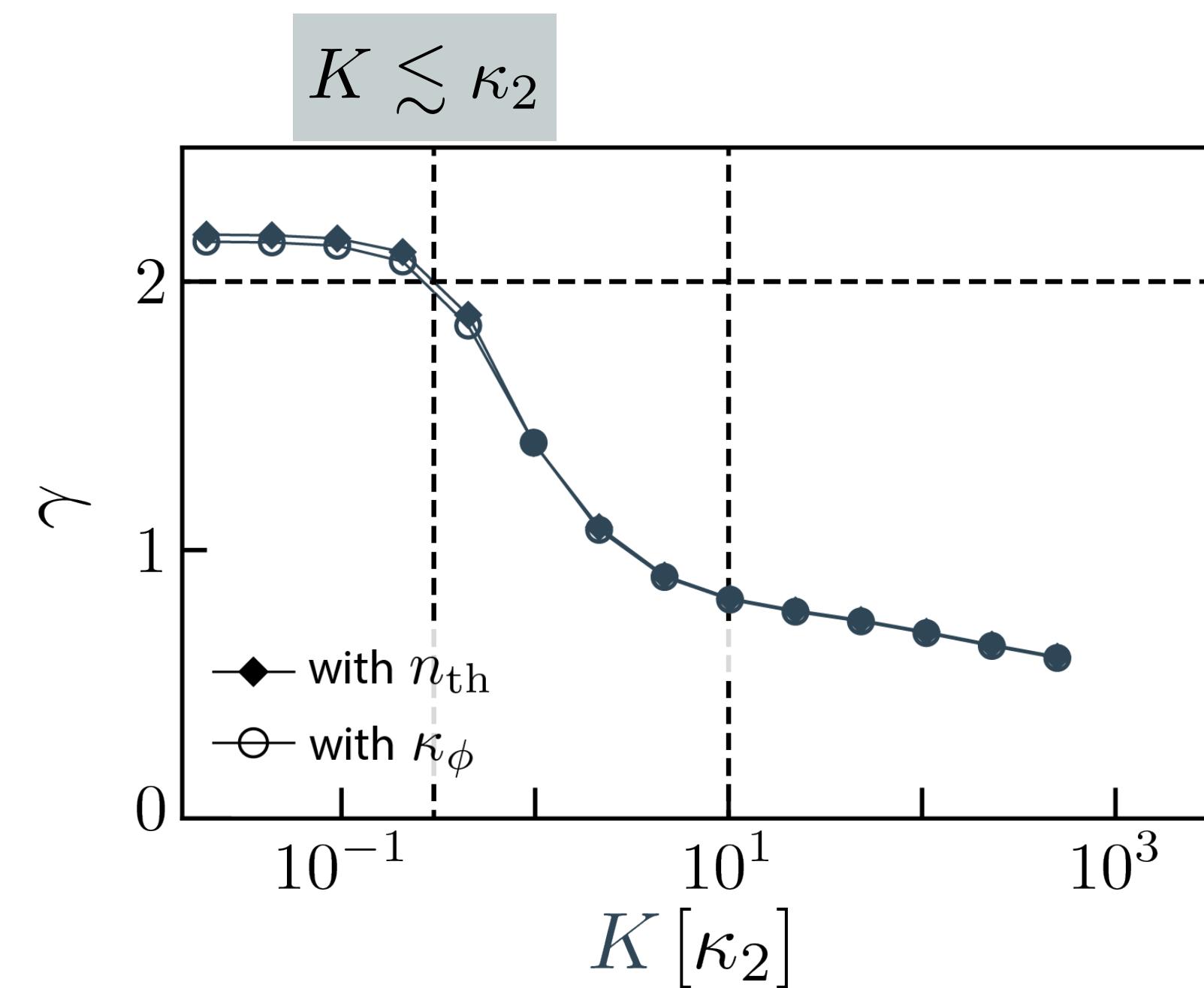
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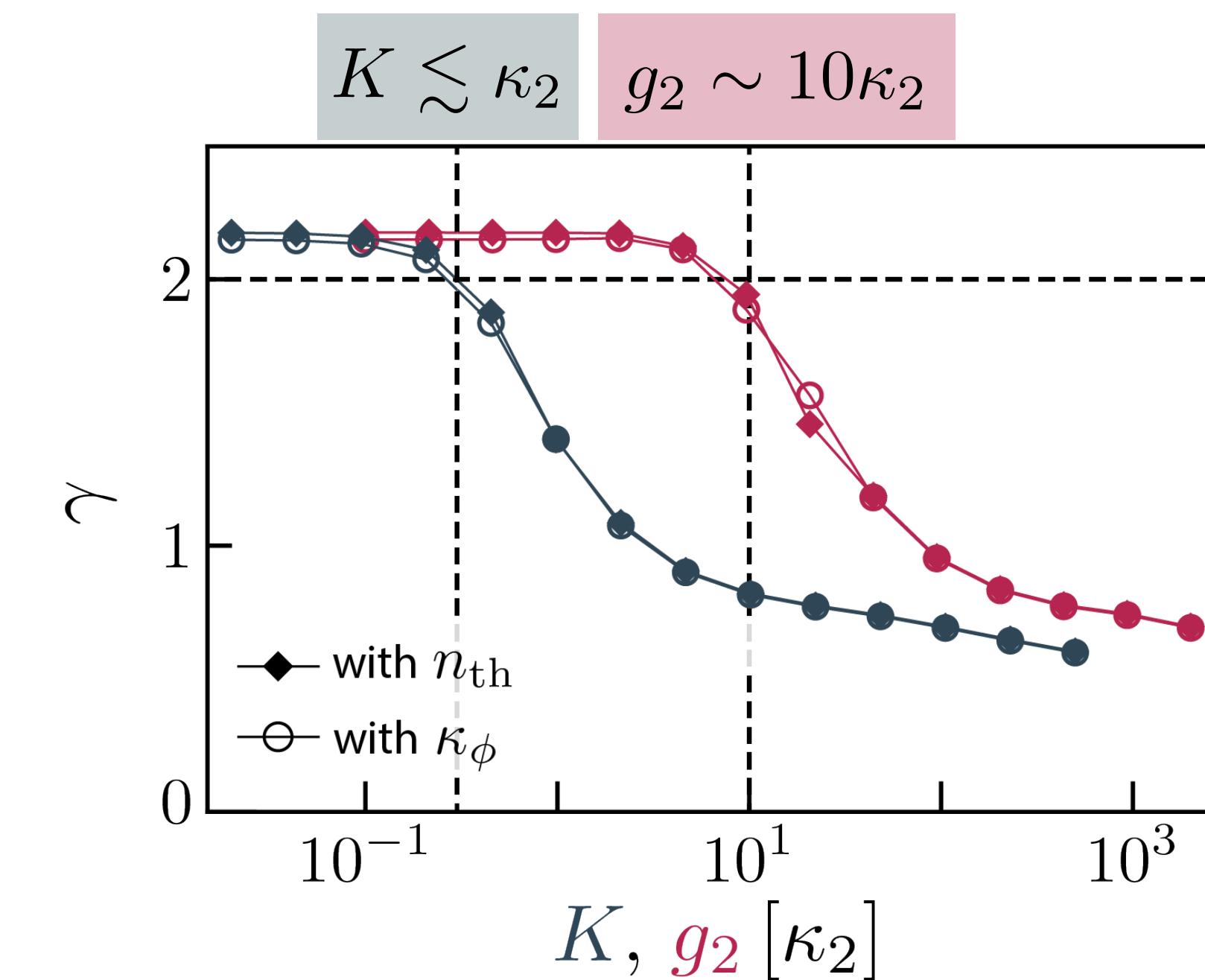
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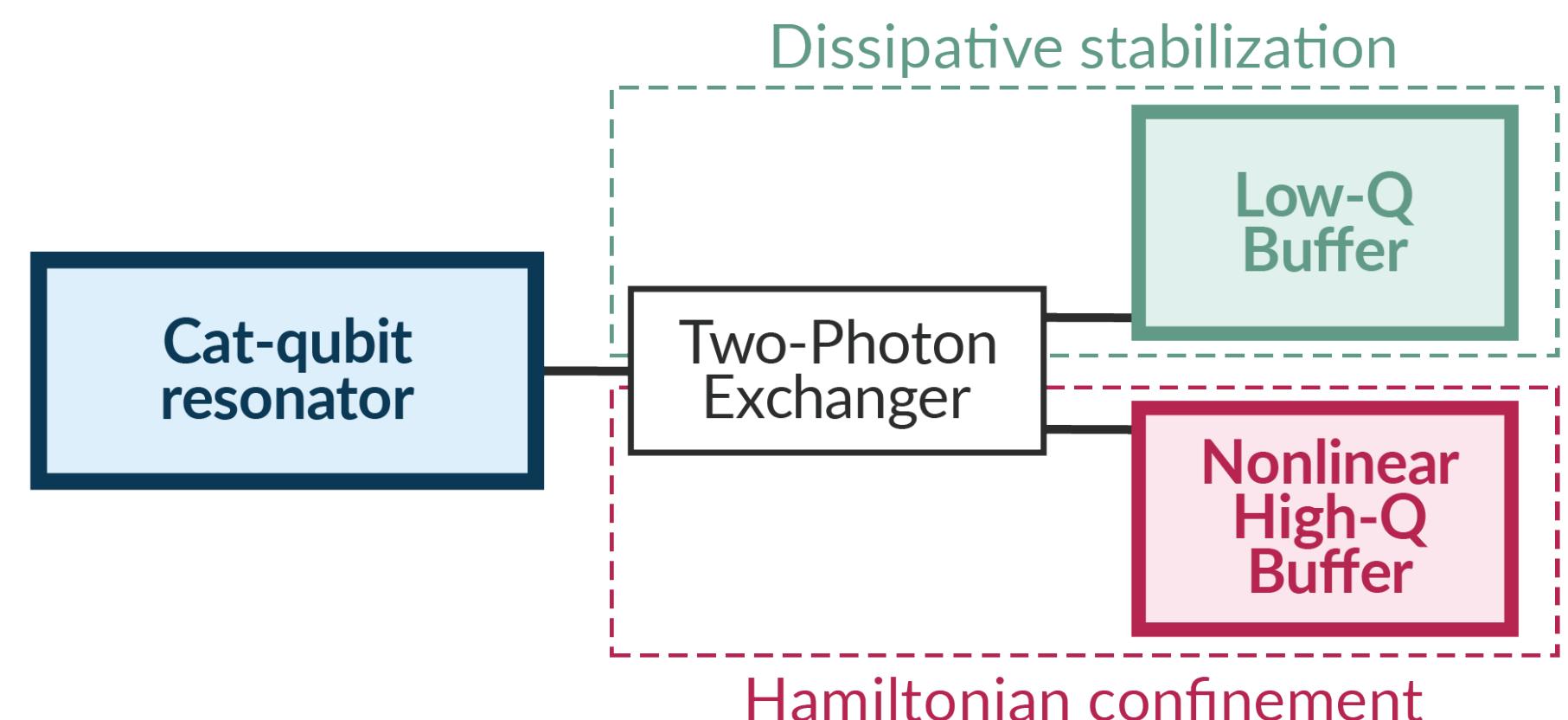
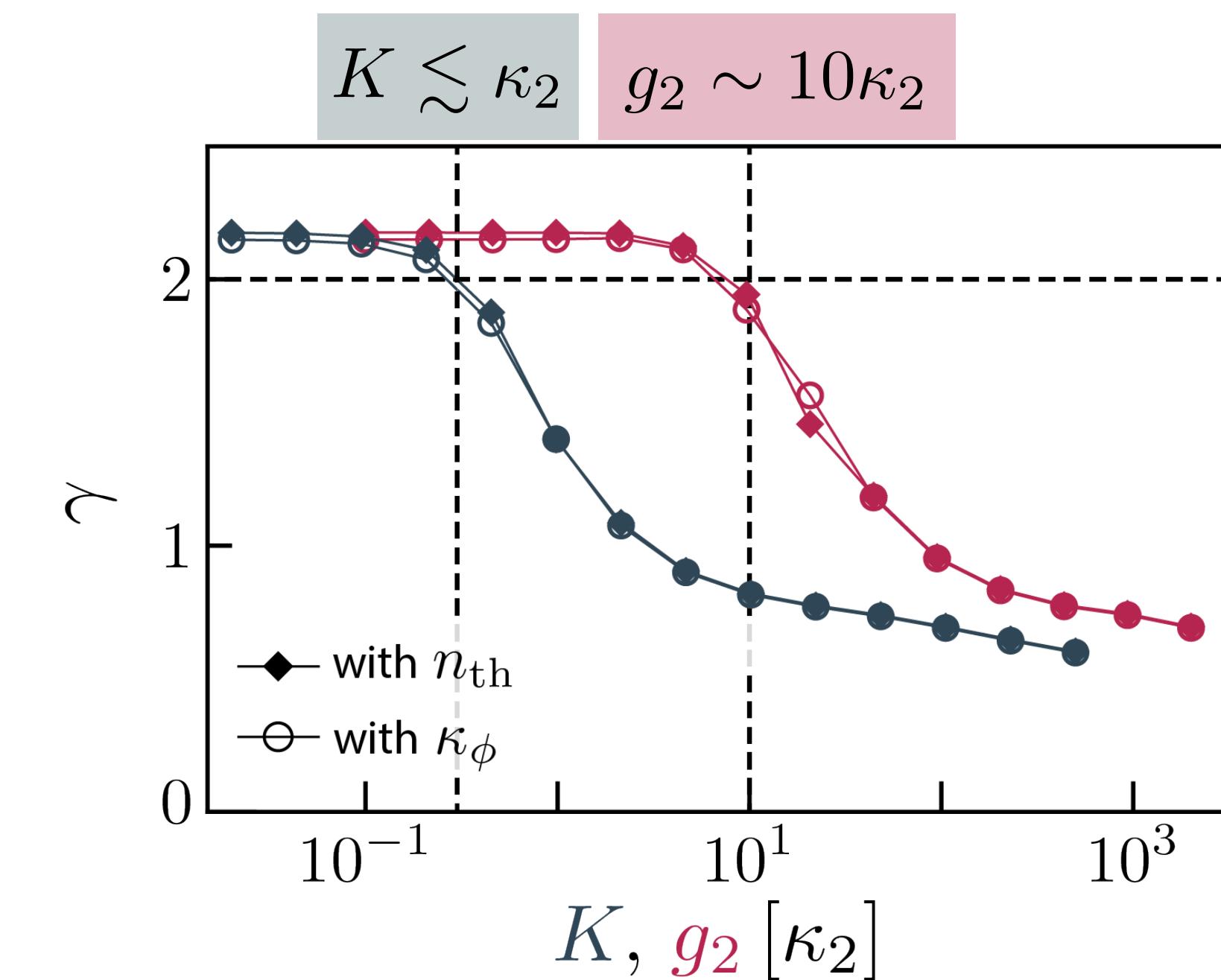
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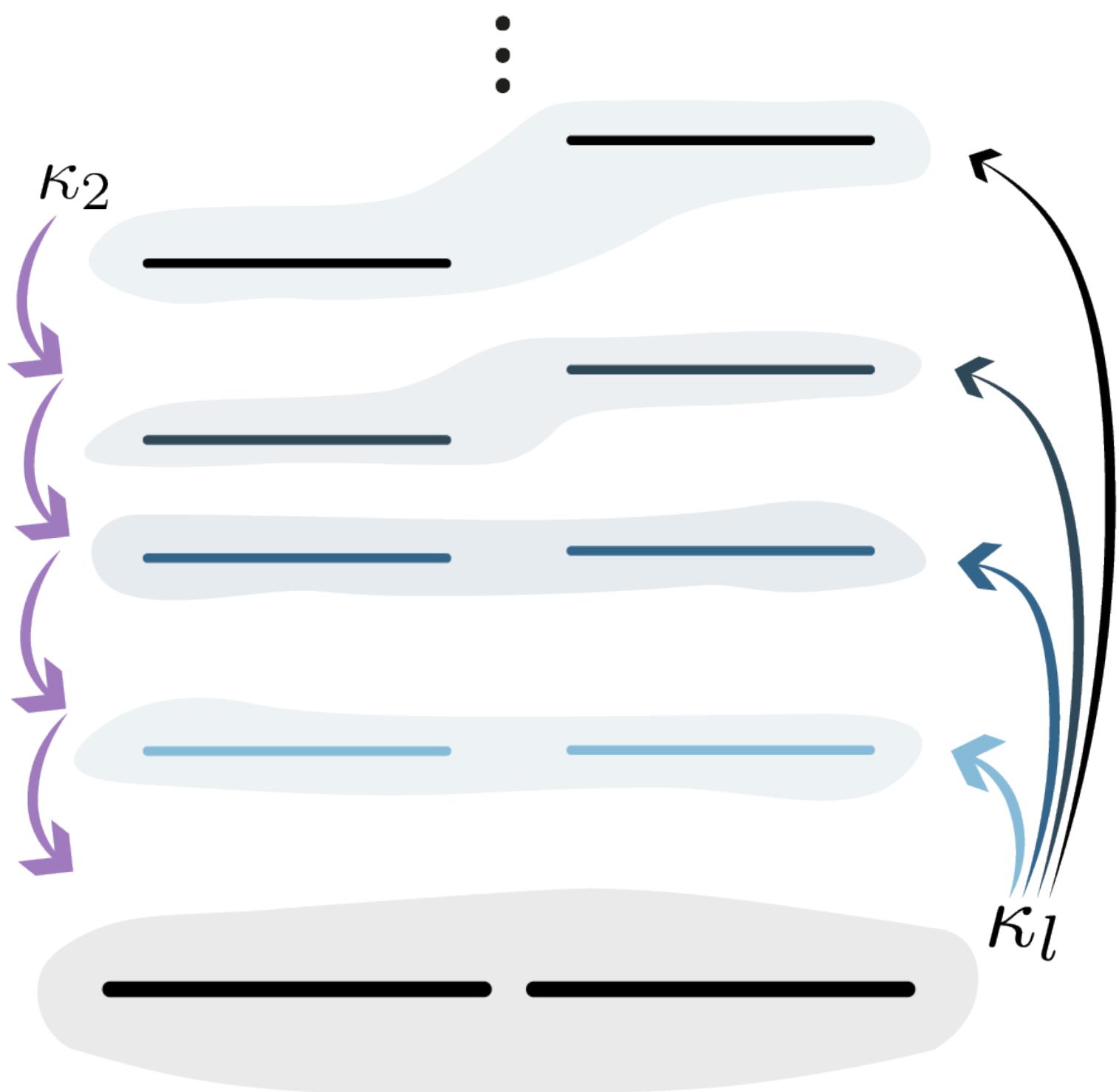
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- Jaynes-Cummings-like interaction between a two-photon memory and a qubit
- Eigenenergies $E_n/g_2 = \pm\sqrt{e_n/K}$
- Easier to engineer together with two-photon dissipation



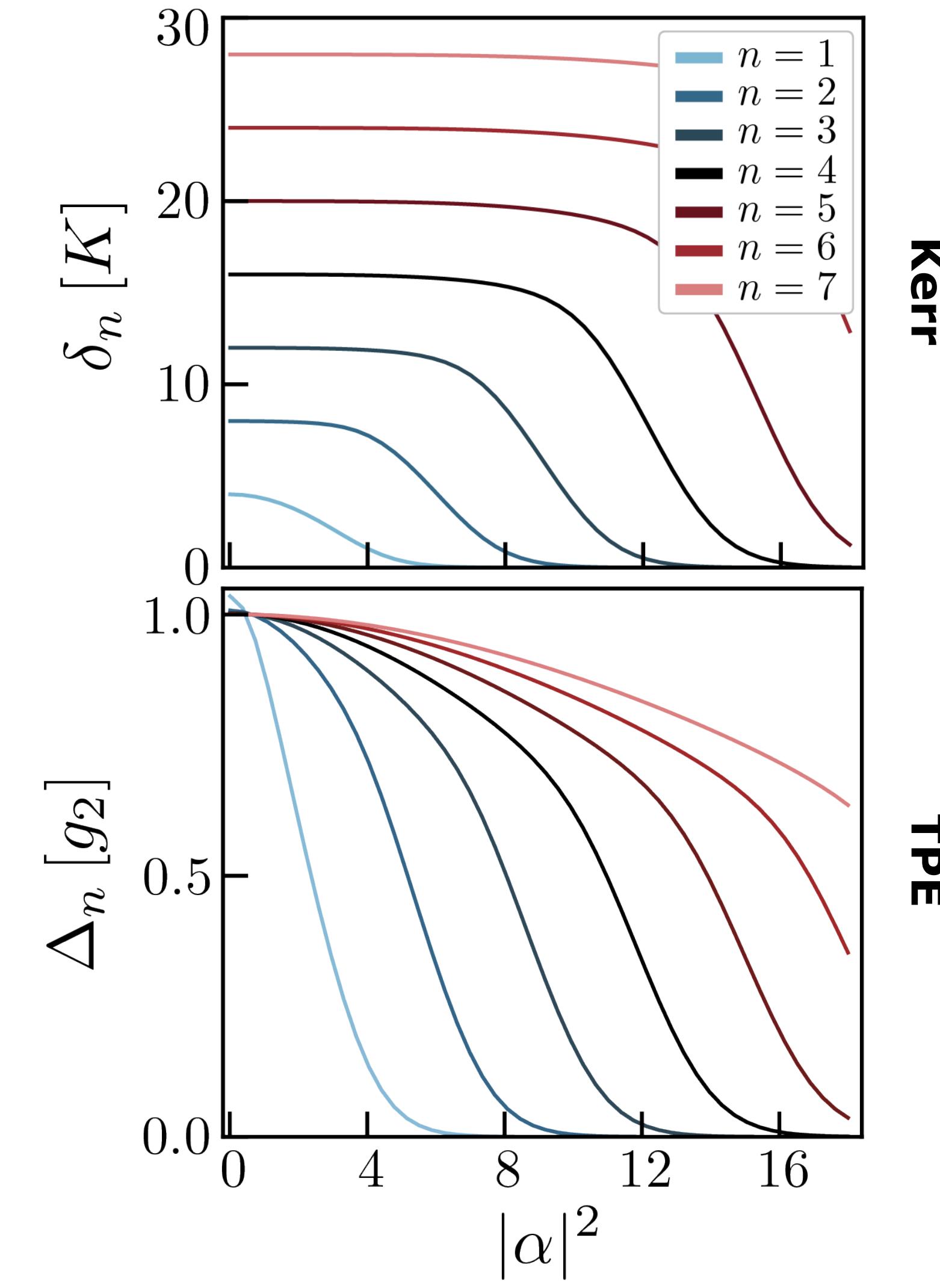
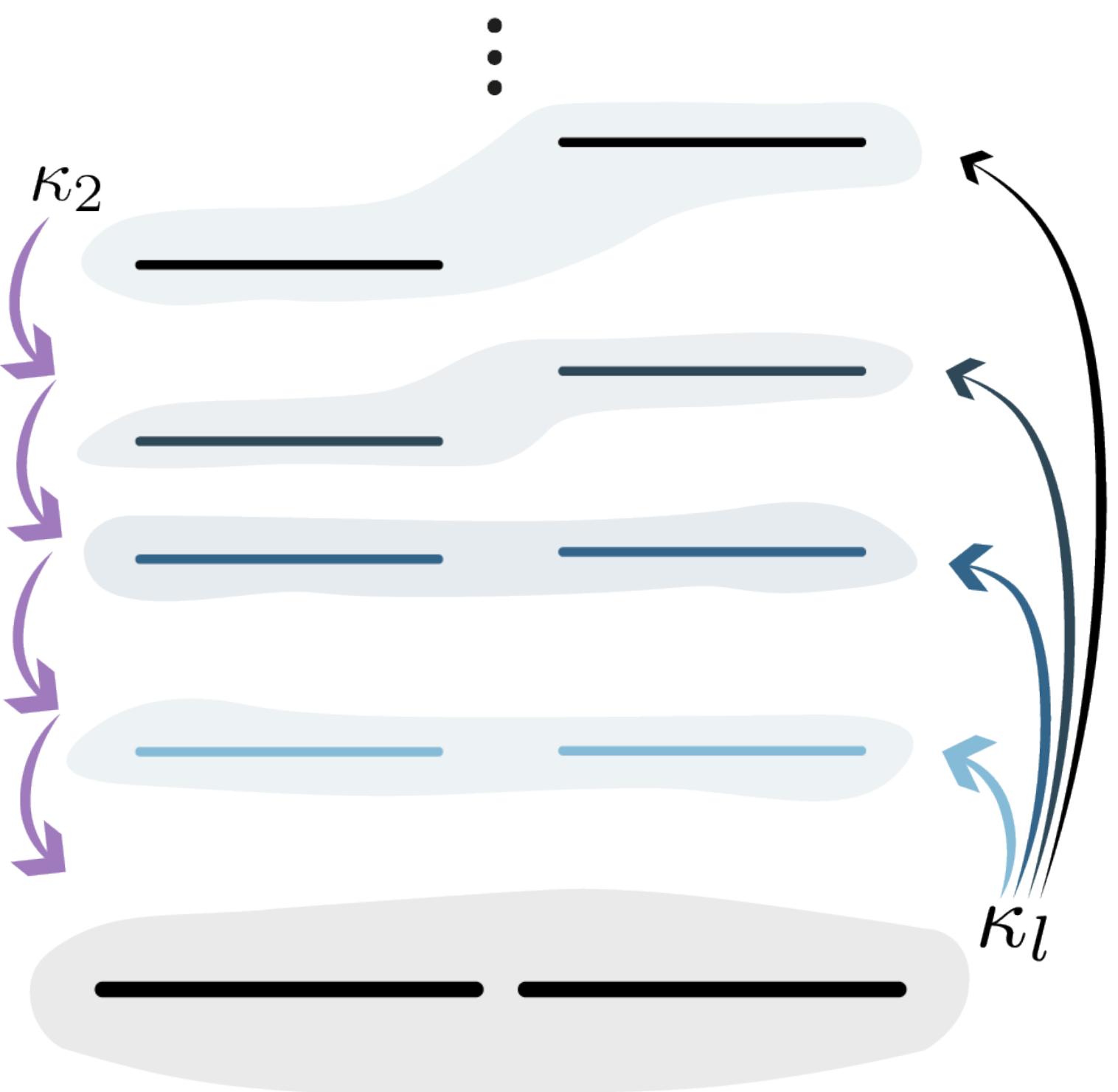
Combined Hamiltonian and dissipative confinement

- Competition between “heating” by leakage, and “cooling” by two-photon dissipation



Combined Hamiltonian and dissipative confinement

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Lifetime and controllability of cat qubits

Kerr cat qubits

Lifetime
→ $|\alpha|^2$

1. Limited by leakage
2. Weak exponential scaling

[Putterman et al., 2021]

[Gautier et al., 2021]

[Frattini et al., 2022]

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Dissipative cat qubits

Exponential scaling

[Mirrahimi et al., 2014]
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Lifetime and controllability of cat qubits

	Kerr cat qubits	Dissipative cat qubits
Lifetime → $ \alpha ^2$	<ol style="list-style-type: none">1. Limited by leakage2. Weak exponential scaling <p>[Puttermann et al., 2021] [Gautier et al., 2021] [Frattini et al., 2022]</p>	Exponential scaling
Controllability → T_{gate}	Adiabatic theorem ↓ Exponential scaling	[Mirrahimi et al., 2014] [Lescanne et al., 2019]

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Exponential scaling

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[Xu et al., 2022]

Dissipative cat qubits

Exponential scaling

[Mirrahimi et al., 2014]
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Linear scaling

[Mirrahimi et al., 2014]
[Guillaud et al., 2019]
[Gautier et al., 2022]

Summary of PhD contributions

Work on cat qubits

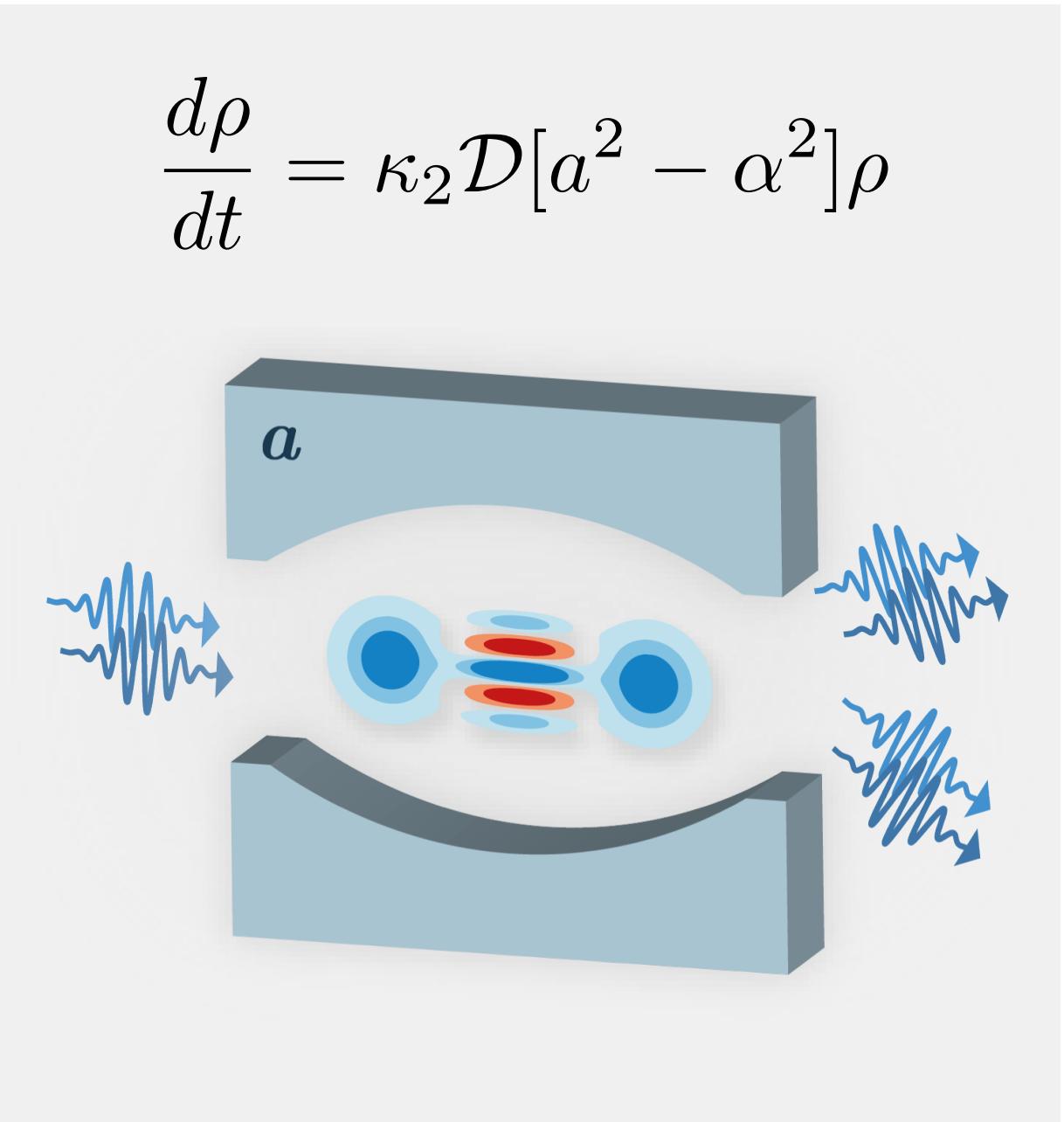
- RG, A. Sarlette, M. Mirrahimi, *Combined dissipative and Hamiltonian confinement of cat qubits*, PRX Quantum (2021)
- D. Ruiz, RG, J. Guillaud, M. Mirrahimi, *Two-photon driven Kerr quantum oscillator with multiple spectral degeneracies*, Phys. Rev. A (2022)
- RG, M. Mirrahimi, A. Sarlette, *Designing high-fidelity Zeno gates for dissipative cat qubits*, PRX Quantum (2022)
- U. Réglade, A. Bocquet, RG, et al., *Quantum control of a cat-qubit with bit-flip times exceeding ten seconds*, arXiv (2023)

Work on optimal control of open quantum systems

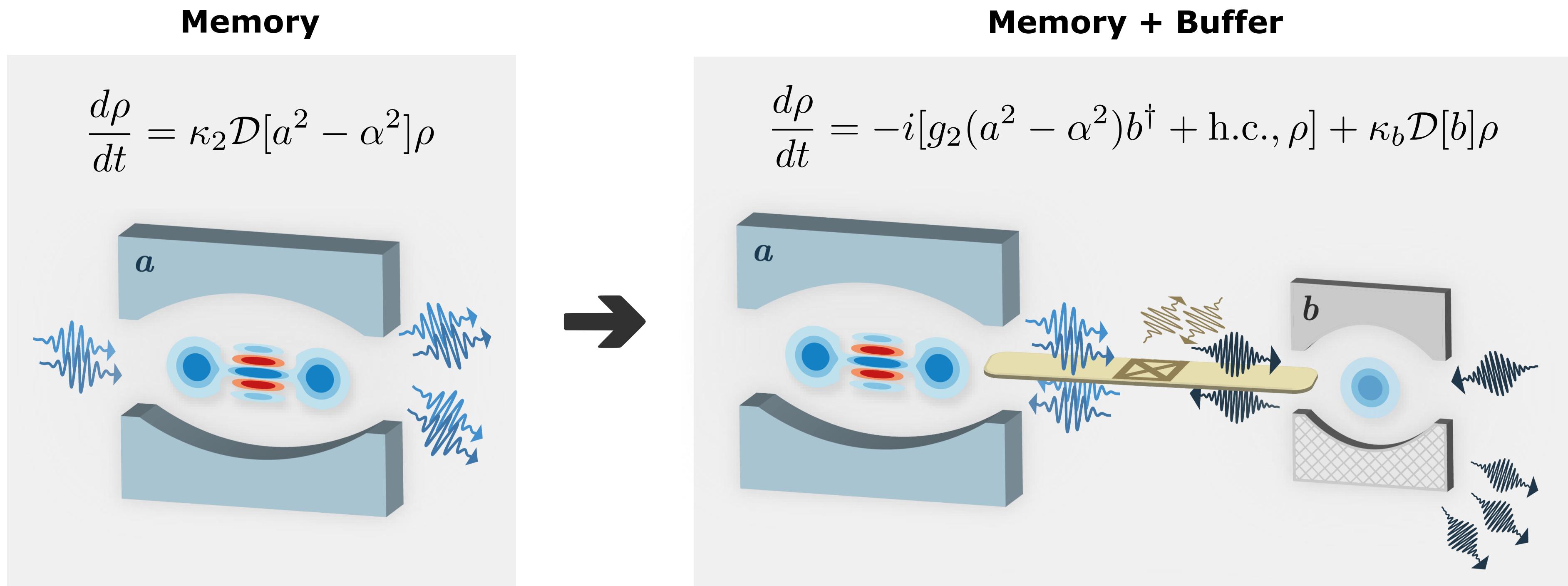
- RG, É. Genois, A. Blais, *Optimal readout and reset of a transmon*, in preparation
- P. Guilmin, RG, A. Bocquet, É. Genois, *dynamiqs: an open-source library for GPU-accelerated and differentiable quantum simulation*, in preparation

Reservoir engineering of two-photon dissipation

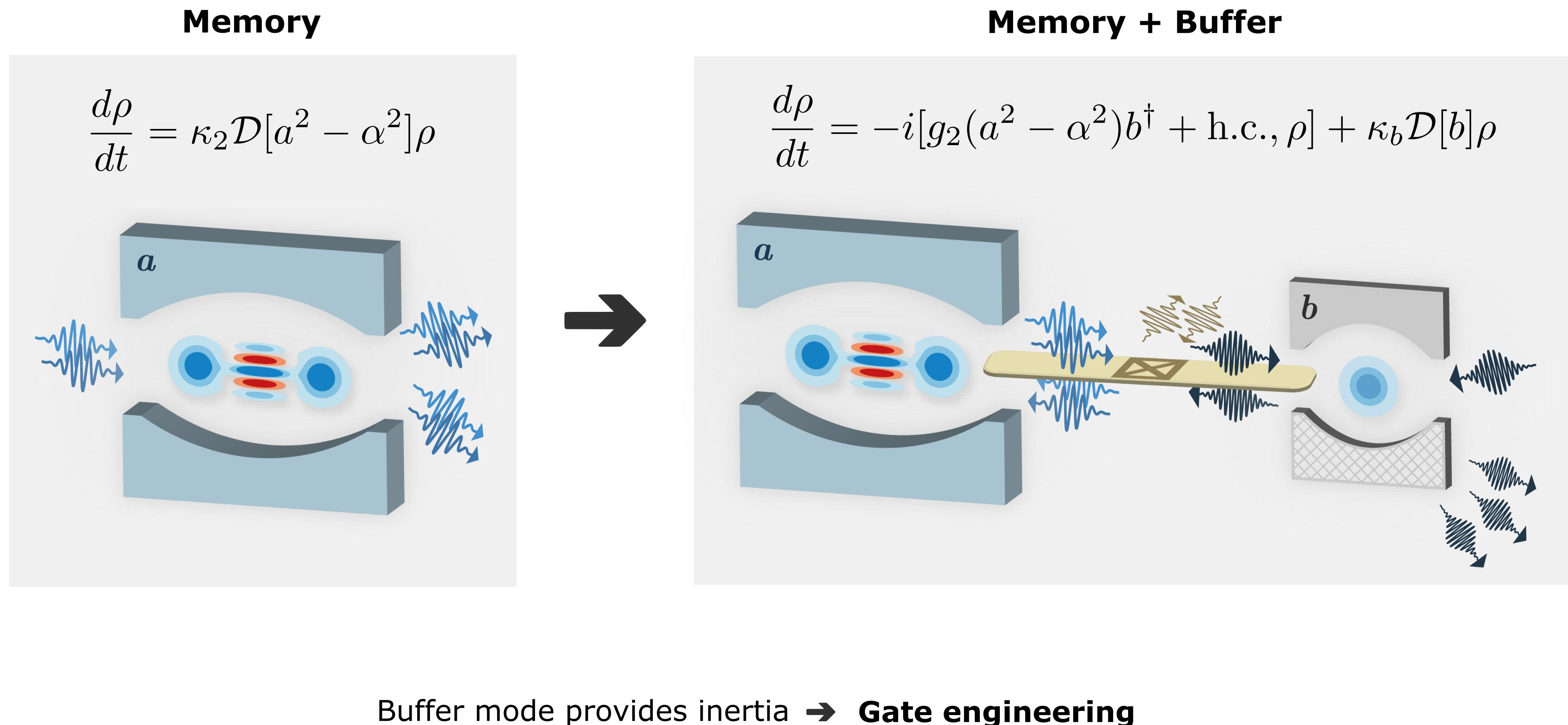
Memory



Reservoir engineering of two-photon dissipation



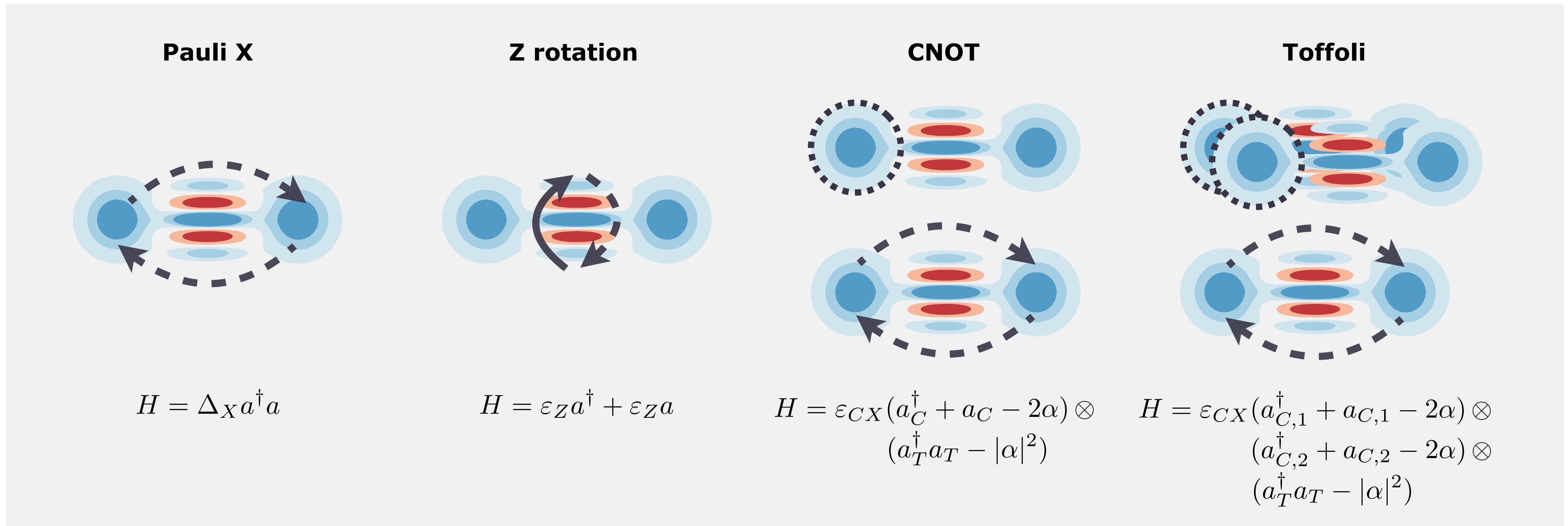
Reservoir engineering of two-photon dissipation



Buffer mode provides inertia → **Gate engineering**

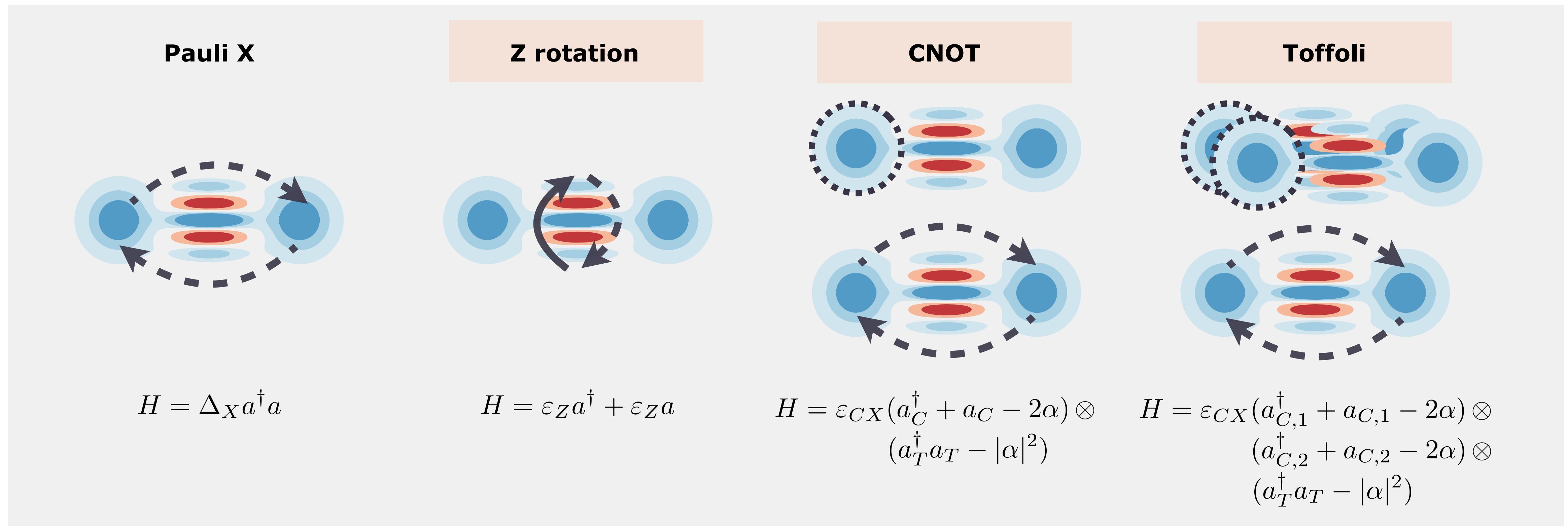
Gates based on the Zeno effect

Guillaud et al. (2019)



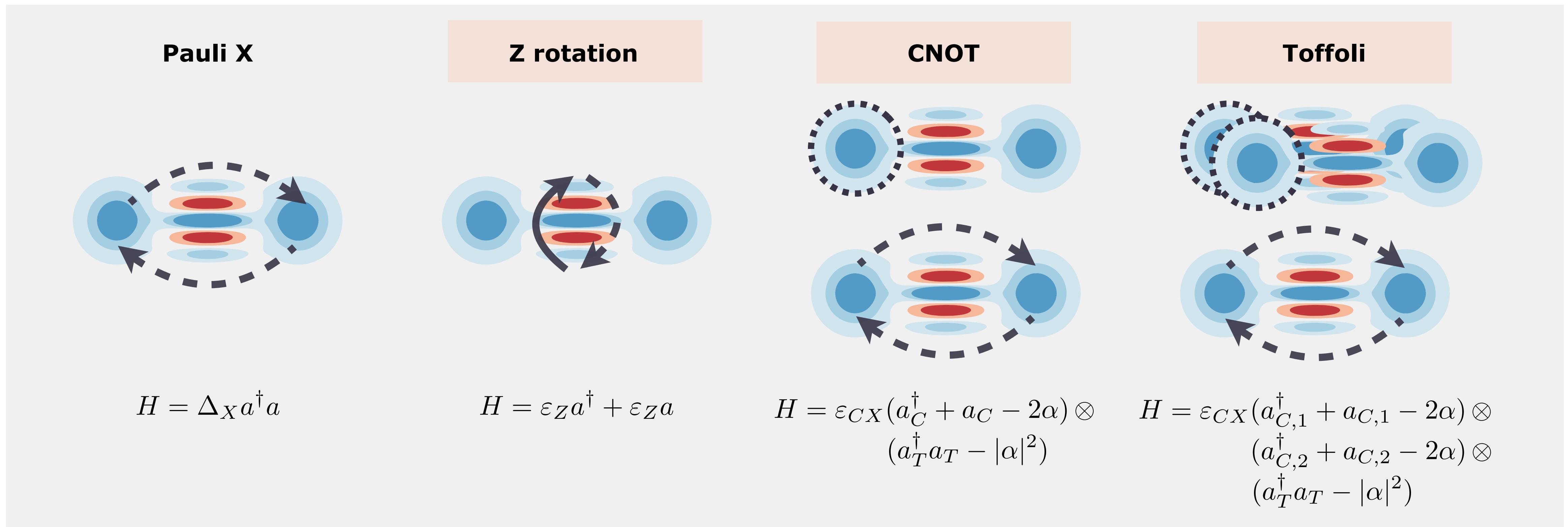
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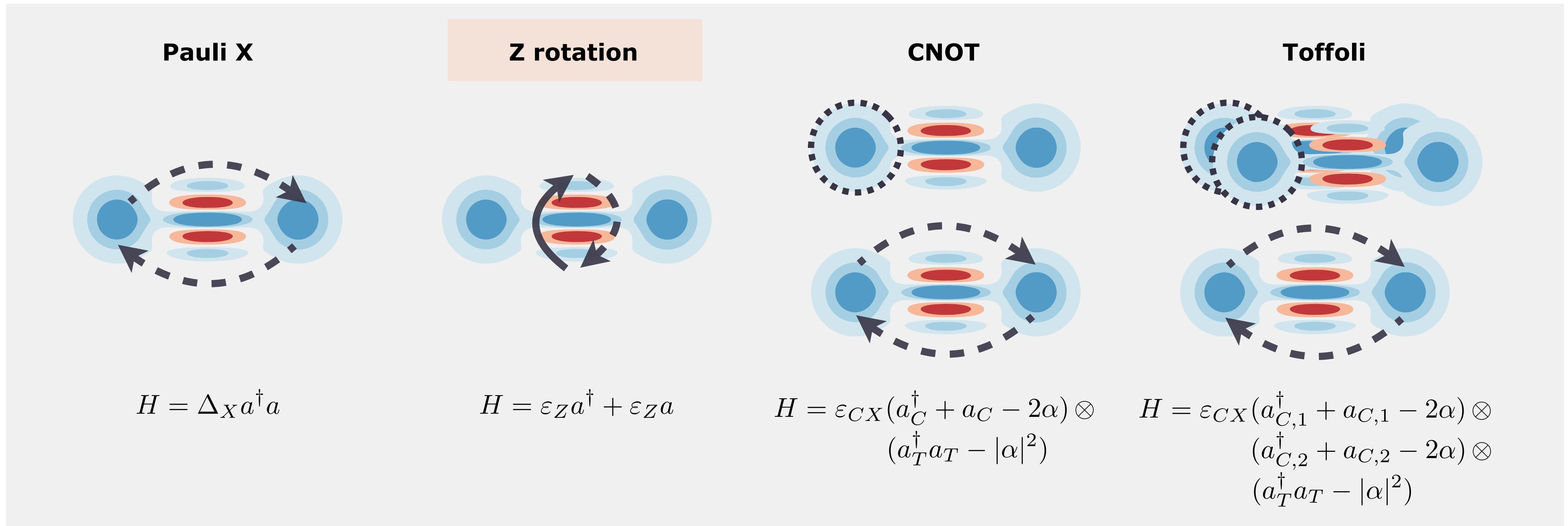


$$p_Z = \frac{\pi^2}{16|\alpha|^4 \kappa_2 T} + \kappa_1 |\alpha|^2 T$$

Gate errors Cavity lifetime

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Gate errors Cavity lifetime

Gate-induced errors are
only on **control** qubits

Unravelling the origin of gate errors

Hamiltonian of Z rotation gate

► $\hat{H} = g_2(\hat{a}^2 - \alpha^2)\hat{b}^\dagger + \varepsilon_Z\hat{a} + \text{h.c.}$

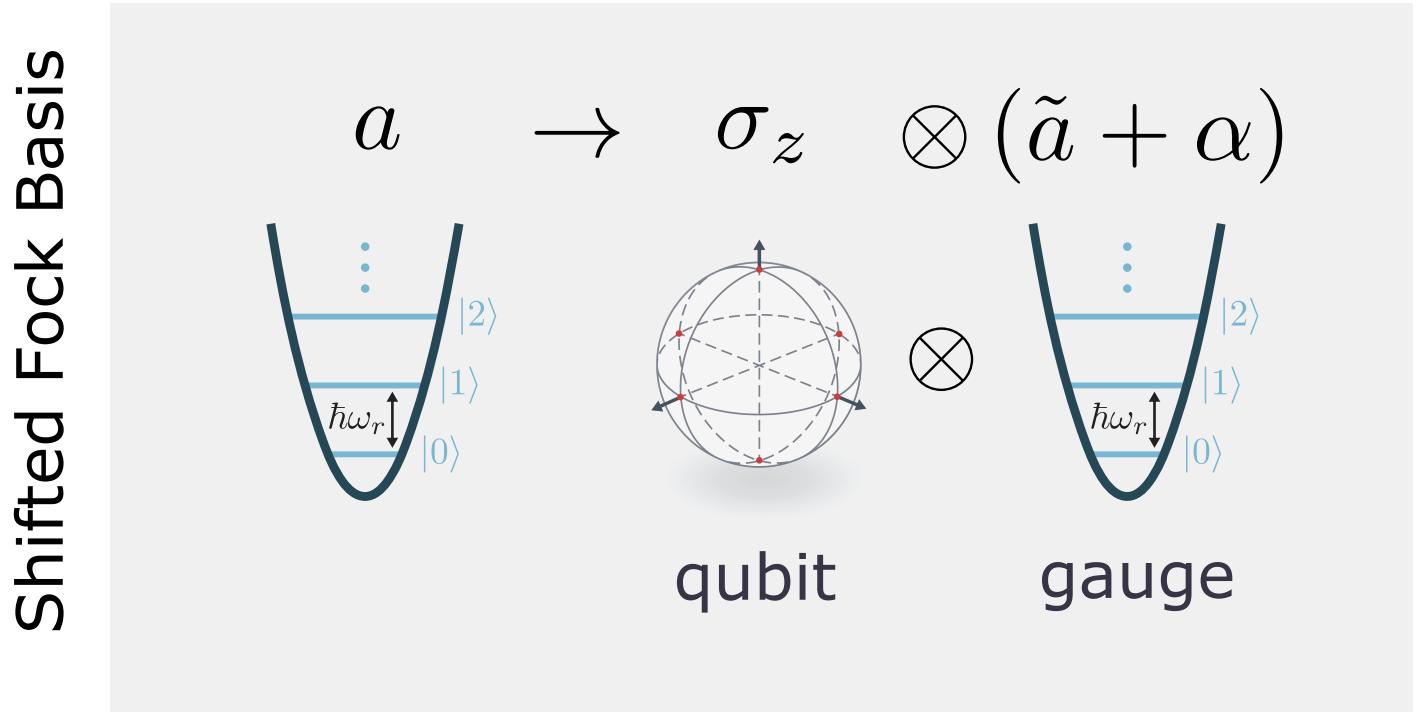
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Move in Shifted Fock Basis

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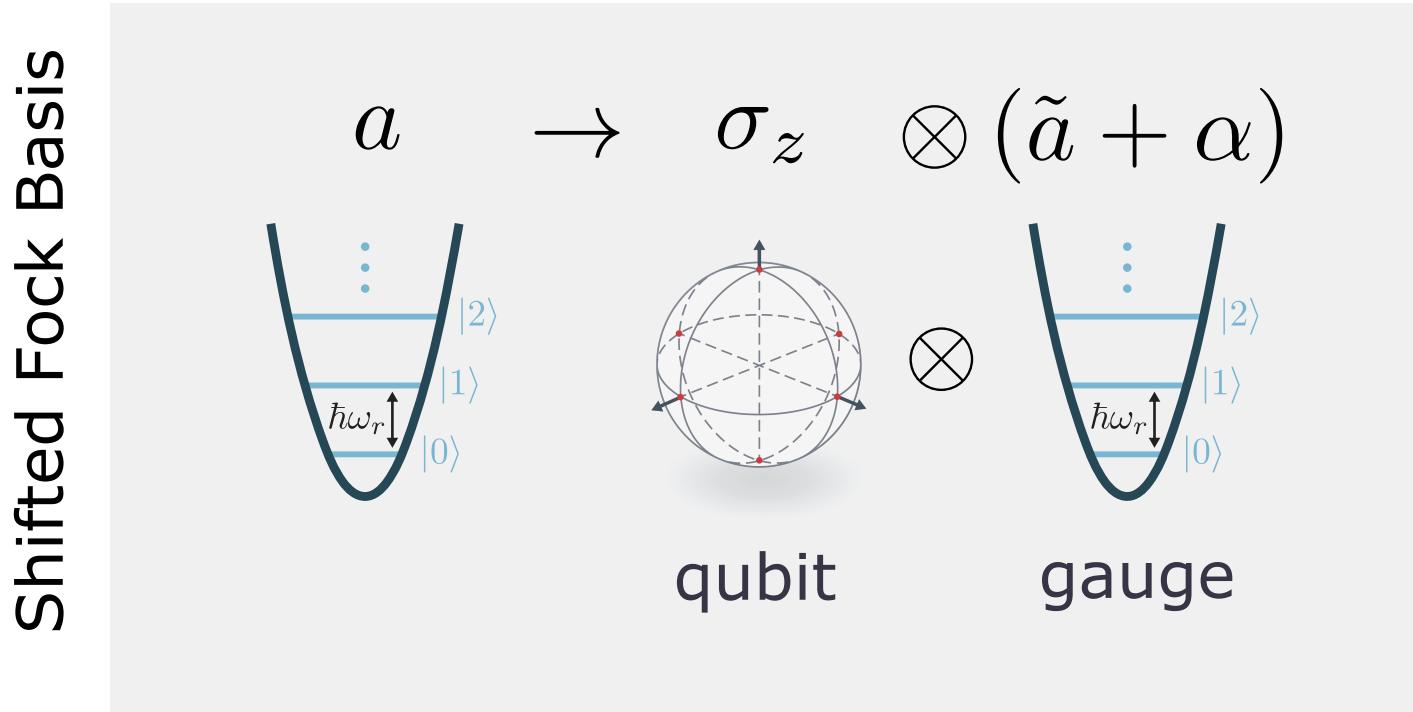
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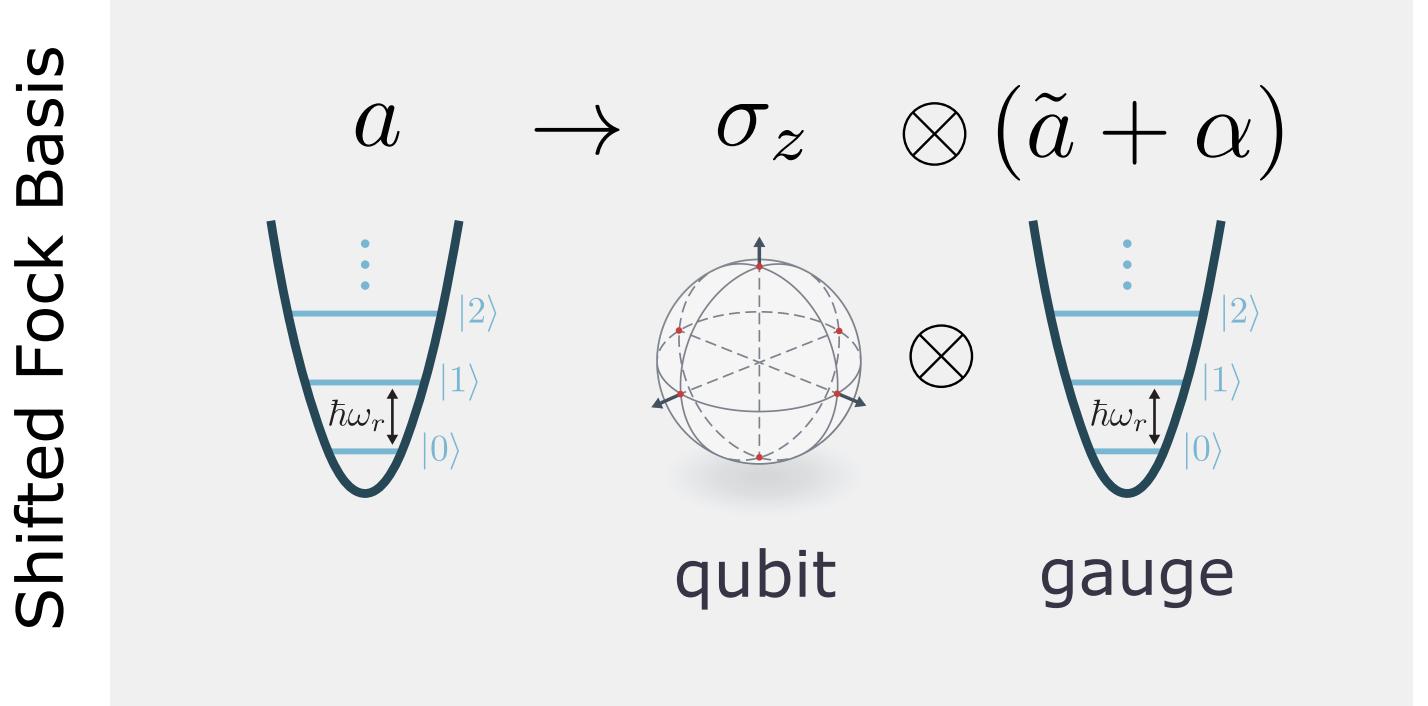
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Mean-field description of buffer dynamics

► $\ddot{b} + \frac{1}{2}\kappa_b \dot{b} + \nu^2 b = -\nu \varepsilon_Z \sigma_z$

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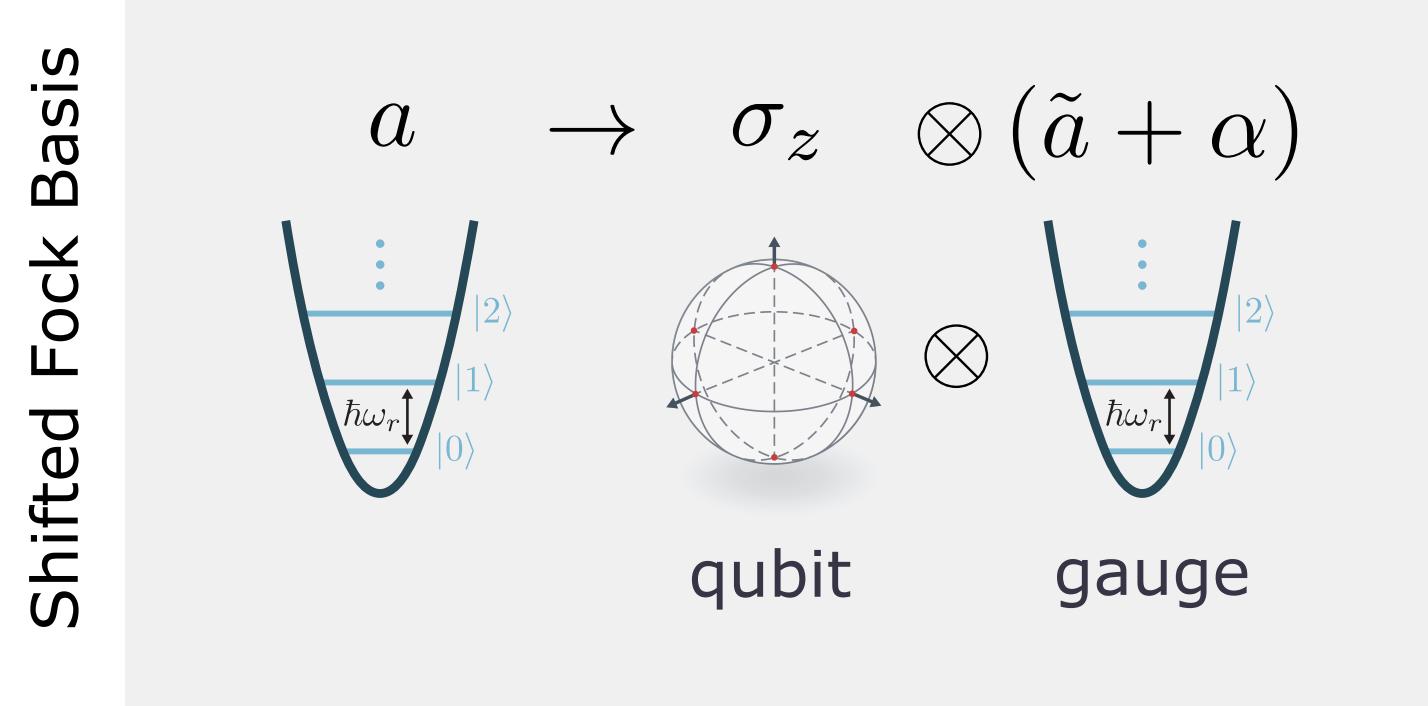
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Mean-field description of buffer dynamics

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Infinite-time dynamics

► $b \xrightarrow{t \rightarrow \infty} -\frac{\varepsilon_Z}{\nu} \sigma_z$



Unravelling the origin of gate errors

Hamiltonian of Z rotation gate

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Move in Shifted Fock Basis

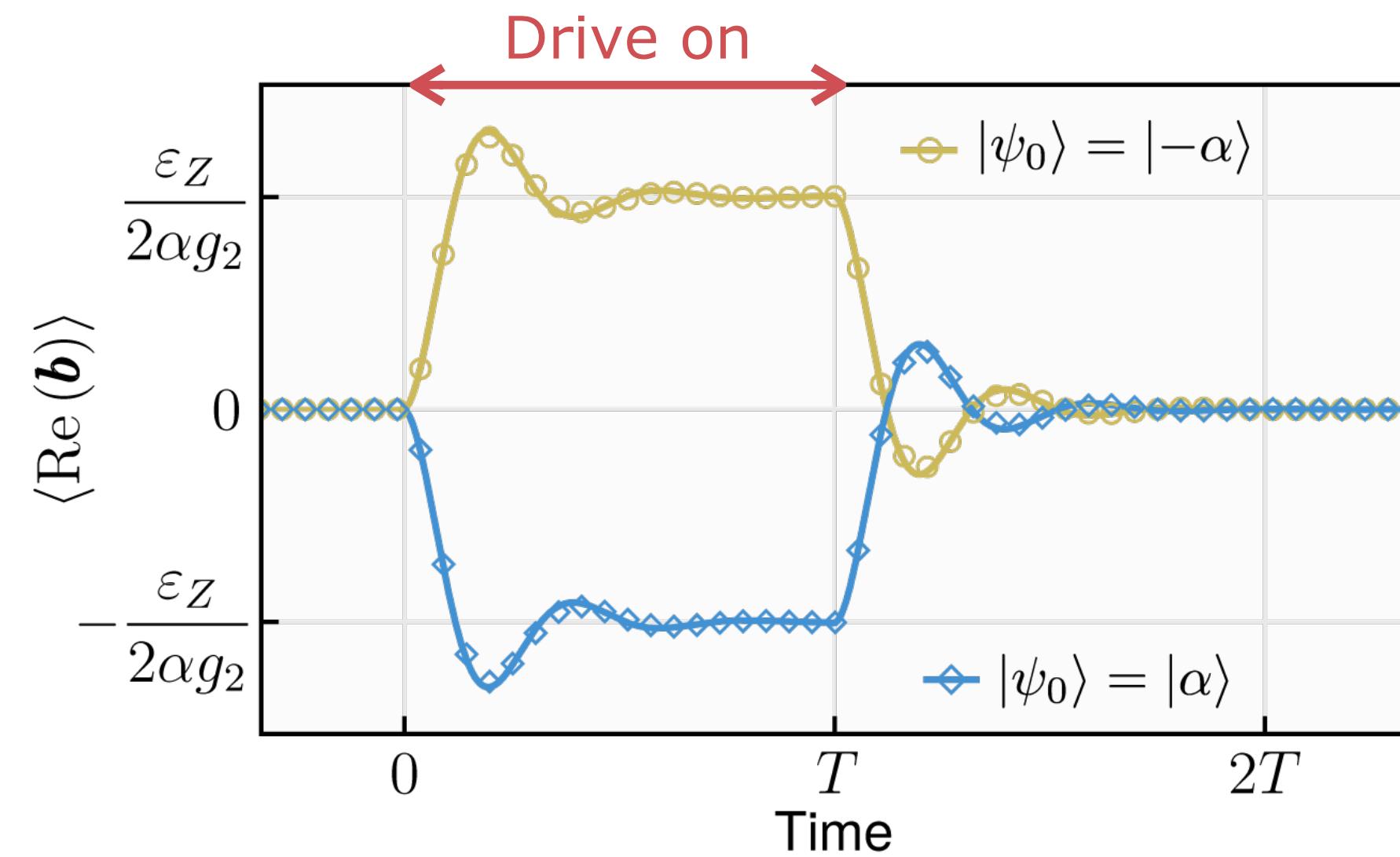
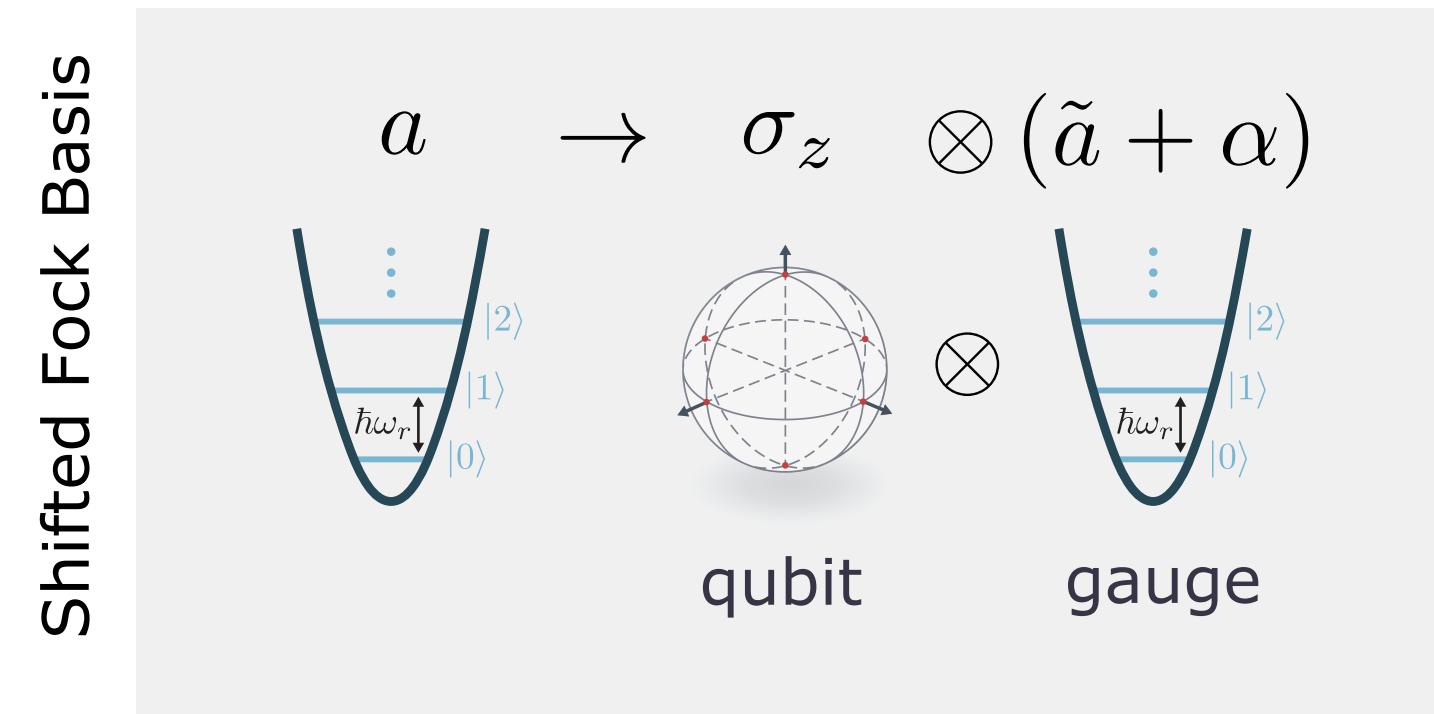
► $\hat{H} = g_2(\hat{\tilde{a}}^2 + 2\alpha\hat{\tilde{a}})\hat{b}^\dagger + \varepsilon_Z \hat{\sigma}_z \hat{a} + \varepsilon_Z \alpha \hat{\sigma}_z + \text{h.c.}$

Mean-field description of buffer dynamics

► $\ddot{b} + \frac{1}{2}\kappa_b \dot{b} + \nu^2 b = -\nu \varepsilon_Z \sigma_z$

Infinite-time dynamics

► $b \xrightarrow{t \rightarrow \infty} -\frac{\varepsilon_Z}{\nu} \sigma_z$



Unravelling the origin of gate errors

Hamiltonian of Z rotation gate

► $\hat{H} = g_2(\hat{a}^2 - \alpha^2)\hat{b}^\dagger + \varepsilon_Z \hat{a} + \text{h.c.}$

Move in Shifted Fock Basis

► $\hat{H} = g_2(\hat{\tilde{a}}^2 + 2\alpha\hat{\tilde{a}})\hat{b}^\dagger + \varepsilon_Z \hat{\sigma}_z \hat{a} + \varepsilon_Z \alpha \hat{\sigma}_z + \text{h.c.}$

Mean-field description of buffer dynamics

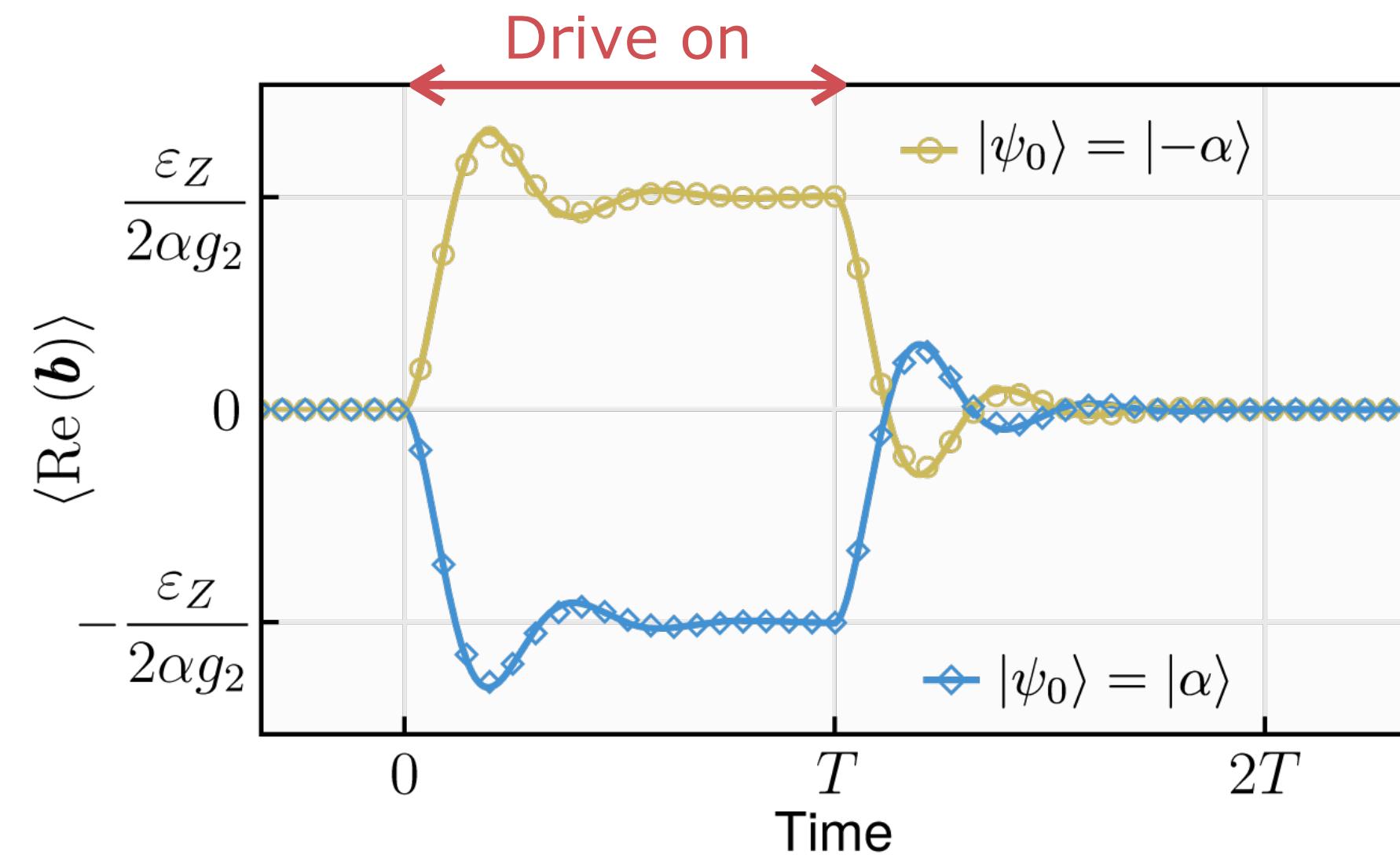
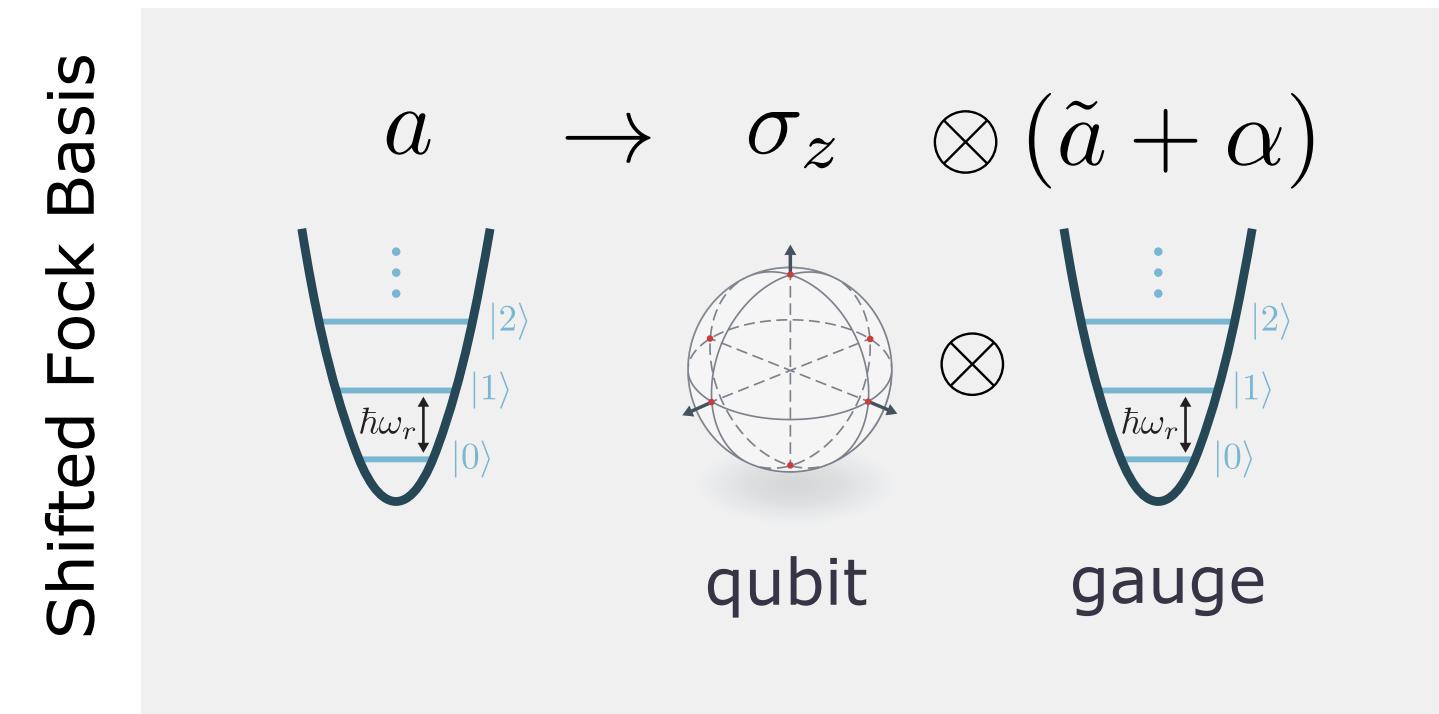
► $\ddot{b} + \frac{1}{2}\kappa_b \dot{b} + \nu^2 b = -\nu \varepsilon_Z \sigma_z$

Infinite-time dynamics

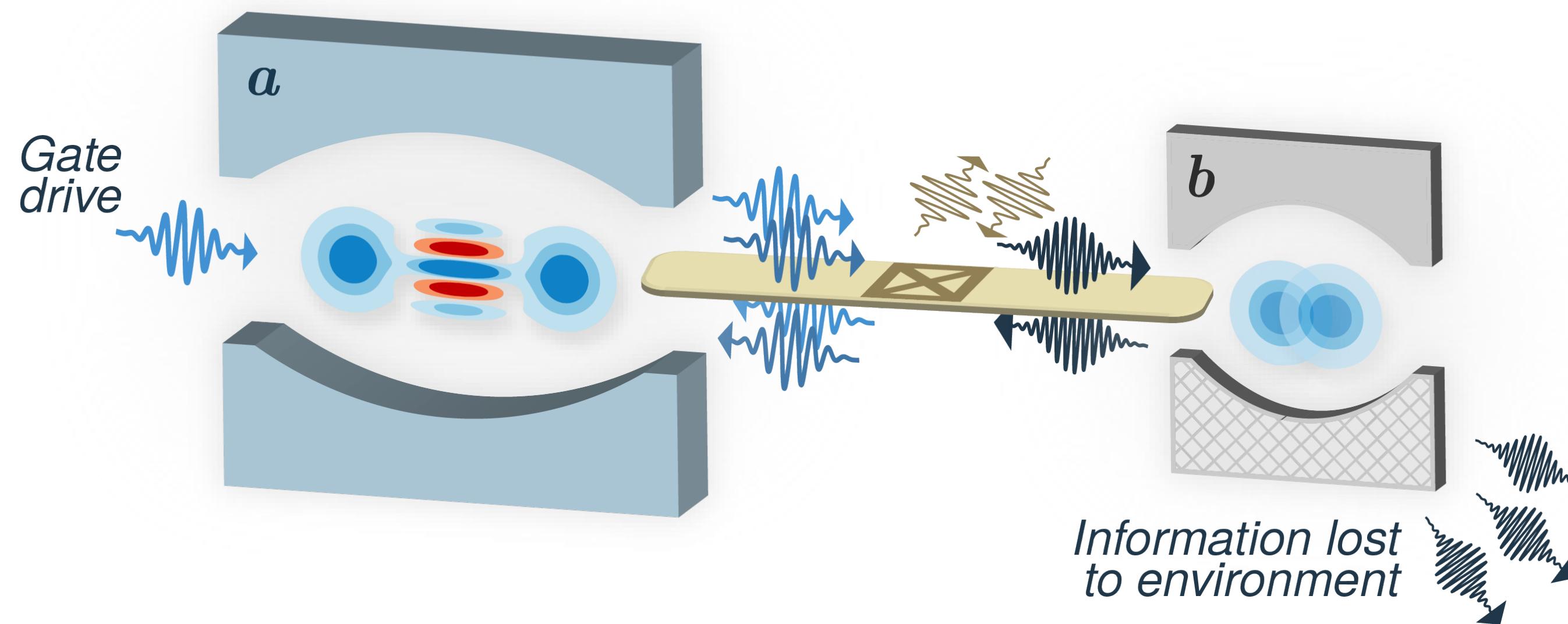
► $b \xrightarrow{t \rightarrow \infty} -\frac{\varepsilon_Z}{\nu} \sigma_z$

Generalises Zeno gate errors to 2 modes

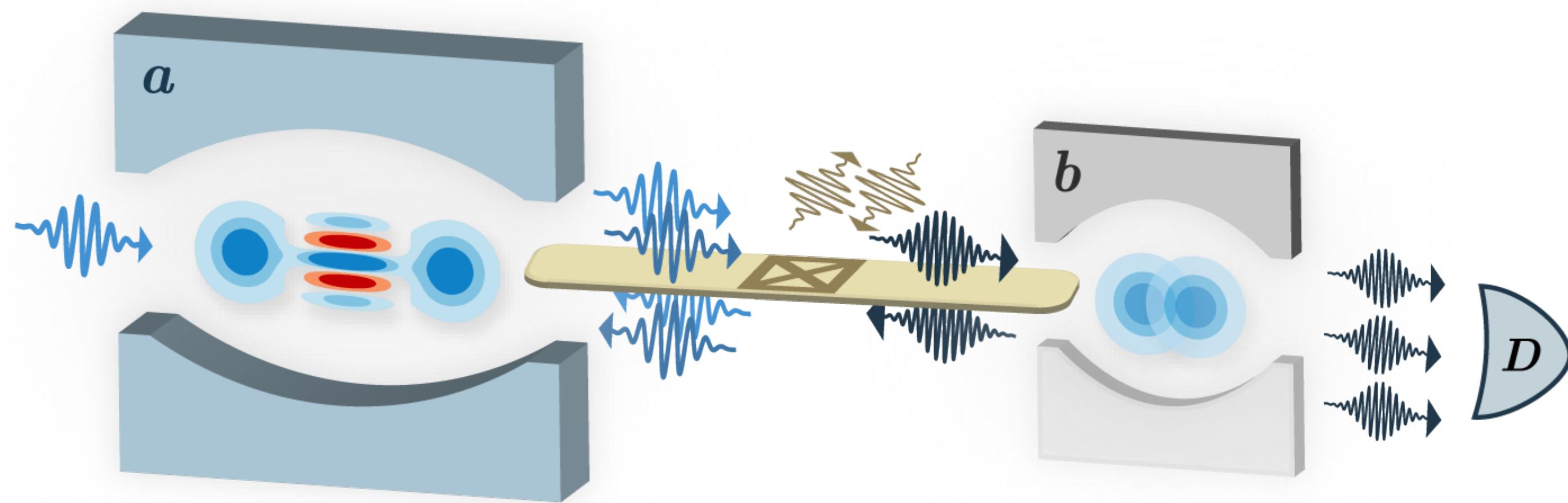
► $p_Z = \frac{\pi^2}{16|\alpha|^4 T} \frac{\kappa_b}{4g_2^2}$



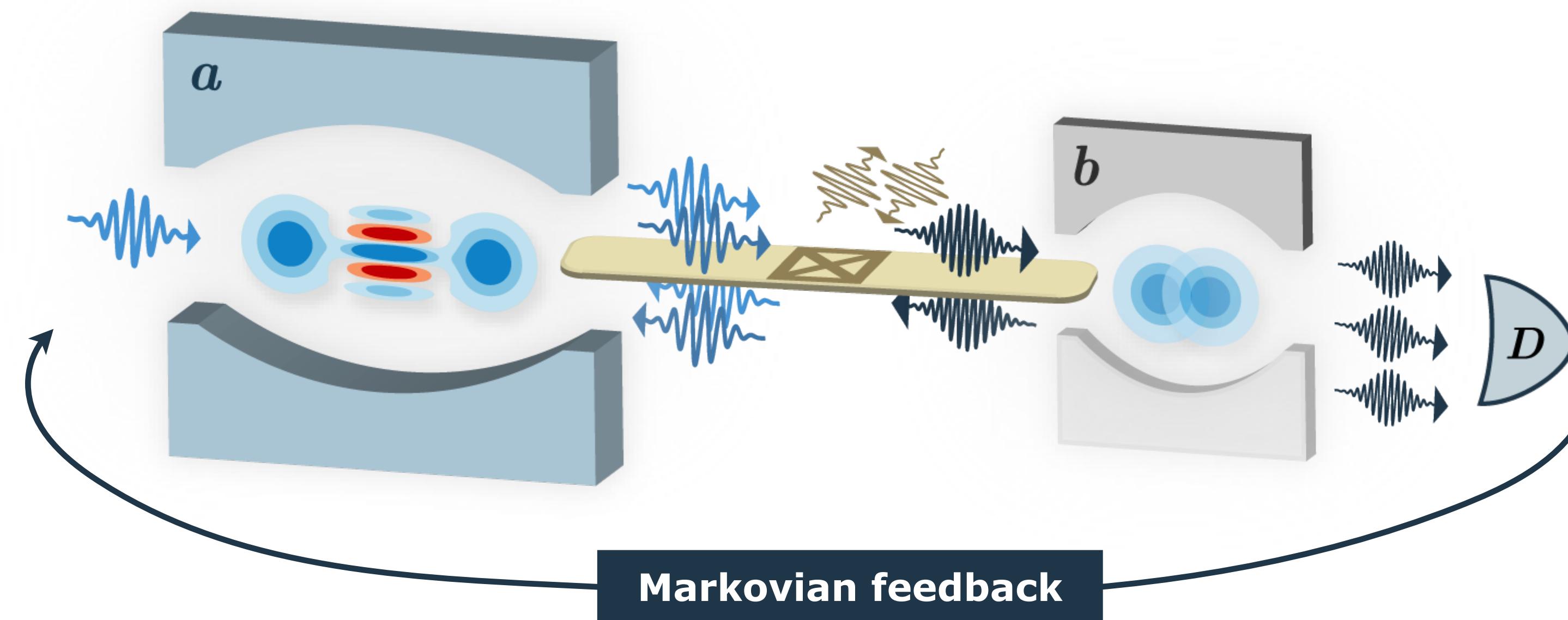
From discovery to answers



From discovery to answers



From discovery to answers



Rabi oscillations with a photodetector

Information is lost through the buffer mode

→ **measure the buffer output** to retrieve it

$$\frac{d\rho}{dt} = -i[H, \rho]dt + \kappa_b (\mathcal{D}_\eta[b]\rho dt + \mathcal{J}[\rho]dN_\eta)$$

no-jump jump

with $0 \leq \eta \leq 1$ (detection efficiency)

Rabi oscillations with a photodetector

Information is lost through the buffer mode

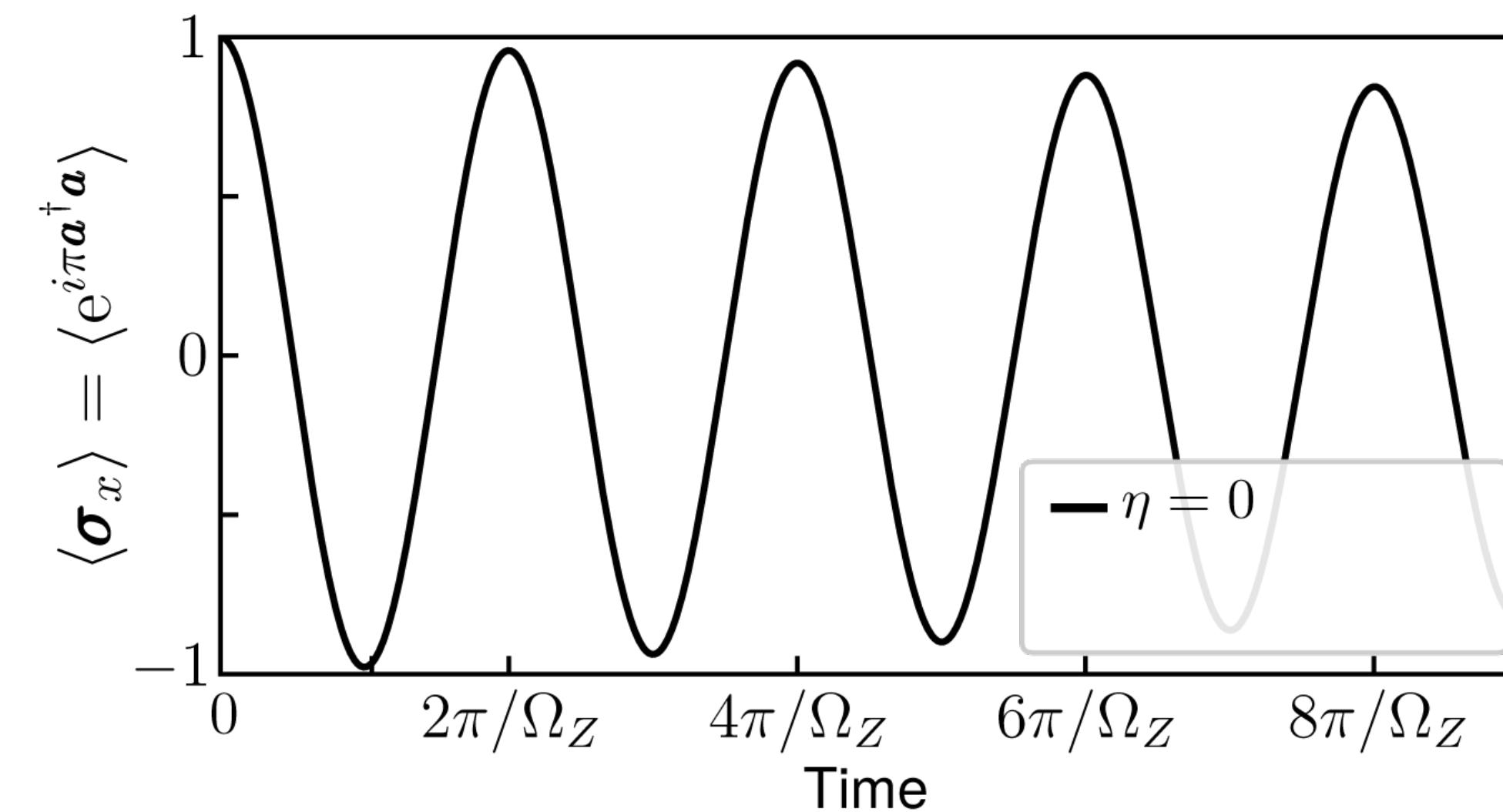
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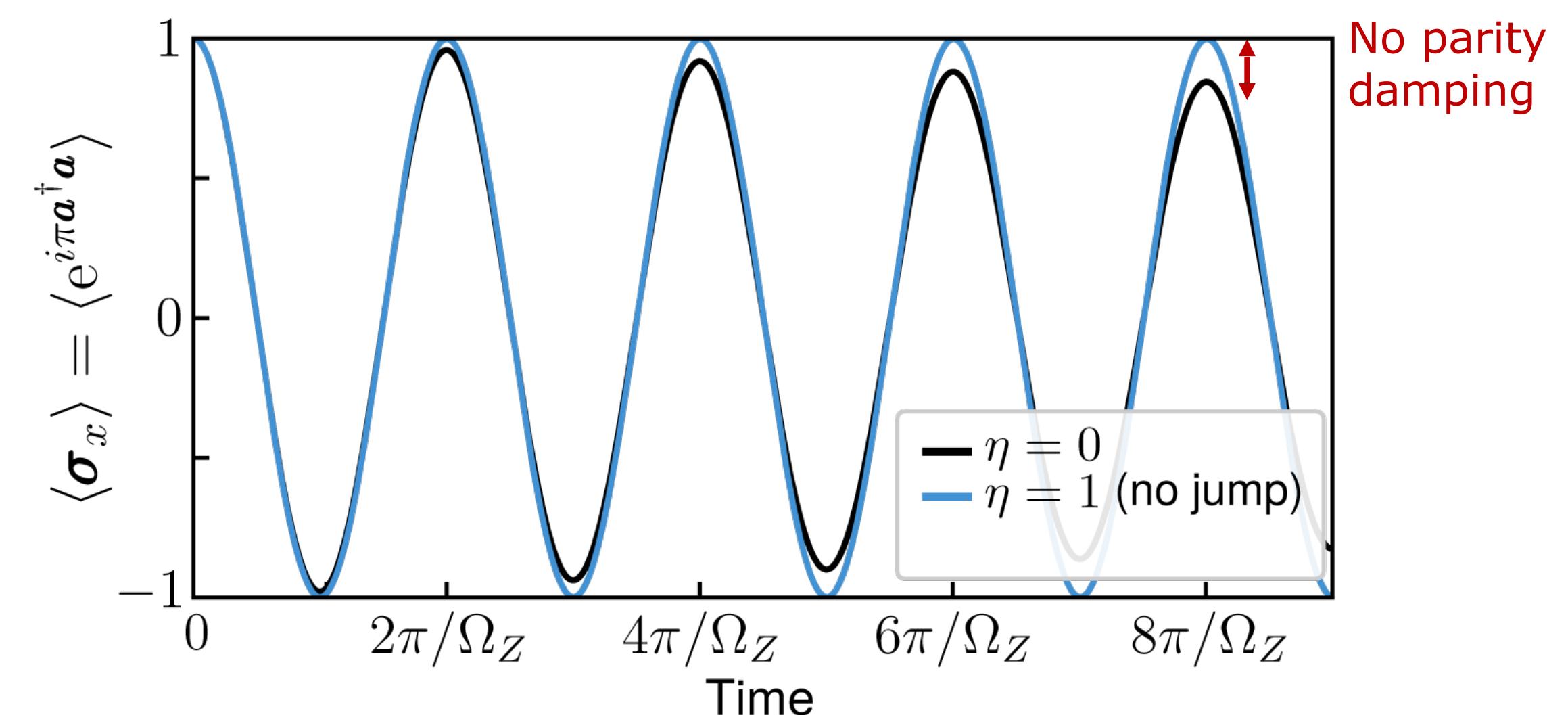
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no-jump jump

with $0 \leq \eta \leq 1$ (detection efficiency)

Assuming $\eta = 1$

► Preserves purity (no information lost)



Rabi oscillations with a photodetector

Information is lost through the buffer mode

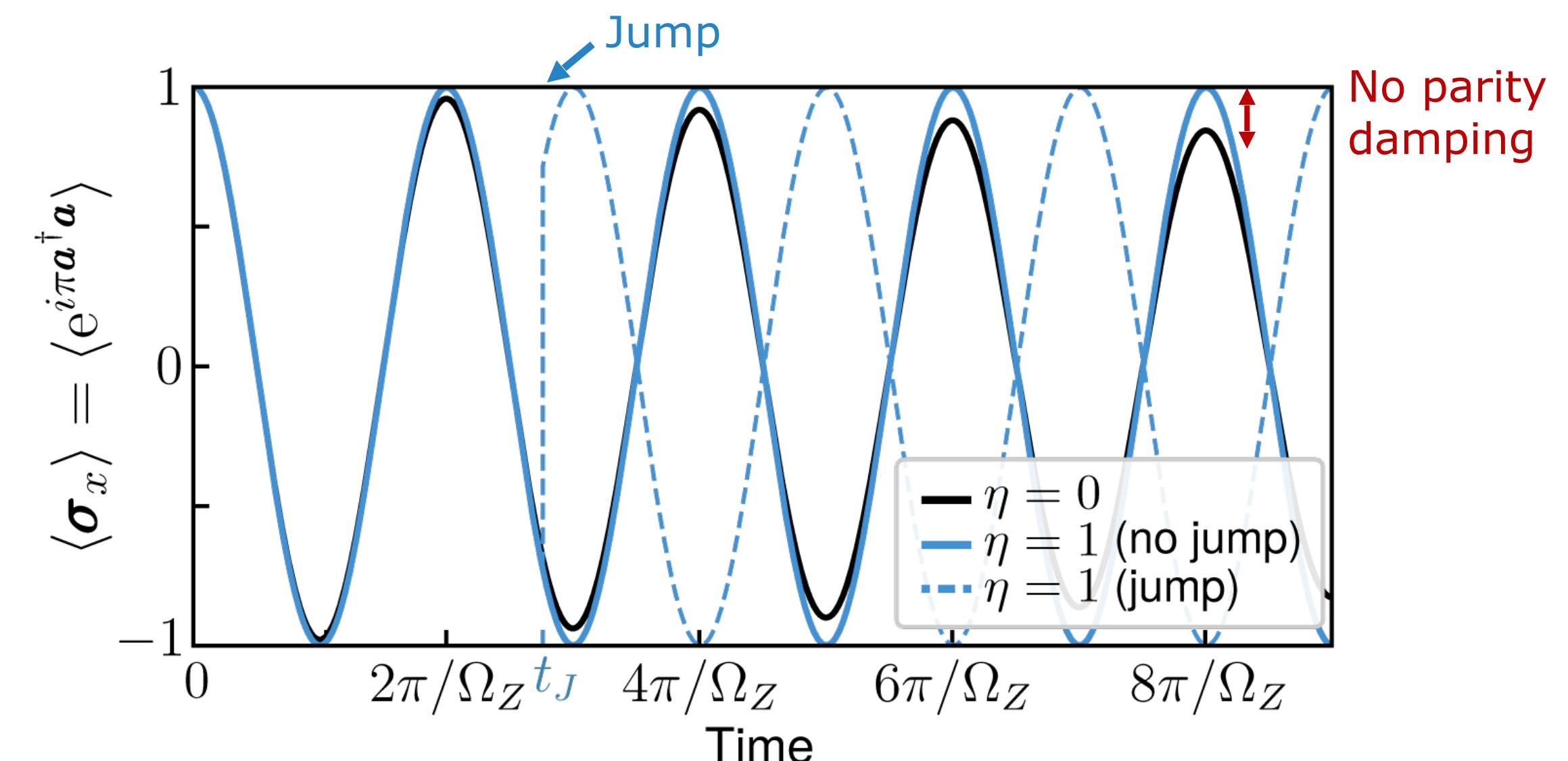
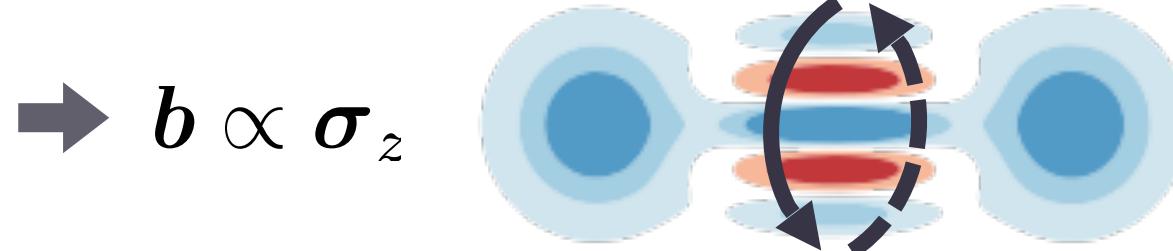
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$$\frac{d\rho}{dt} = -i[H, \rho]dt + \kappa_b (\mathcal{D}_\eta[b]\rho dt + \mathcal{J}[\rho]dN_\eta)$$

no-jump jump
with $0 \leq \eta \leq 1$ (detection efficiency)

Assuming $\eta = 1$

- Preserves purity (no information lost)
- Jump detected = parity swap



Gate engineering with a photodetector

Information is lost through the buffer mode

→ measure the buffer output to retrieve it + **markovian feedback**

$$\frac{d\rho}{dt} = -i[H, \rho]dt + \kappa_b (\mathcal{D}_\eta[b]\rho dt + \mathcal{J}[\rho]dN_\eta)$$

no-jump jump

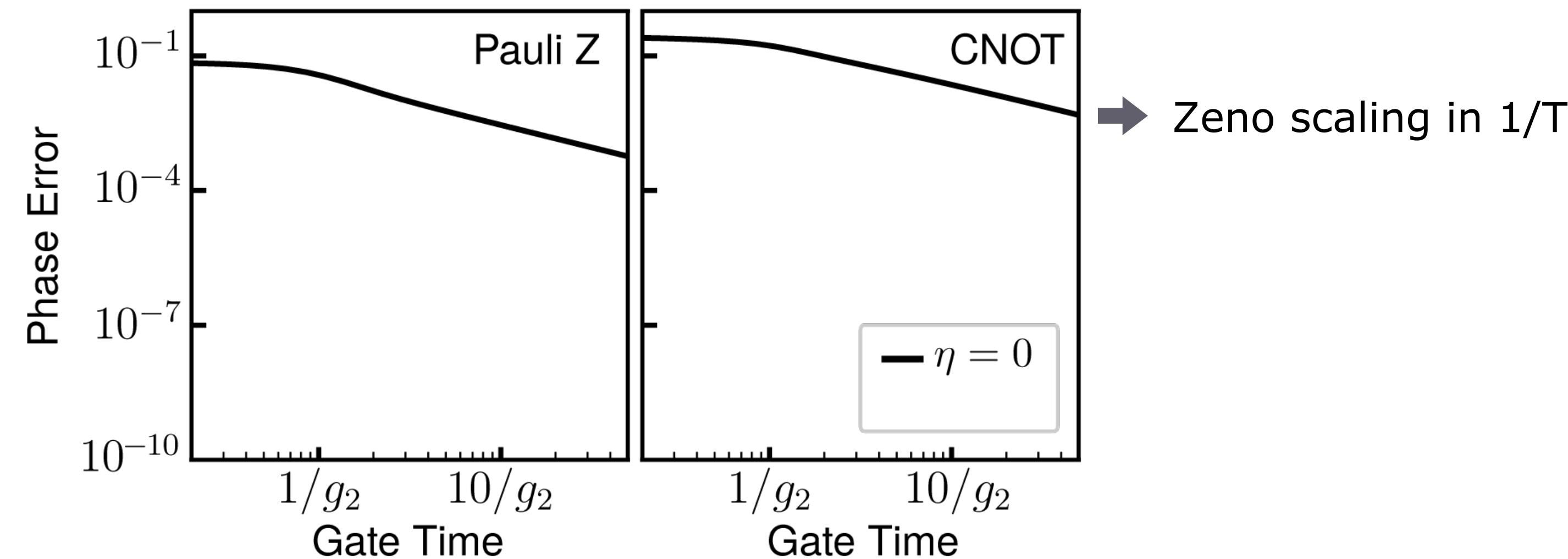
Gate engineering with a photodetector

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no-jump jump



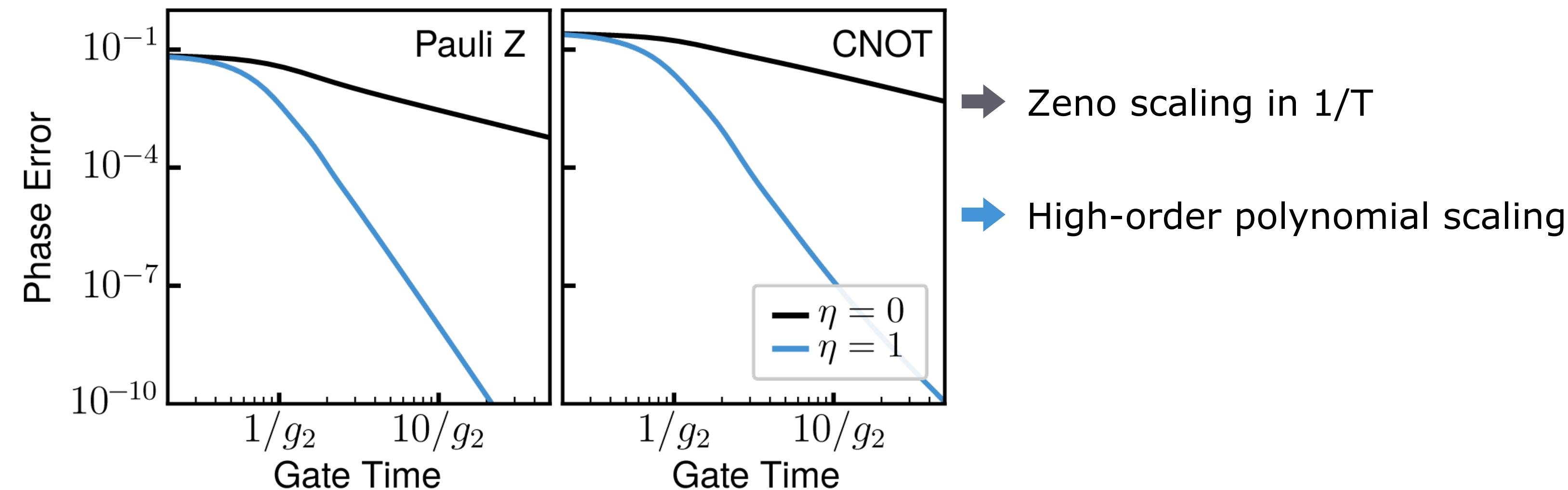
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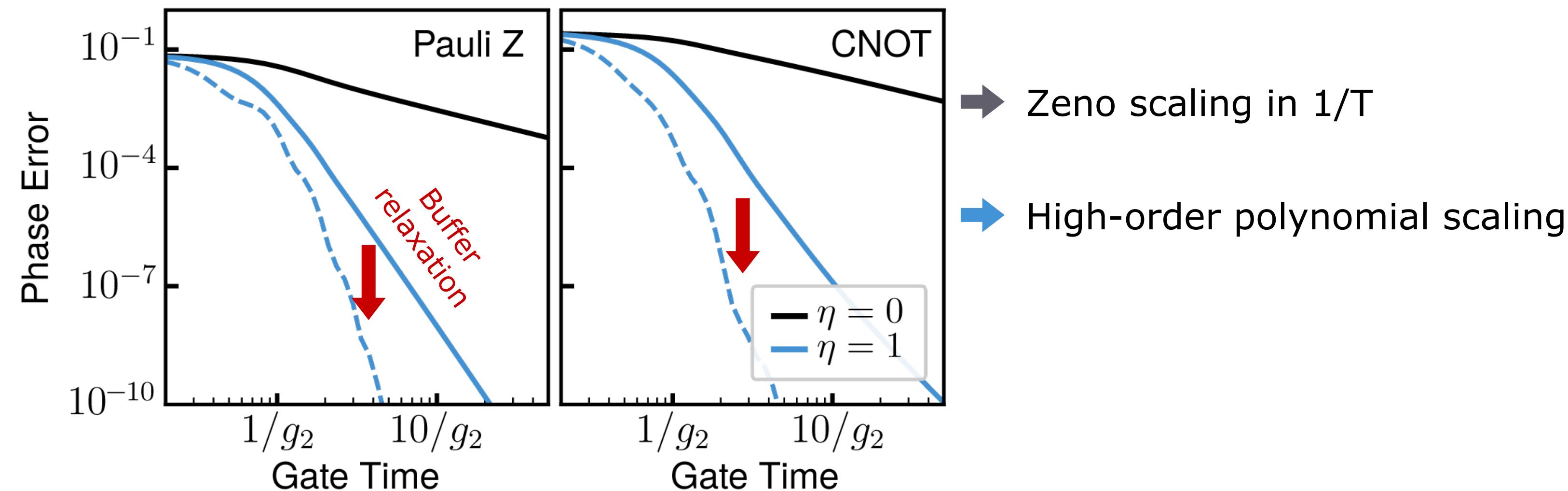
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no-jump jump



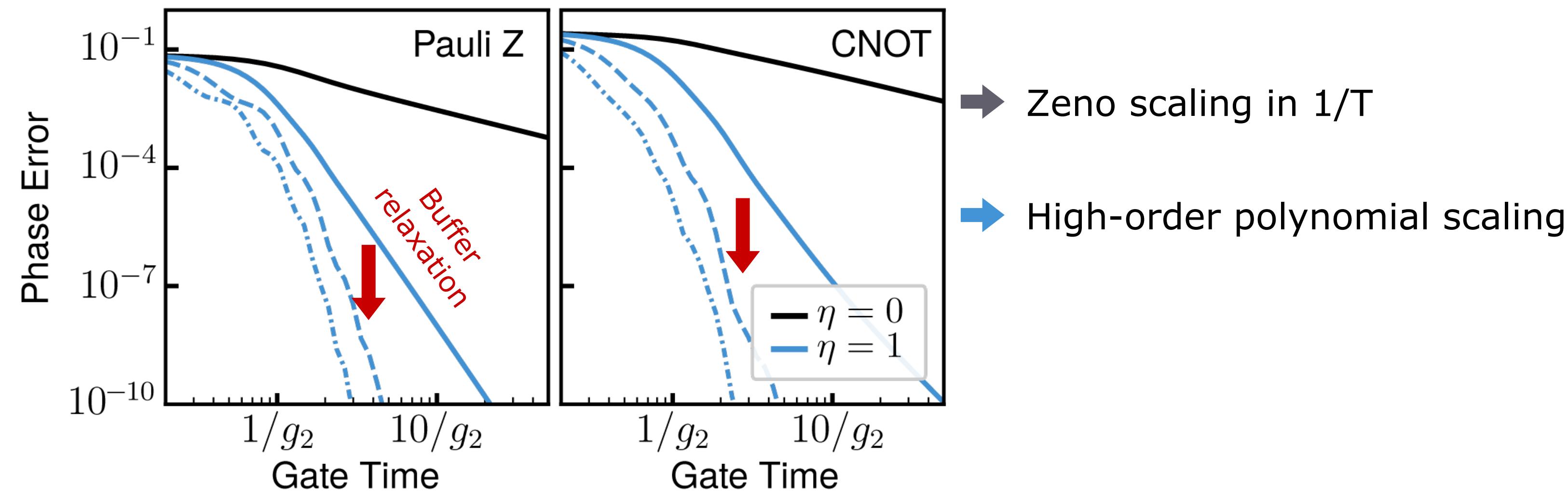
Gate engineering with a photodetector

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no-jump jump



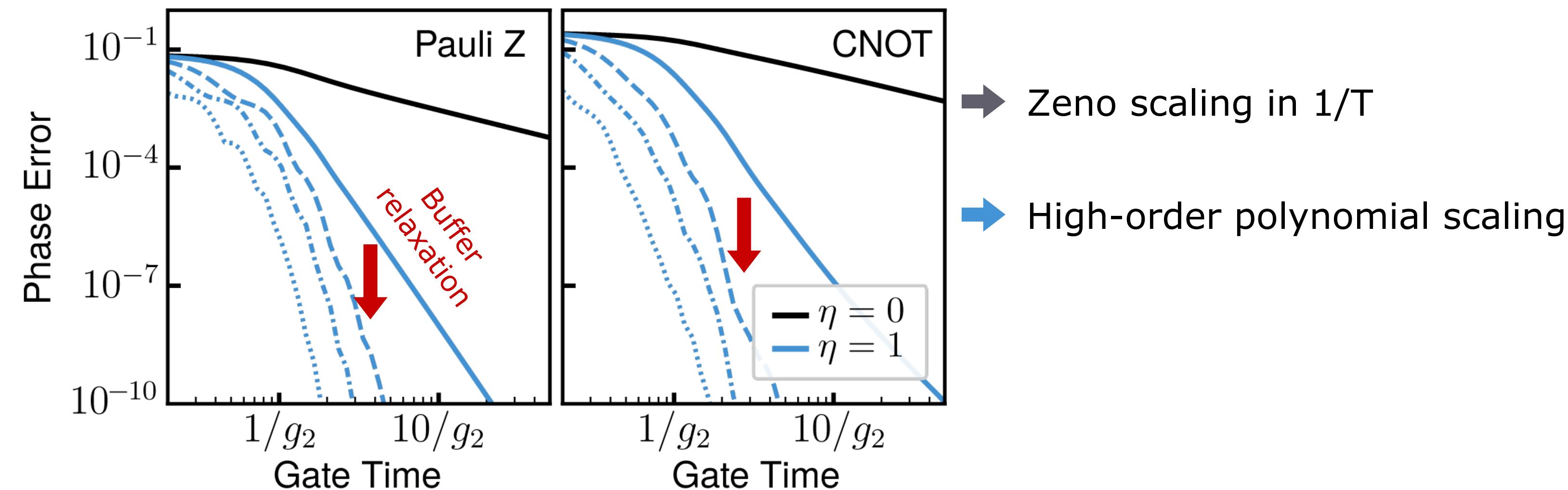
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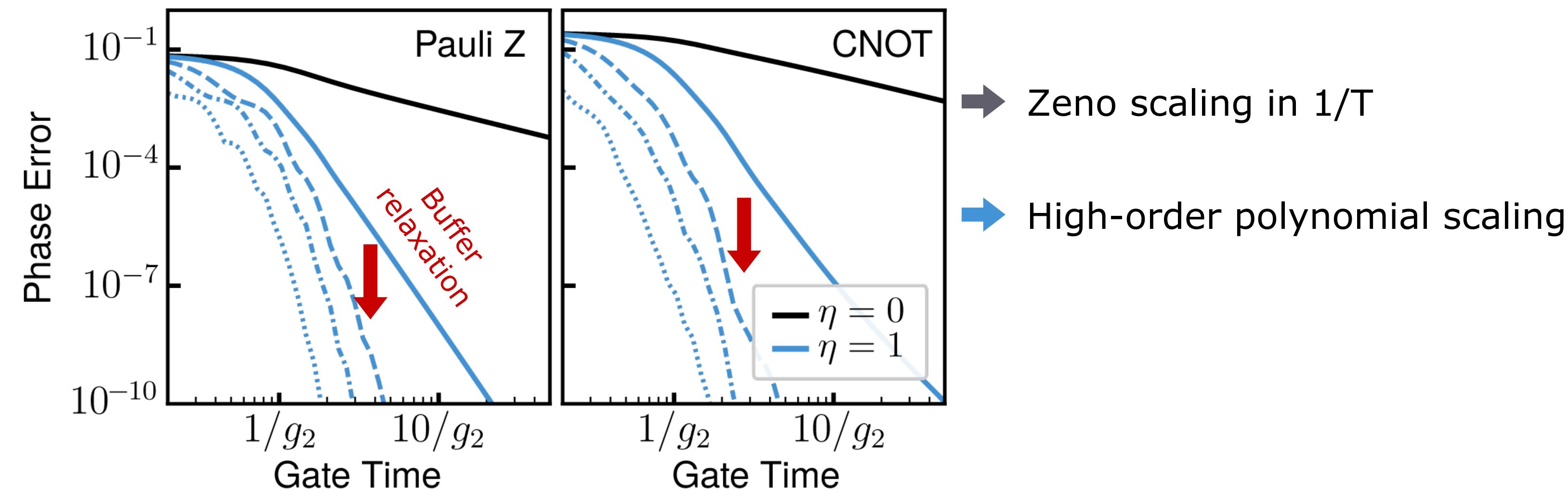
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no-jump jump



- Can apply feedback once per QEC cycle

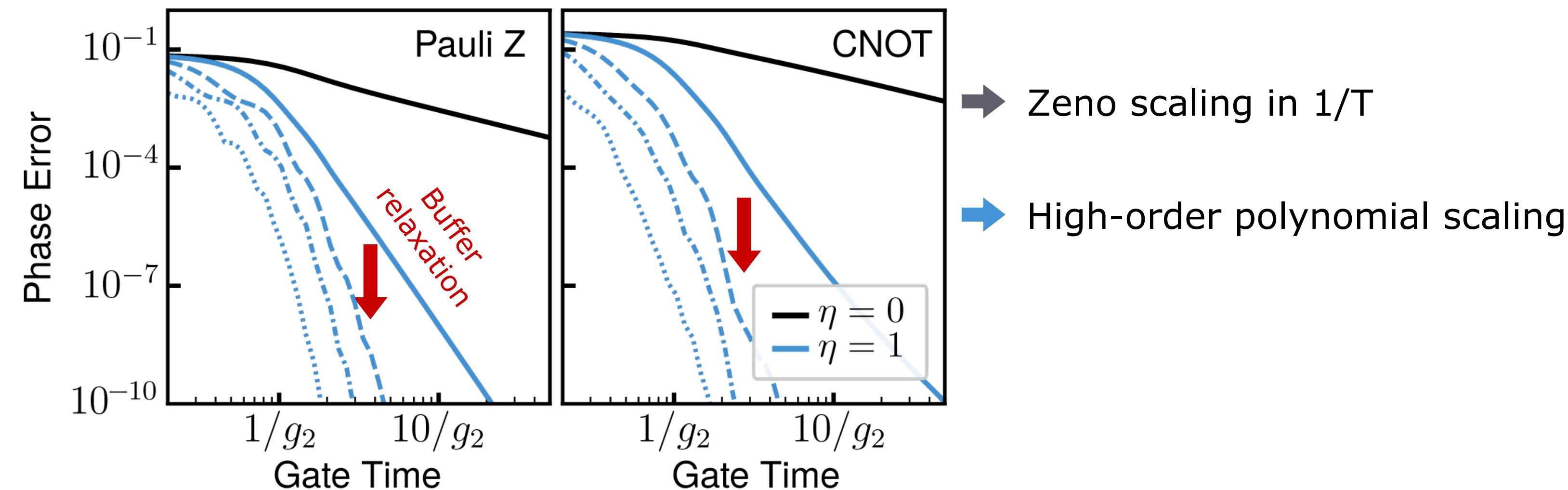
Gate engineering with a photodetector

Information is lost through the buffer mode

→ measure the buffer output to retrieve it + **markovian feedback**

$$\frac{d\rho}{dt} = -i[H, \rho]dt + \kappa_b (\mathcal{D}_\eta[b]\rho dt + \mathcal{J}[\rho]dN_\eta)$$

no-jump jump



- Can apply feedback once per QEC cycle
- Gate fidelity limited by detection efficiency

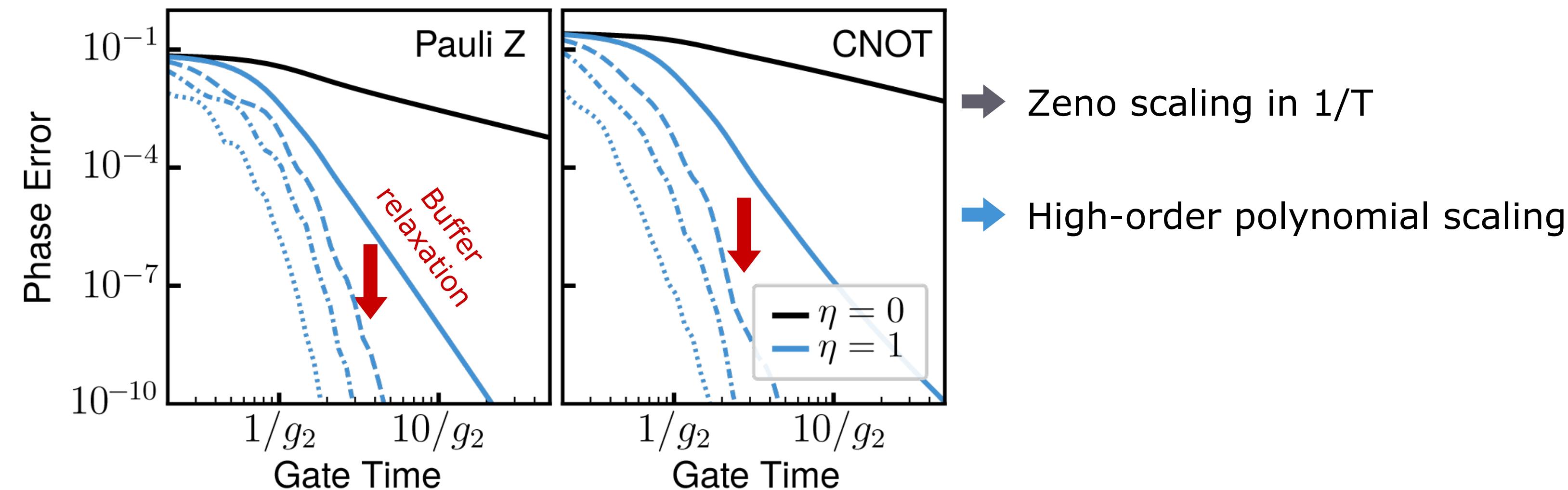
Gate engineering with a photodetector

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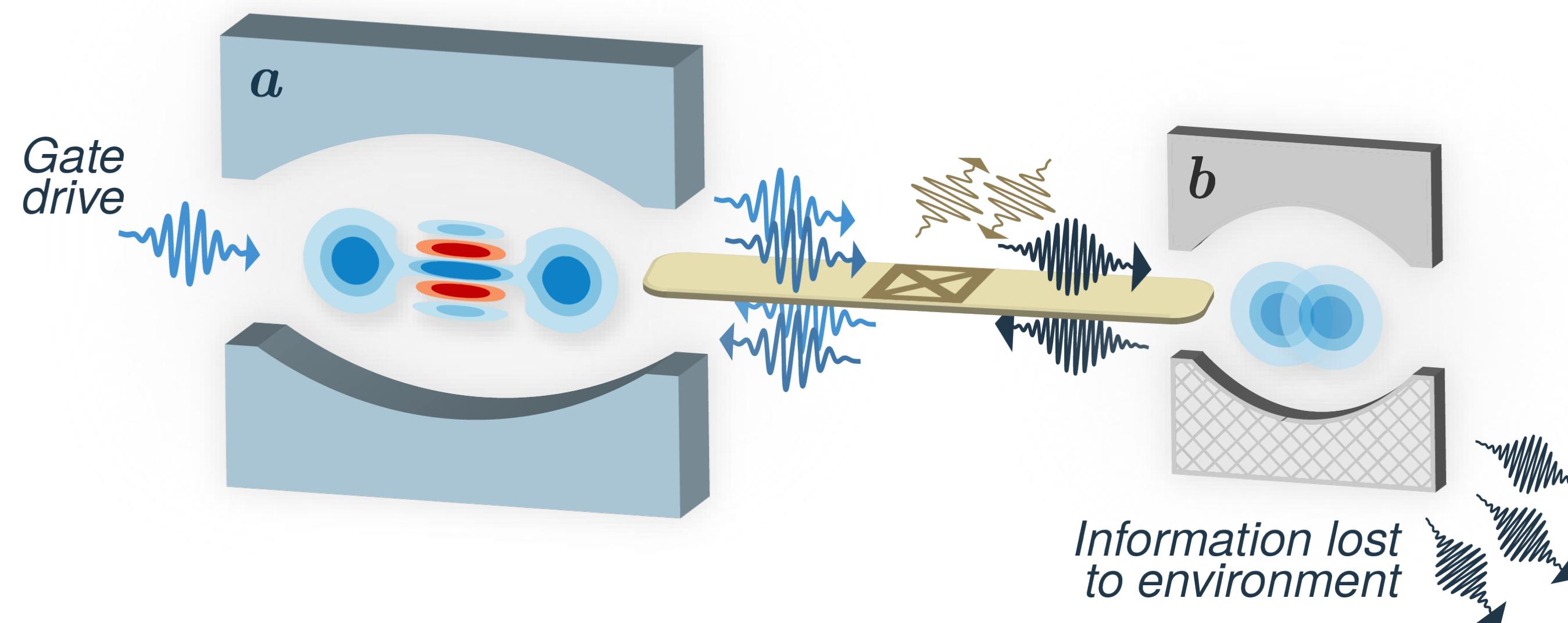
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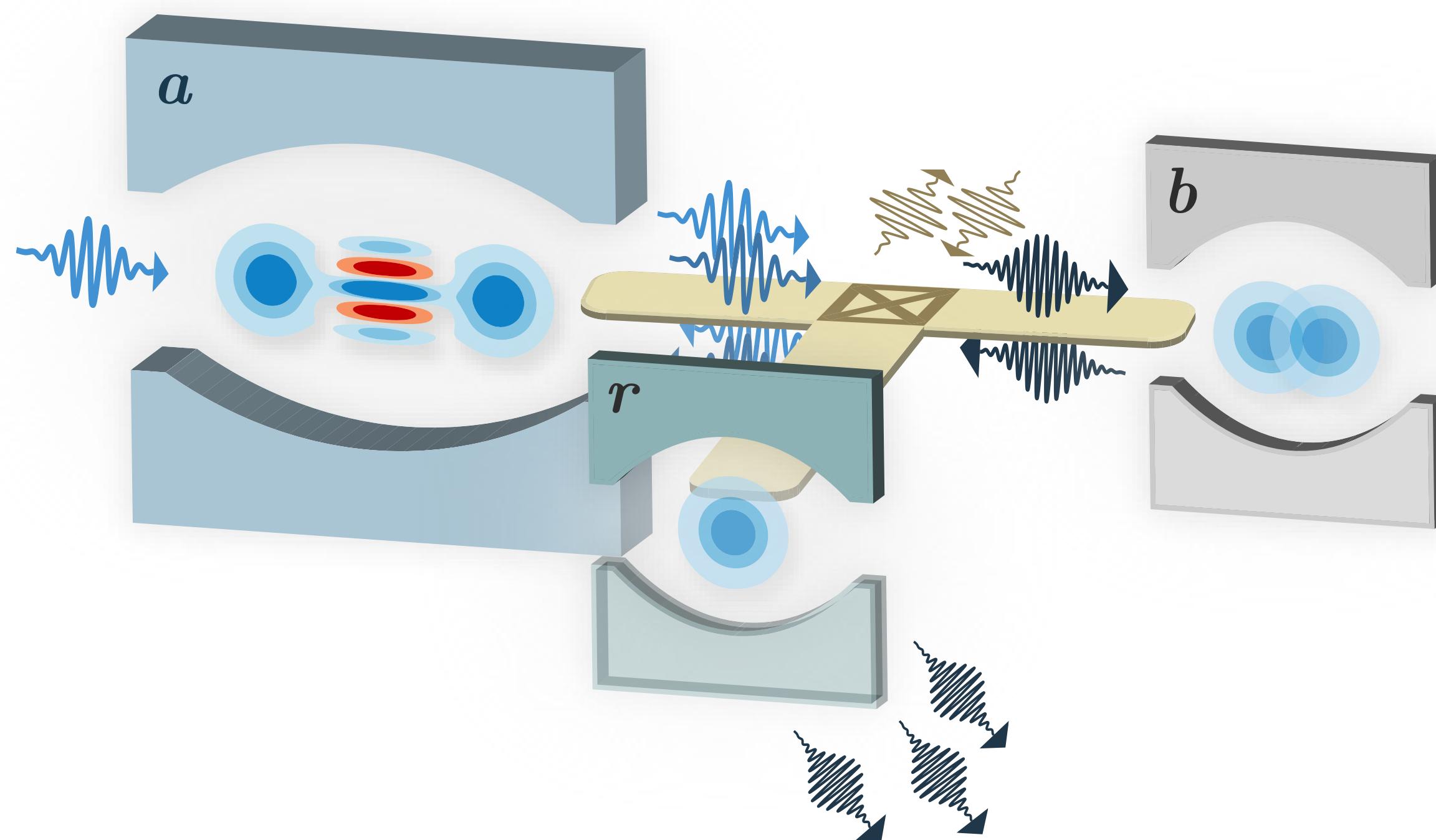


- Can apply feedback once per QEC cycle
- Gate fidelity limited by detection efficiency → **Autonomous feedback**

Autonomous feedback



Autonomous feedback



Autonomous feedback

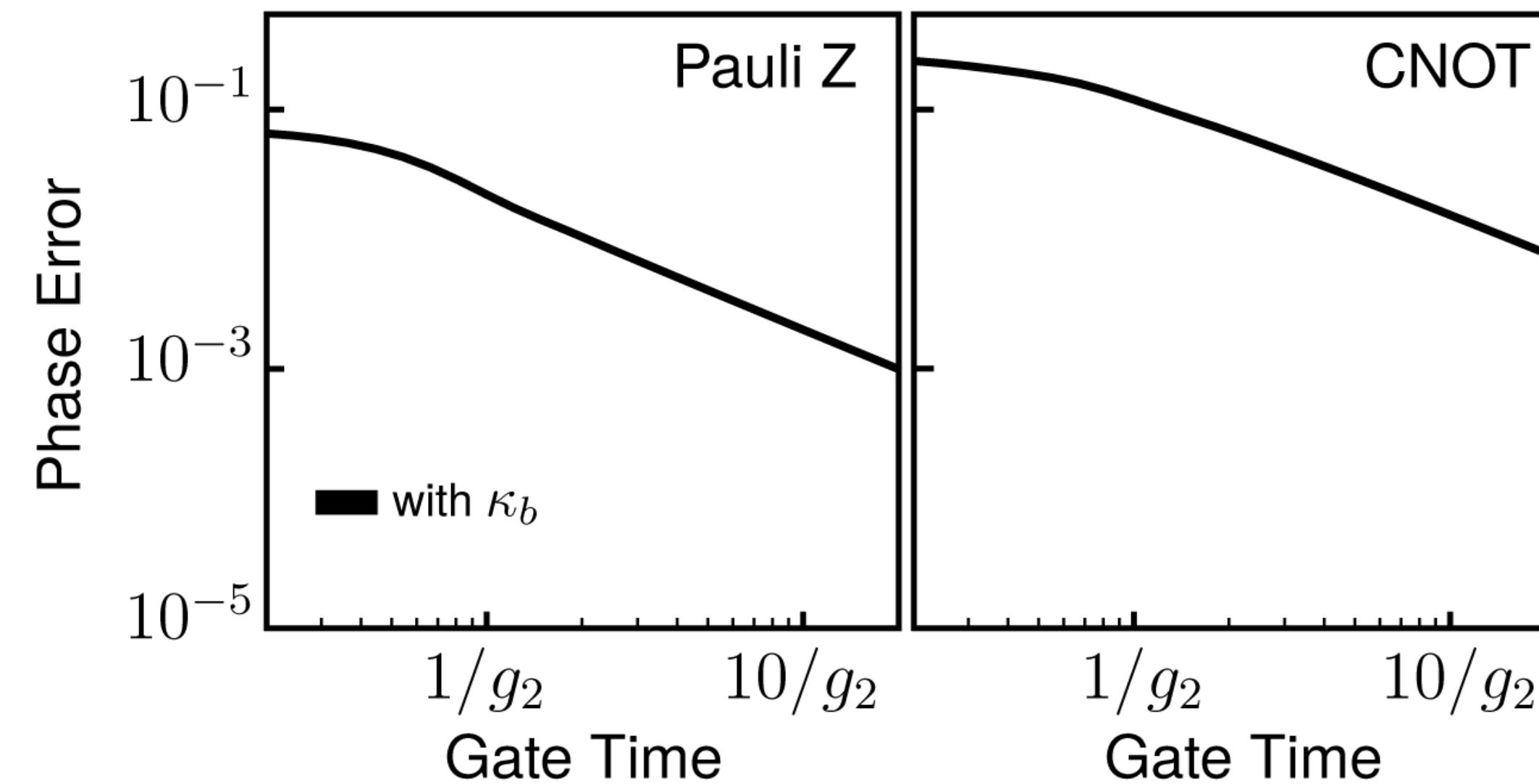
- Correlate buffer photon losses with parity-swaps on memory

$$\frac{d\rho}{dt} = -i[H_{AB} + H_Z, \rho] + \kappa_{ab}\mathcal{D}[ab]\rho$$

Autonomous feedback

- Correlate buffer photon losses with parity-swaps on memory

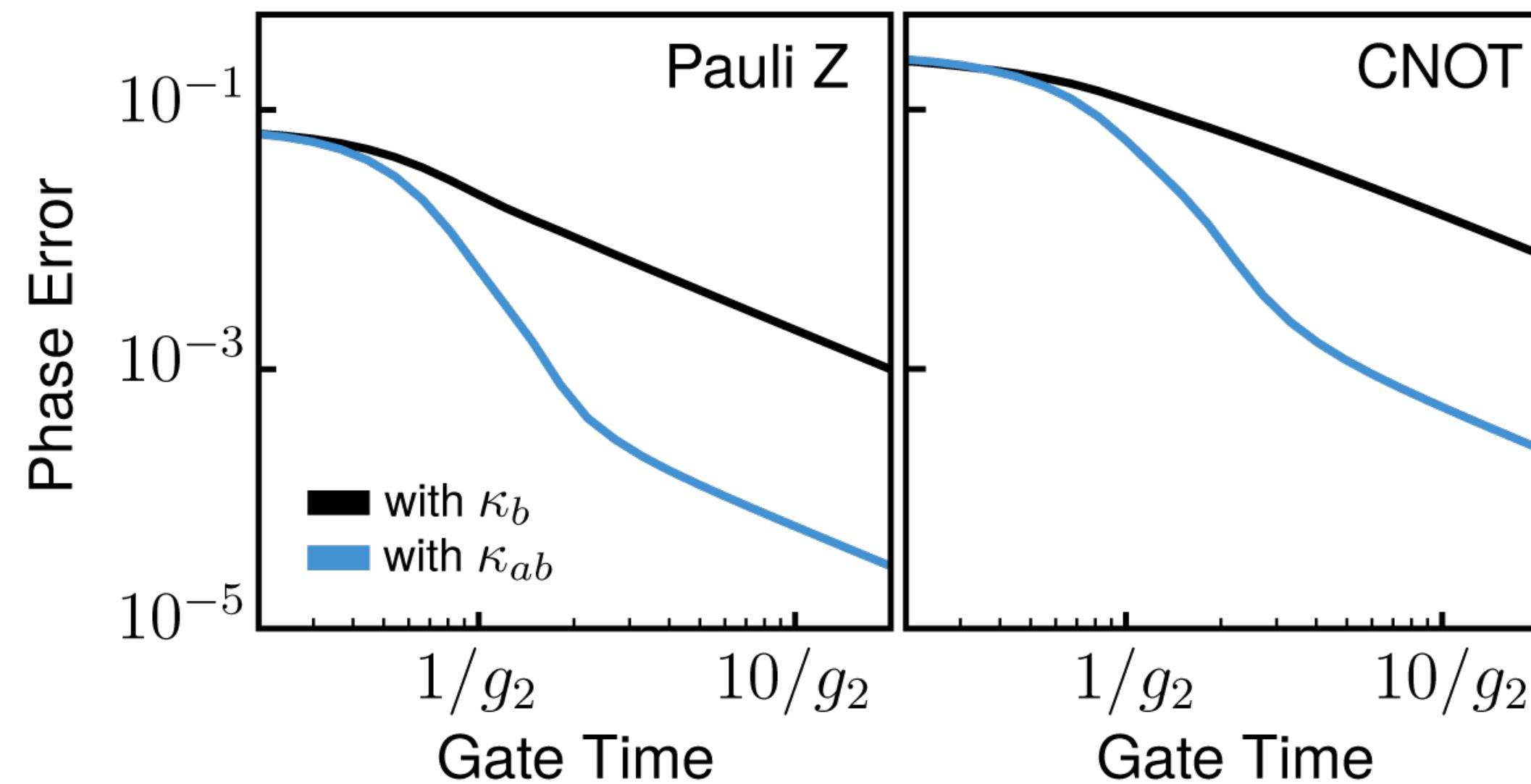
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Autonomous feedback

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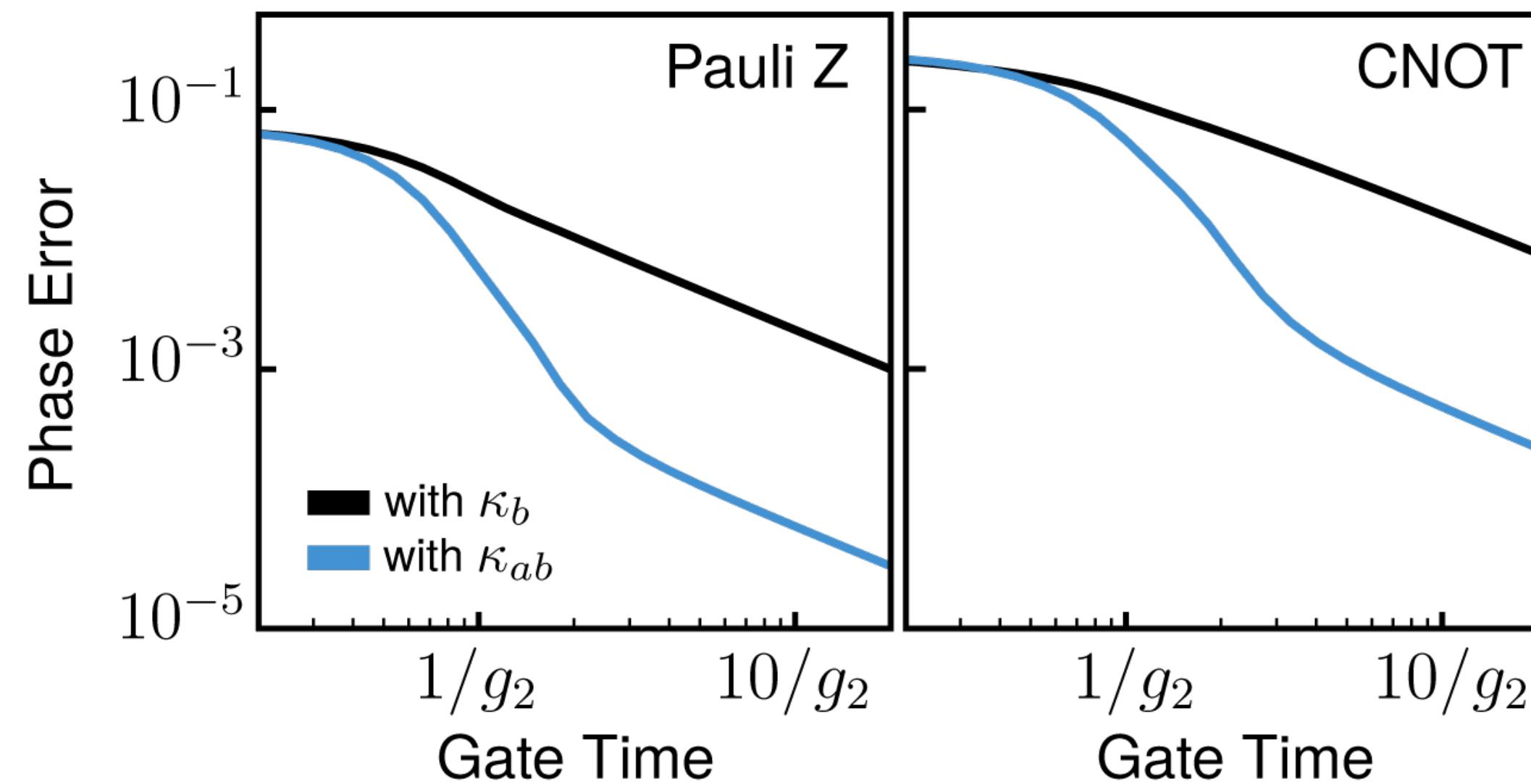


- **x50** gate fidelity improvement (limited by second-order effects), independent of $|\alpha|^2$

Autonomous feedback

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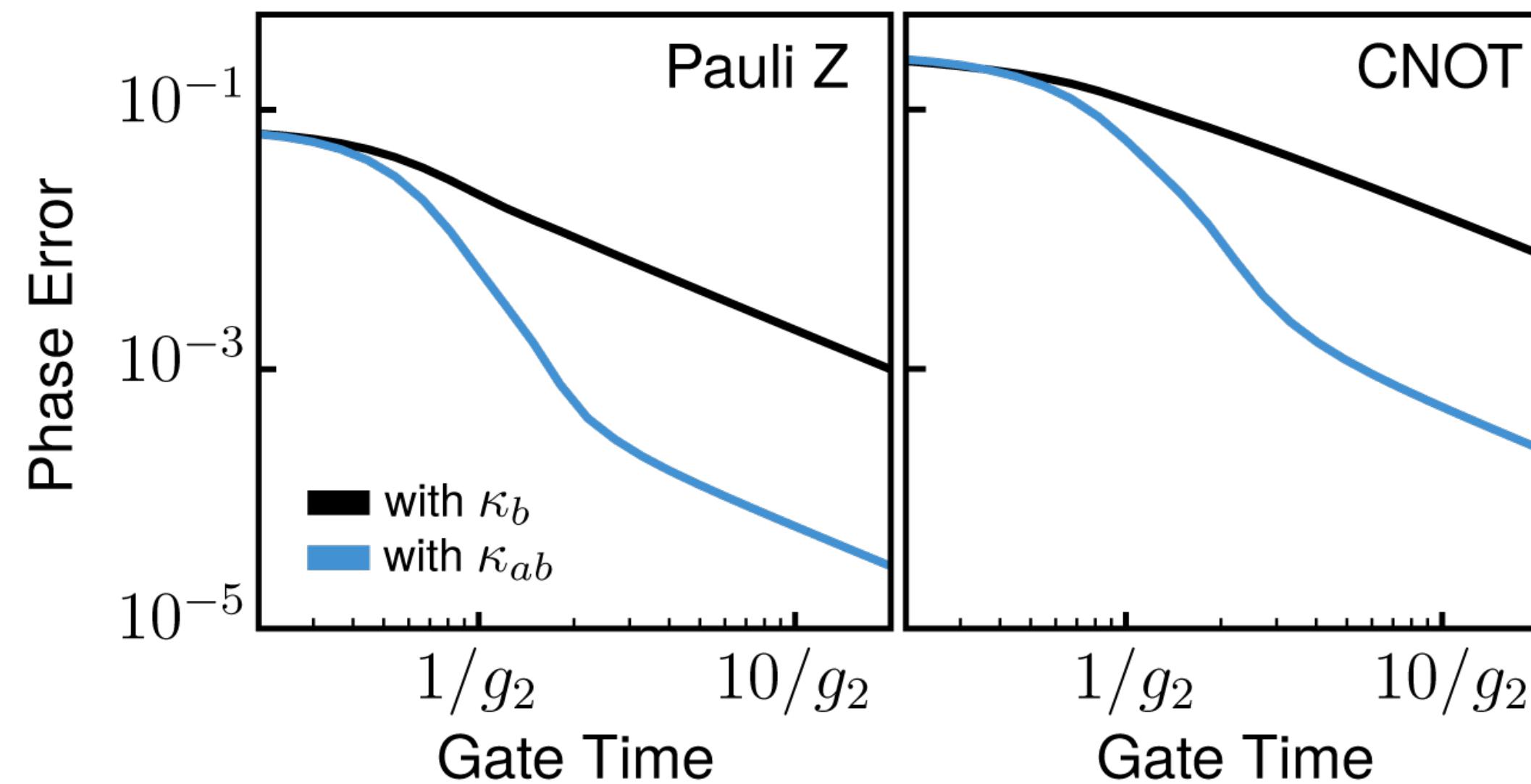


- **x50** gate fidelity improvement (limited by second-order effects), independent of $|\alpha|^2$
- Generalizable to any C^nX gate with no additional experimental overhead

Autonomous feedback

- Correlate buffer photon losses with parity-swaps on memory

$$\frac{d\rho}{dt} = -i[H_{AB} + H_Z, \rho] + \kappa_{ab}\mathcal{D}[ab]\rho$$

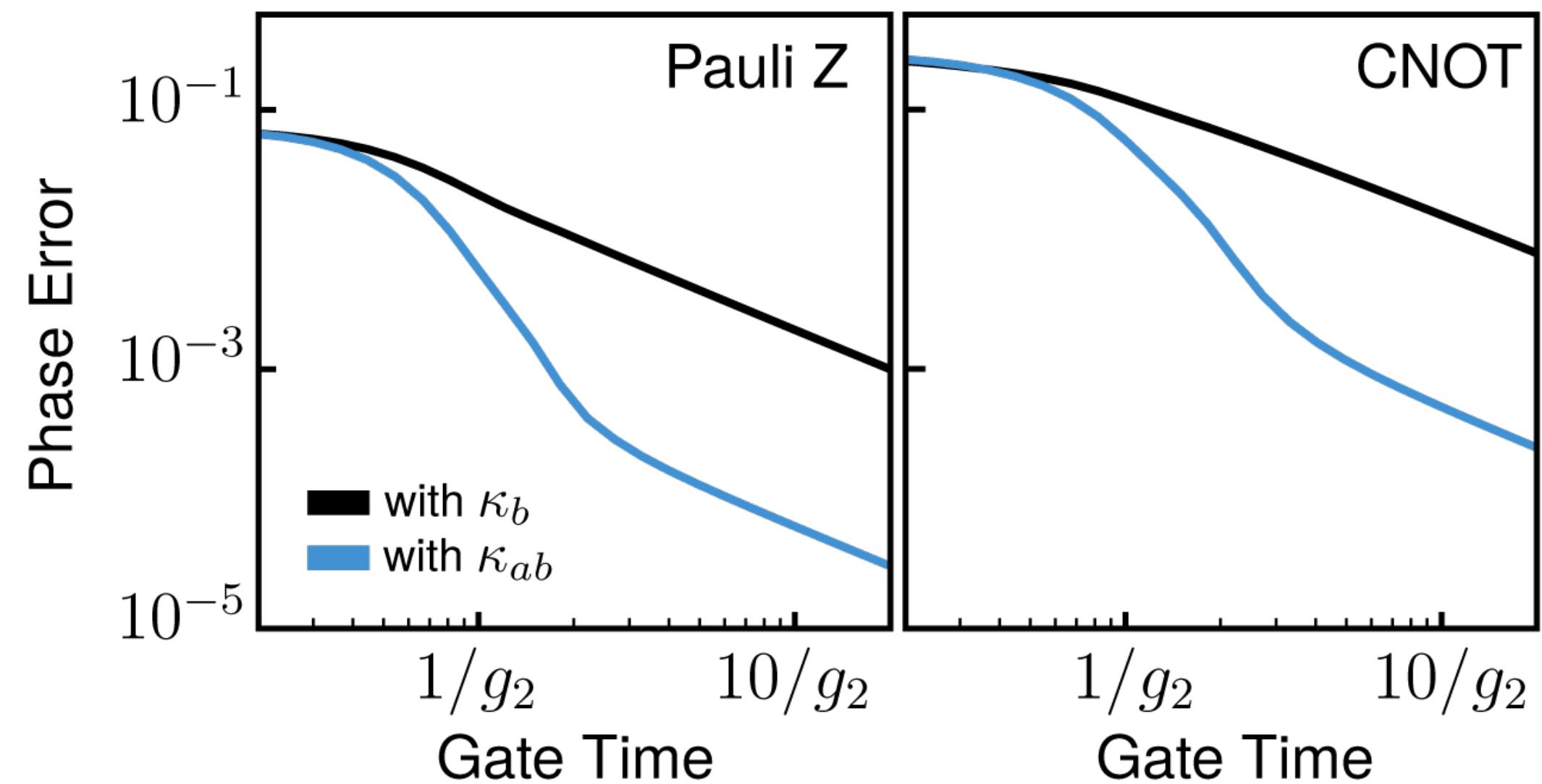


- **x50** gate fidelity improvement (limited by second-order effects), independent of $|\alpha|^2$
- Generalizable to any C^nX gate with no additional experimental overhead
- Autonomous correction of **cavity losses** with squeezed cats

Autonomous feedback

- Correlate buffer photon losses with parity-swaps on memory

$$\frac{d\rho}{dt} = -i[H_{AB} + H_Z, \rho] + \kappa_{ab}\mathcal{D}[ab]\rho$$



- **x50** gate fidelity improvement (limited by second-order effects), independent of $|\alpha|^2$
- Generalizable to any C^nX gate with no additional experimental overhead
- Autonomous correction of **cavity losses** with squeezed cats
- Can tune dissipation in situ

Summary of gate designs

Gautier et al. (2022)

	Hamiltonian \mathbf{H}	Dissipator \mathcal{D}	Gate Errors
Ref. [30, 37] Standard Zeno	$\mathbf{H}_{AB} + \mathbf{H}_Z$	$\kappa_b \mathcal{D}[\mathbf{b}]$	$p_Z^{(0)} \equiv \frac{\pi^2}{16 \alpha ^4 T} \frac{\kappa_b}{4g_2^2}$
Ref. [50]	Combined dissipation and TPE Hamiltonian $\mathbf{H}_{TPE} \equiv g'_2(\mathbf{a}^2 - \alpha^2)\boldsymbol{\sigma}_+ + \text{h.c.}$	$\kappa_b \mathcal{D}[\mathbf{b}]$	$p_Z = \frac{1}{1 + (2g'_2/\kappa_2)^2} p_Z^{(0)}$
Sec. V	Buffer photodetection with classical feedback	$\kappa_b \mathcal{D}[\mathbf{b}]$ (photodetected)	$p_Z \gtrsim (1 - \eta) p_Z^{(0)}$ (detection efficiency η)
Sec. VI	Cat-buffer autonomous feedback	$\kappa_{ab} \mathcal{D}[\mathbf{a}\mathbf{b}]$	$p_Z = \mu p_Z^{(0)}$ with $\mu \gtrsim 0.02$
Sec. VII	Locally flat Hamiltonian $\mathbf{H}_{Z,N} = \varepsilon_Z \sum_{n=0}^N c_n (\mathbf{a} + \mathbf{a}^\dagger)^{2n+1}$	$\kappa_b \mathcal{D}[\mathbf{b}]$	$p_Z = \nu \alpha ^{-2N} p_Z^{(0)}$ with $\nu \sim 1$
Sec. VIII	Discrete jump	$\kappa_b \mathcal{D}[\mathbf{b}] + \kappa_Z \mathcal{D}[\mathbf{a}_\theta \boldsymbol{\sigma}_+]$	$p_Z = \exp(-\kappa_Z \alpha ^2 T)$

High-Fidelity Control and stabilization of Cat Qubits

UDS

Université de
Sherbrooke



ENS



Work on cat qubits

- ▶ RG, A. Sarlette, M. Mirrahimi, *Combined dissipative and Hamiltonian confinement of cat qubits*, PRX Quantum (2021)
- ▶ D. Ruiz, RG, J. Guillaud, M. Mirrahimi, *Two-photon driven Kerr quantum oscillator with multiple spectral degeneracies*, Phys. Rev. A (2022)
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Work on optimal control of open quantum systems

- ▶ RG, É. Genois, A. Blais, *Optimal readout and reset of a transmon*, in preparation
- ▶ P. Guilmin, RG, A. Bocquet, É. Genois, *dynamiqs: an open-source library for GPU-accelerated and differentiable quantum simulation*, in preparation

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cnrs

Inria



Thank you to all colleagues @ Inria, ENS, Alice & Bob, and Institut Quantique

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Supplementary slides

Quantum Optimal Control

QOC Finding an optimal set of parameters for a given operation

► Gradient-free

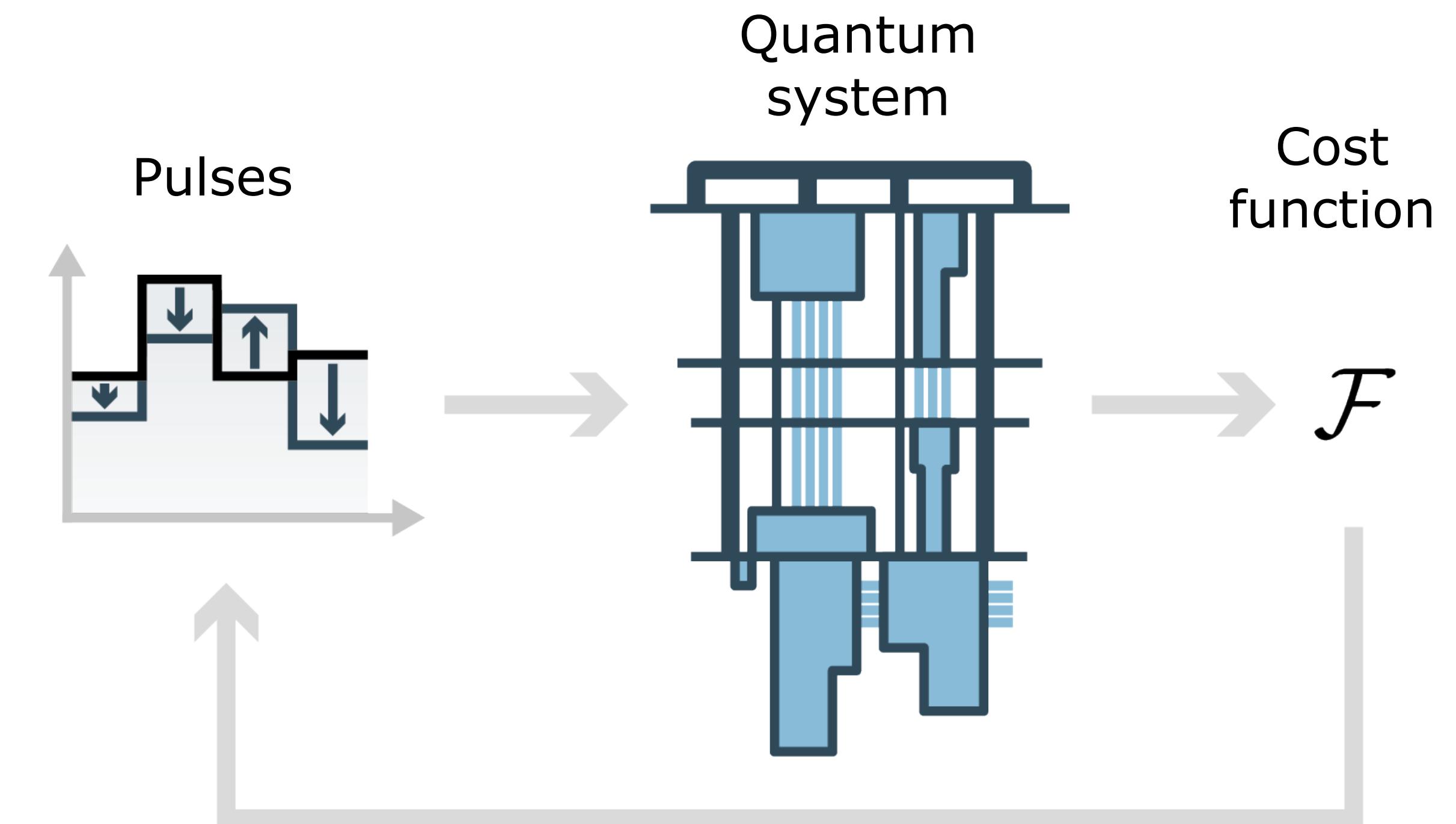
- DRAG
- Chopped Random Basis (CRAB)
- Nelder-Mead
- STIRAP
- Reinforcement Learning

► Gradient-based

- Krotov
- GRAPE
- Automatic Differentiation
- Adjoint state



Any closed or open system, any cost function, low memory overhead, fast



Optimal Control with Automatic Differentiation

Objective Minimize a cost function

$$C = C(\theta, \hat{\rho}(t_0), \dots, \hat{\rho}(t_n))$$

▶ Master Equation

$$\frac{d\hat{\rho}}{dt} = \mathcal{L}\hat{\rho} \equiv -i[\hat{H}, \hat{\rho}] + \sum_k \mathcal{D}[\hat{L}_k]\hat{\rho}$$

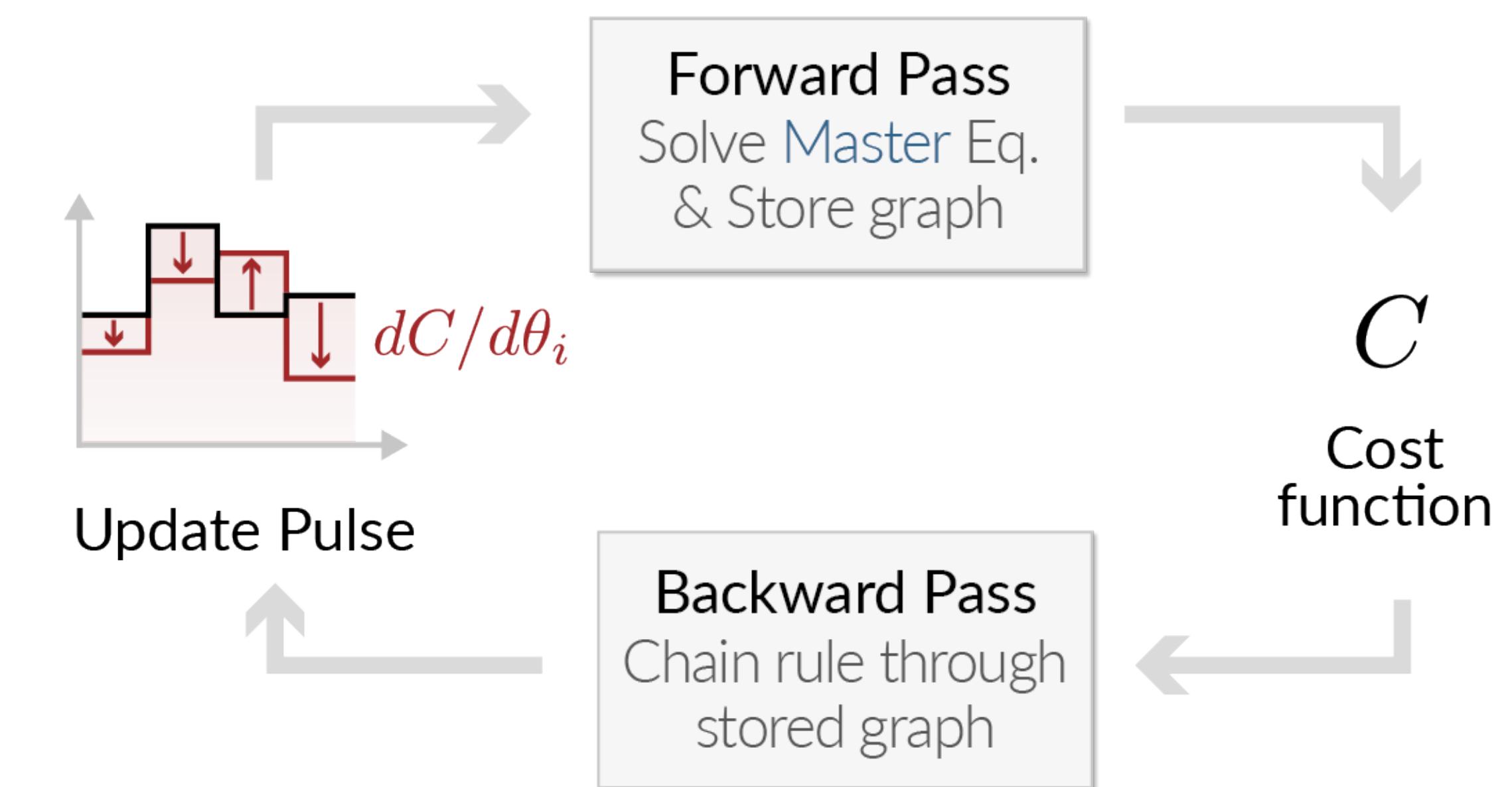
▶ Gradients

$$\frac{dC}{d\theta_i} = \frac{\partial C}{\partial \theta_i} + \sum_k \frac{\partial C}{\partial \hat{\rho}(t_k)} \frac{d\hat{\rho}(t_k)}{d\theta_i}$$

- Store graph of operations
- Backward through graph using chain rule



Memory overhead in $\mathcal{O}(N_t \times N^2)$



Transmon
readout



$$\frac{N_t}{N} \sim 10^4$$



$$\sim 1,16 \text{ TB}$$

Adjoint State Quantum Optimal Control

Objective Minimize a cost function

$$C = C(\theta, \hat{\rho}(t_0), \dots, \hat{\rho}(t_n))$$

▶ Master Equation

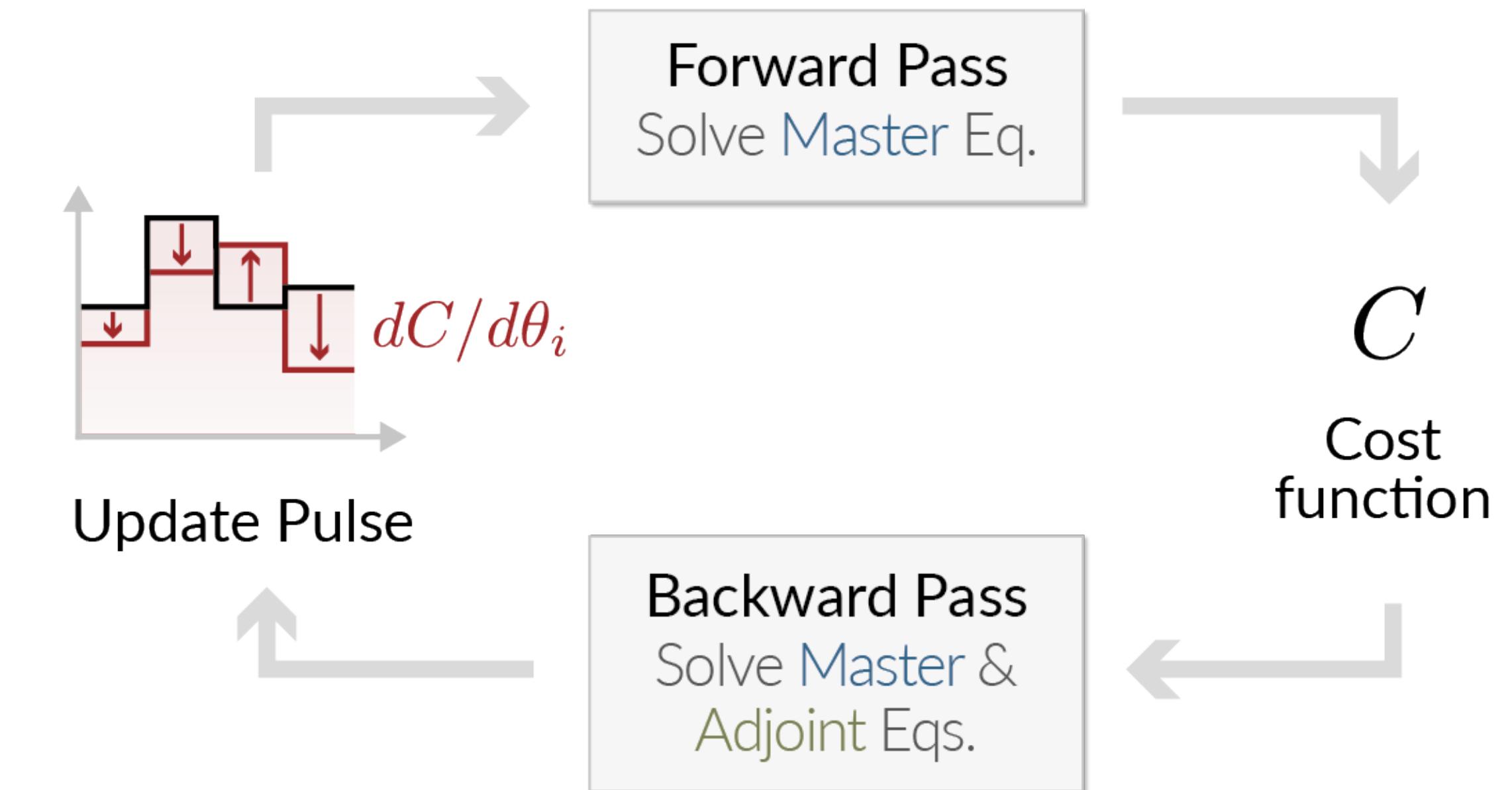
$$\frac{d\hat{\rho}}{dt} = \mathcal{L}\hat{\rho} \equiv -i[\hat{H}, \hat{\rho}] + \sum_k \mathcal{D}[\hat{L}_k]\hat{\rho}$$

▶ Adjoint state $\hat{\phi}(t) = dC/d\hat{\rho}(t)$

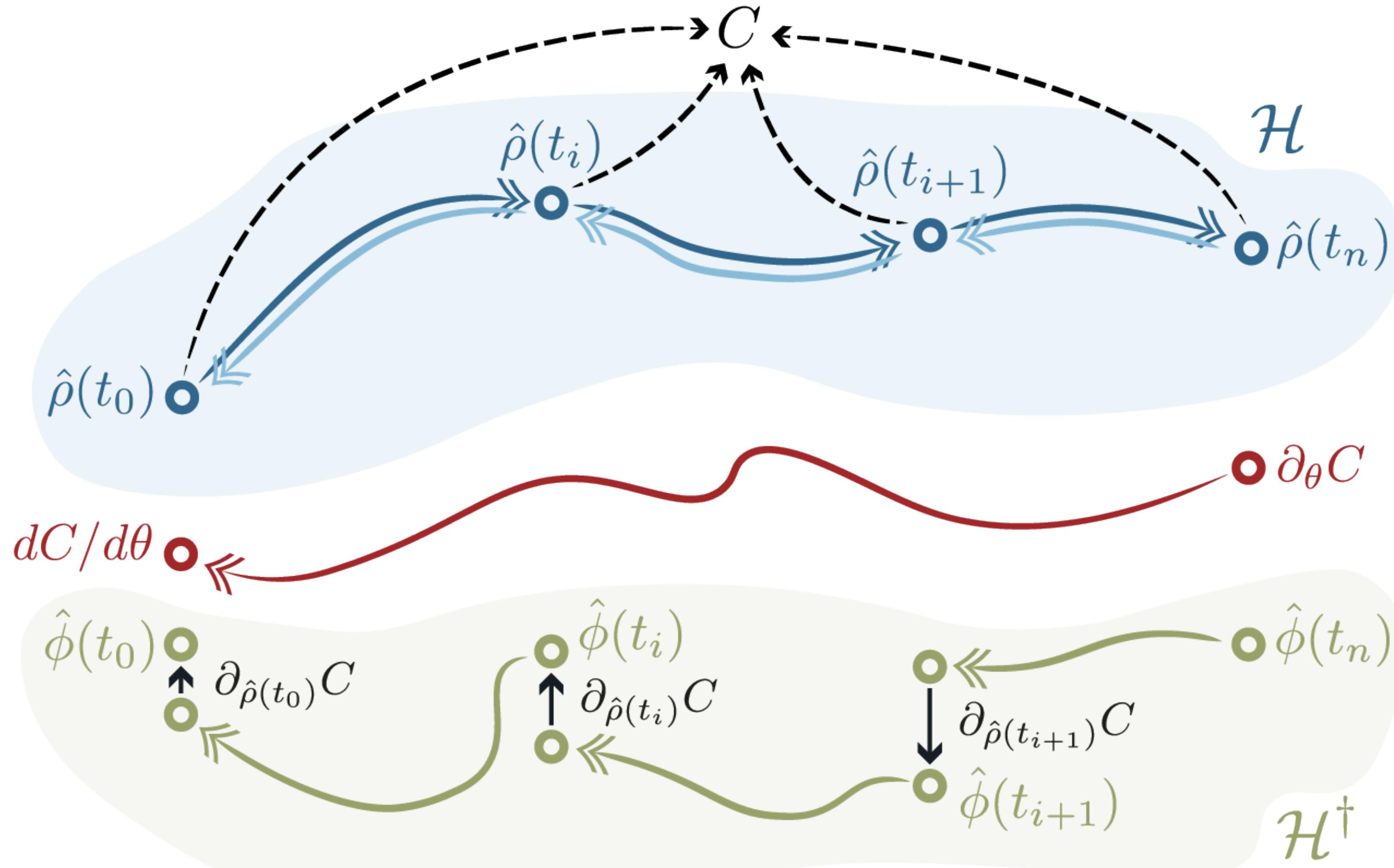
$$\frac{d\hat{\phi}}{dt} = -\mathcal{L}^\dagger \hat{\phi} \equiv -i[\hat{H}, \hat{\phi}] - \sum_k \mathcal{D}^\dagger[\hat{L}_k]\hat{\phi}$$

▶ Gradients

$$\frac{dC}{d\theta} = \frac{\partial C}{\partial \theta} - \int_{t_n}^{t_0} \partial_\theta \text{Tr} \left[\hat{\phi}^\dagger(t) \mathcal{L}(t, \theta) \hat{\rho}(t) \right] dt$$



Adjoint State Quantum Optimal Control



Adjoint State Quantum Optimal Control

Objective Minimize a cost function

$$C = C(\theta, \hat{\rho}(t_0), \dots, \hat{\rho}(t_n))$$

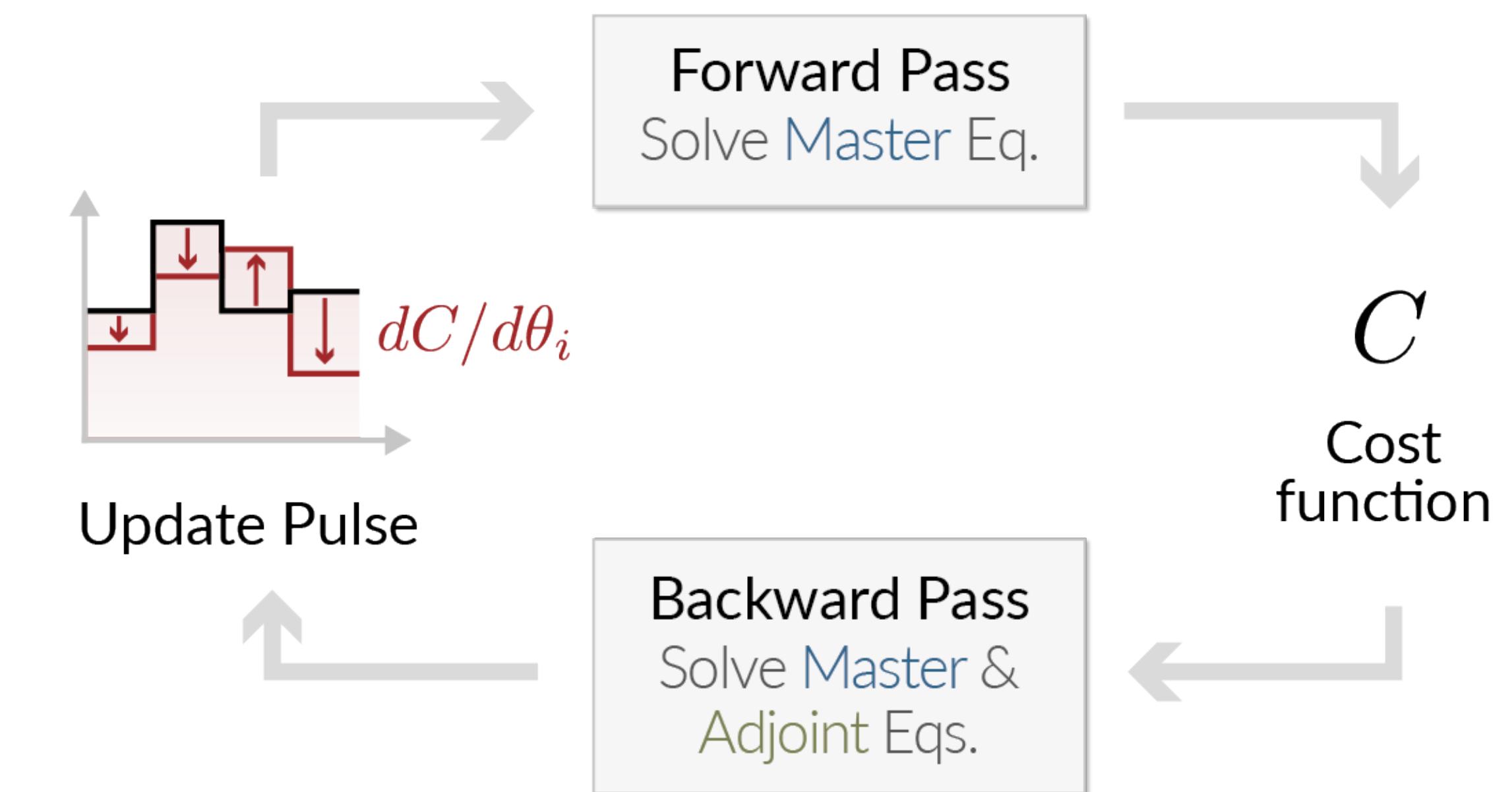
▶ Gradient descent

- Adam
- Stochastic Gradient Descent
- L-FBGS

▶ Space-time costs

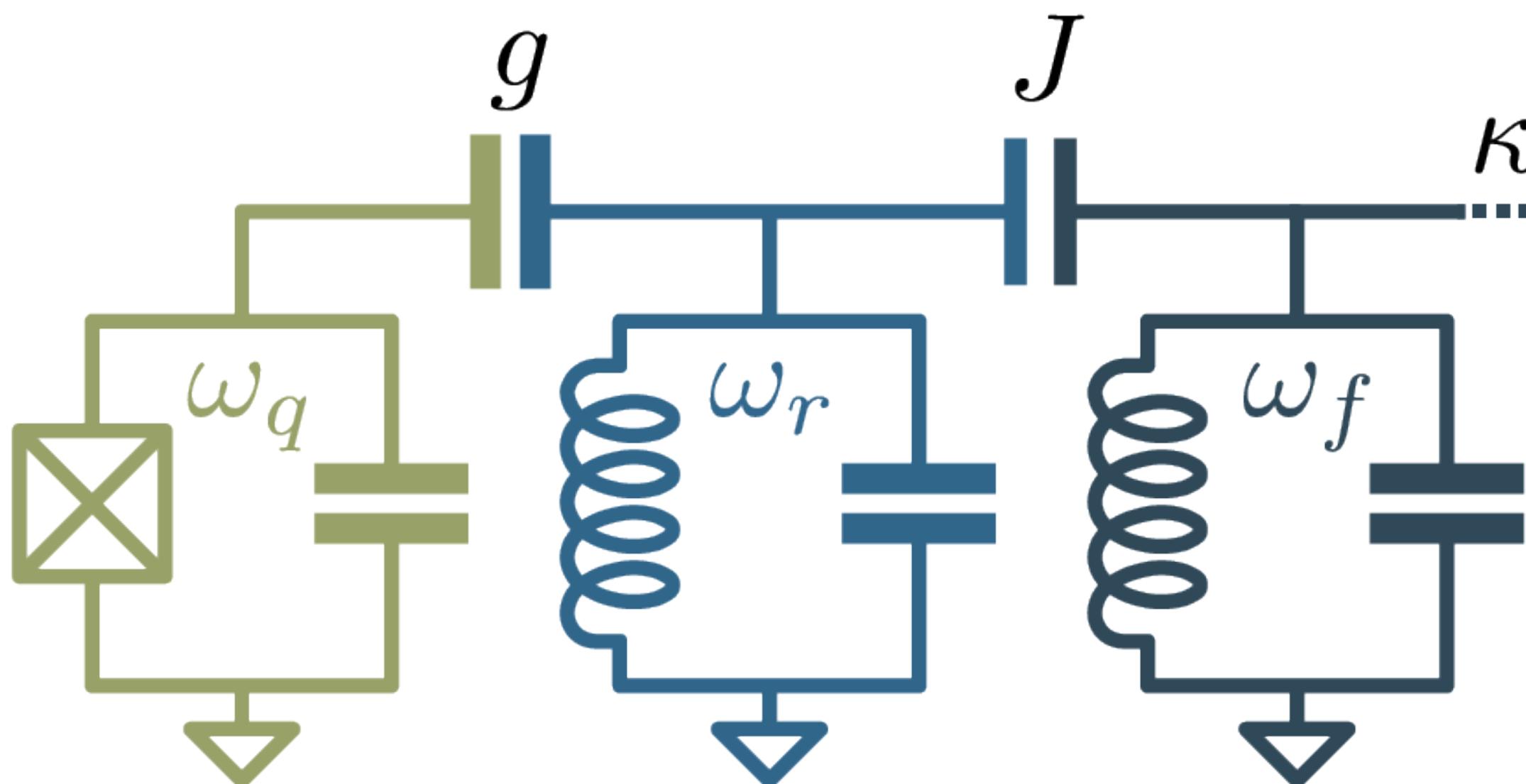
- Memory cost $\mathcal{O}(N_{cp} \times N^2)$
- Time cost $\mathcal{O}(4T_{ME})$
- Can trade memory for numerical stability

▶ Demonstration on transmon readout & reset



Transmon-resonator-filter model

$$\hat{H} = 4E_C\hat{n}_t^2 - E_J \cos(\hat{\varphi}_t) + \omega_r \hat{a}^\dagger \hat{a} + \omega_f \hat{f}^\dagger \hat{f} + i g \hat{n}_t (\hat{a}^\dagger - \hat{a}) + J (\hat{f}^\dagger \hat{a} + \hat{a}^\dagger \hat{f})$$



$E_J/2\pi = 16$ GHz
$E_c/2\pi = 315$ MHz
$E_J/E_c \approx 51$
$\omega_t/2\pi = 6$ GHz
$\omega_r/2\pi = 7.2$ GHz
$\omega_p/2\pi = 7.21$ GHz
$g/2\pi = 150$ MHz
$J/2\pi = 30$ MHz
$\kappa_p/2\pi = 30$ MHz
$\kappa_q/2\pi = 8$ KHz
$\bar{n}_{\text{crit}} = 16$

Readout

Drive Purcell filter $\Omega_f(t)\hat{f}^\dagger + \Omega_f(t)^*\hat{f}$

f0g1 reset

Drive transmon f0-g1 and e-f transitions $\Omega_{f0g1}(t)\hat{n}_t + \Omega_{ef}(t)\hat{n}_t$

Optimizing readout

Objective Maximize SNR in shortest time

Cost function

$$\text{SNR} = \sqrt{2\kappa\eta \int_0^T |\beta_e - \beta_g|^2 dt}$$

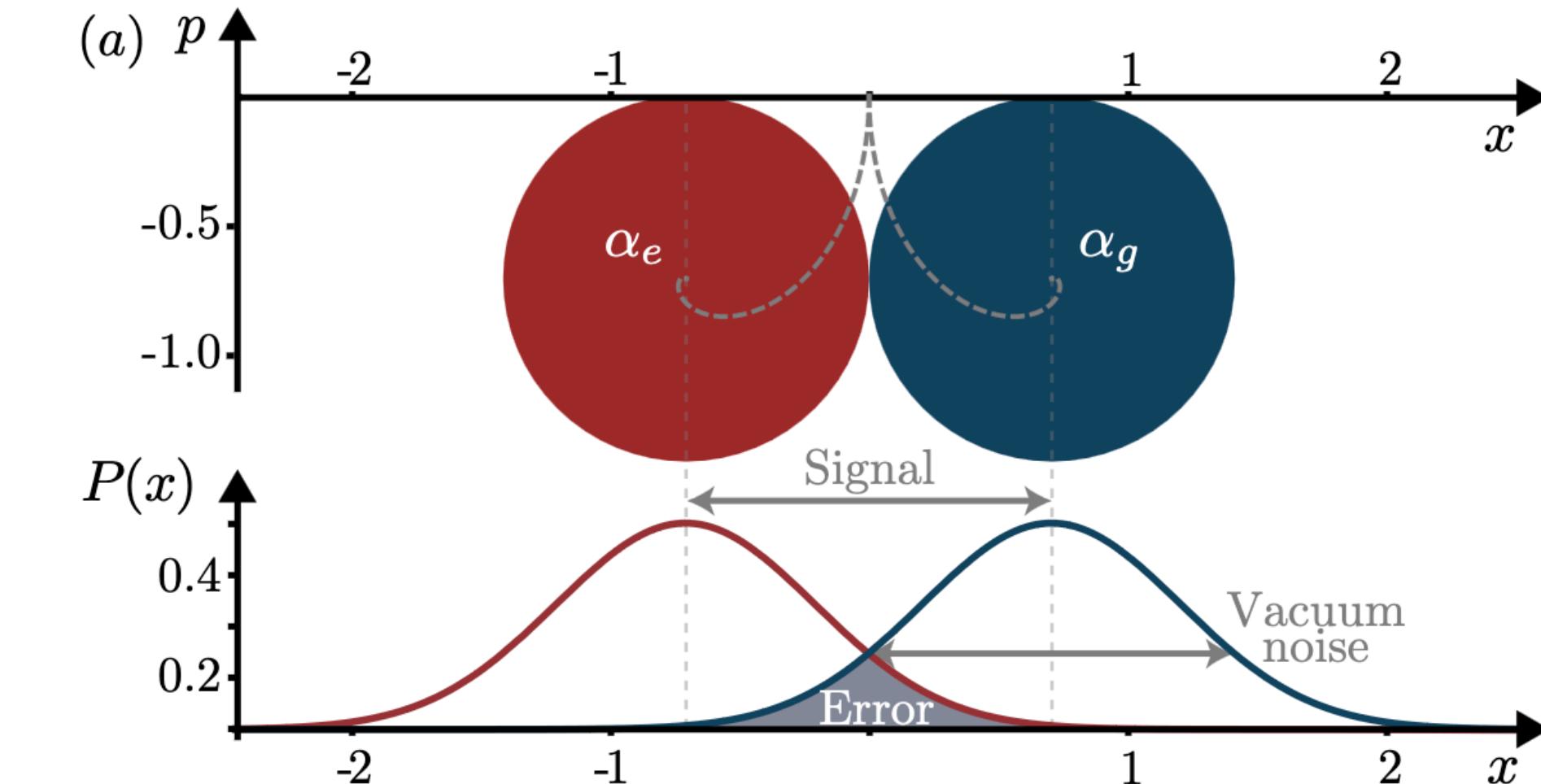
(from Bultink et al., APL 2018)

Parameters

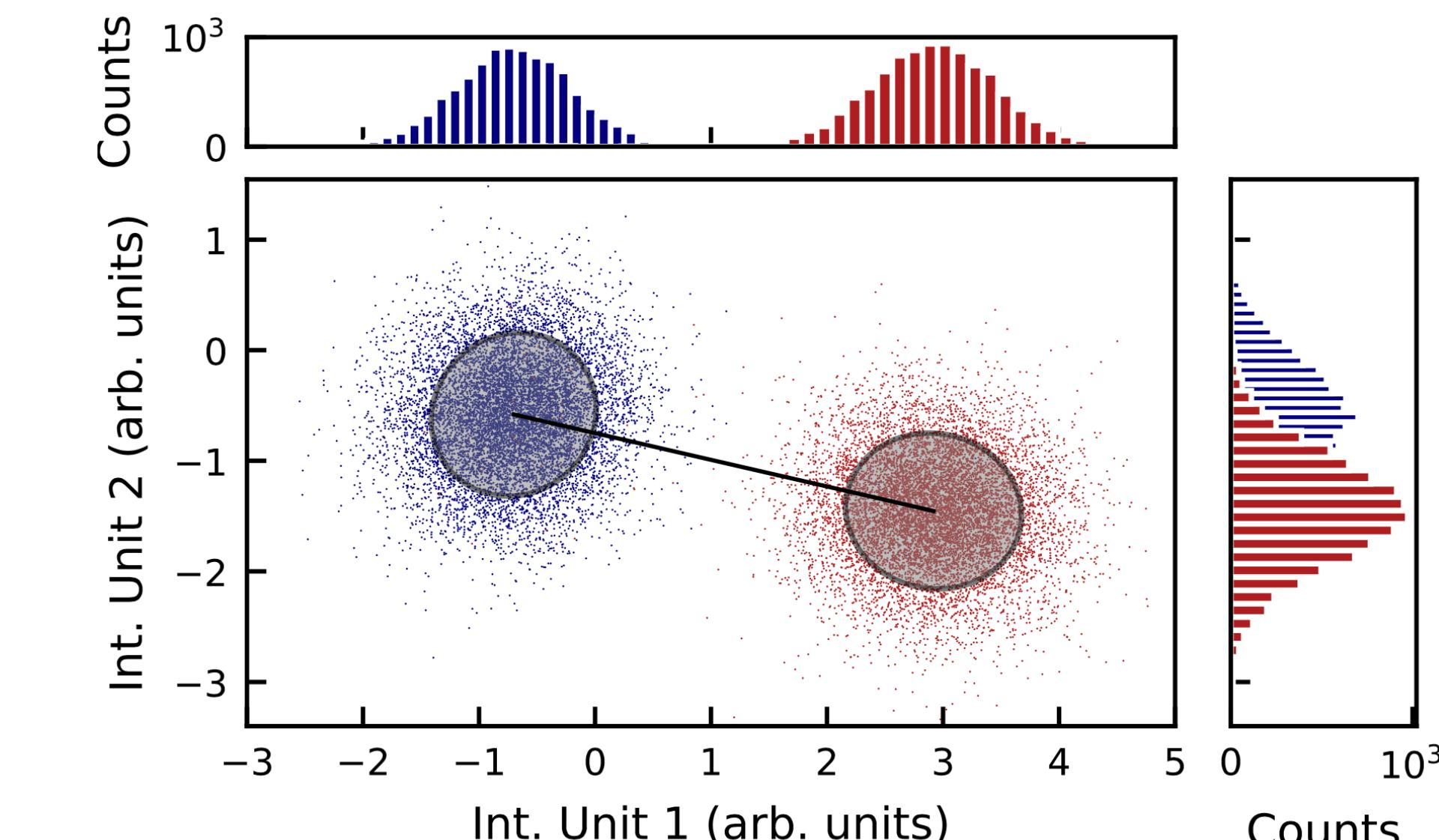
- Pulses (1ns bins, continuous filter)
- Drive frequencies

Other costs

- Resonator population $< \bar{n}_{\text{crit}}$
- Forbidden resonator states
- Maximum pulse amplitudes

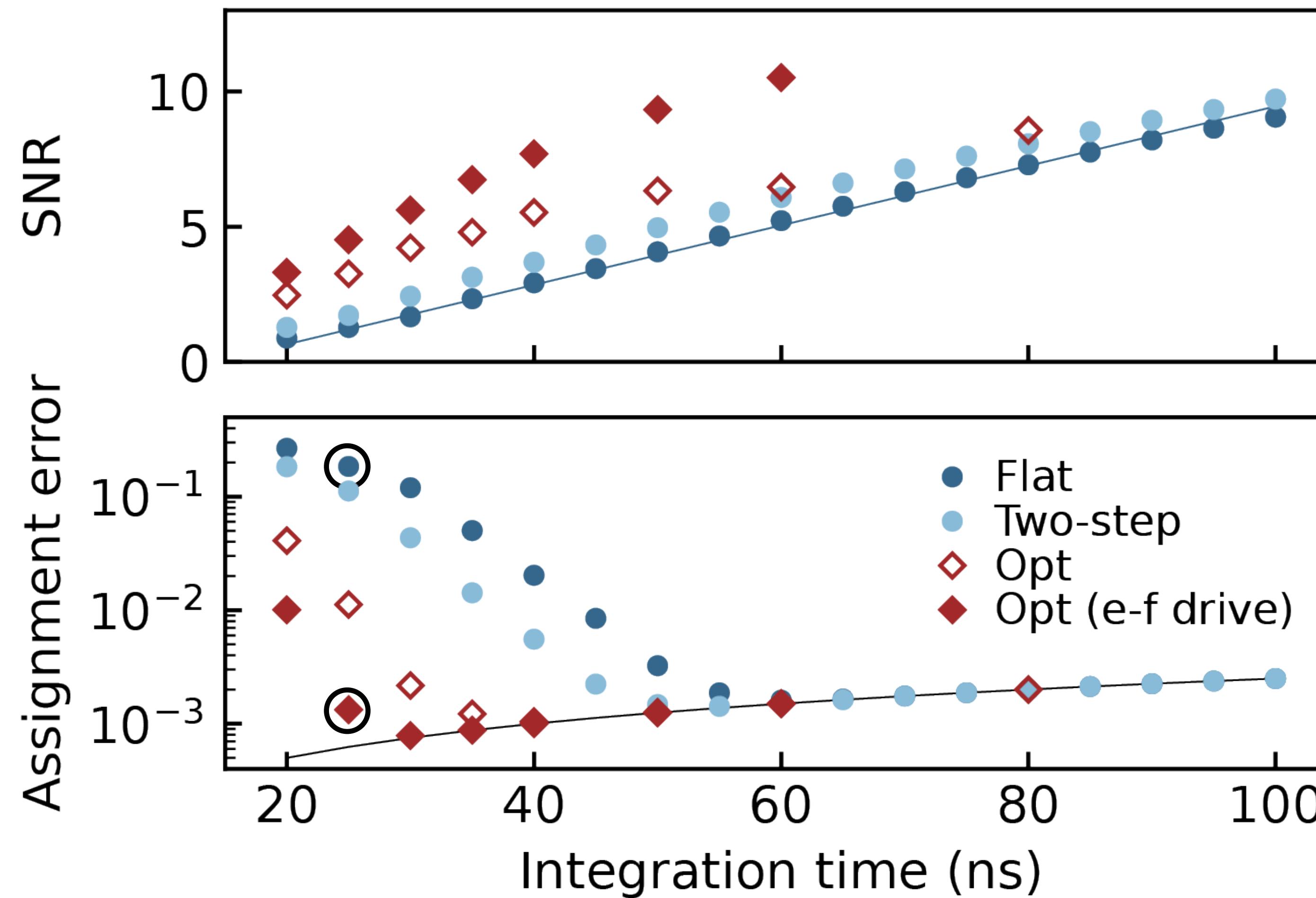


(from Blais et al., RMP 2020)



(from Swiadek et al., in prep.)

Optimizing readout - Preliminary results



Signal-to-Noise Ratio

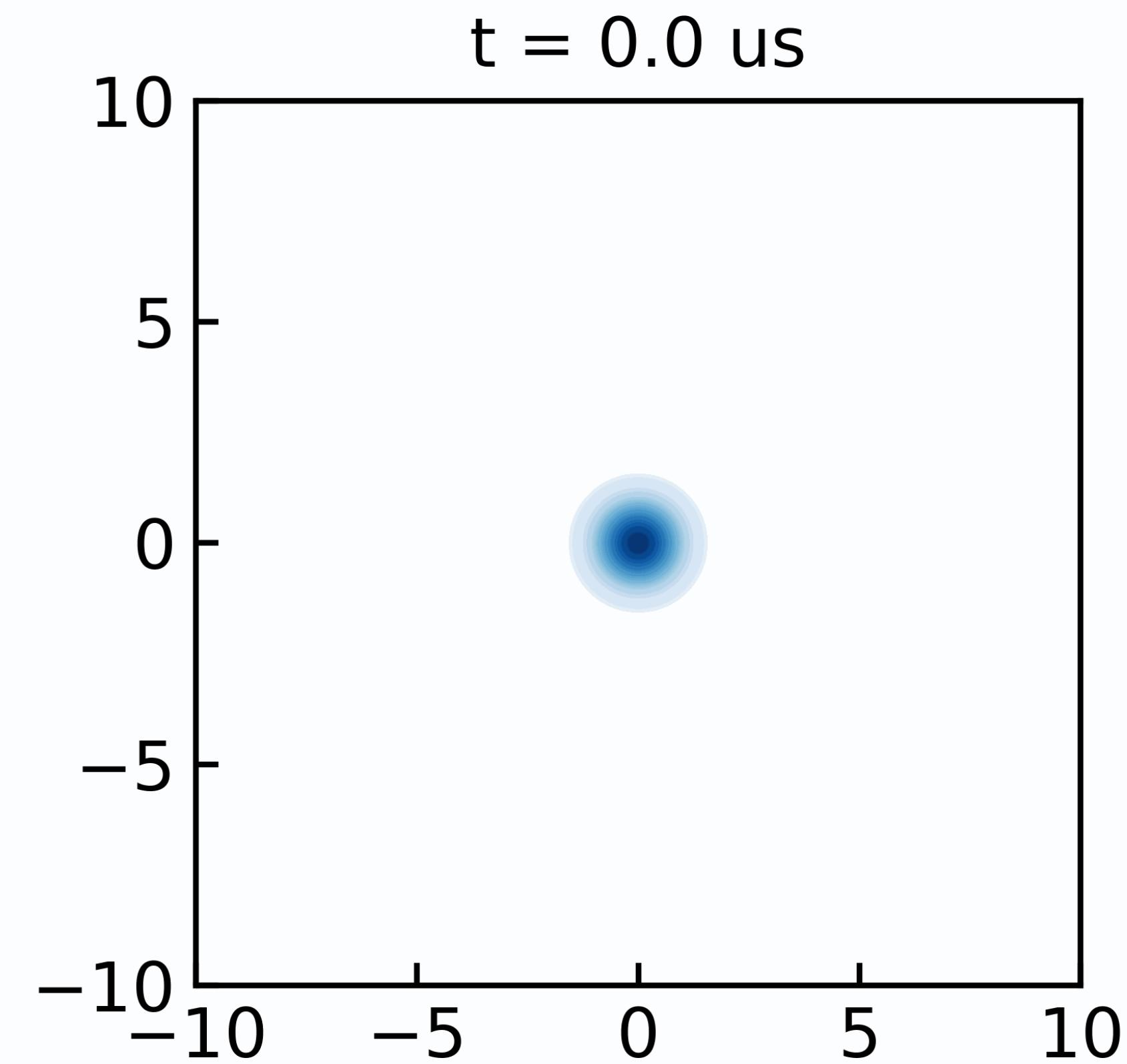
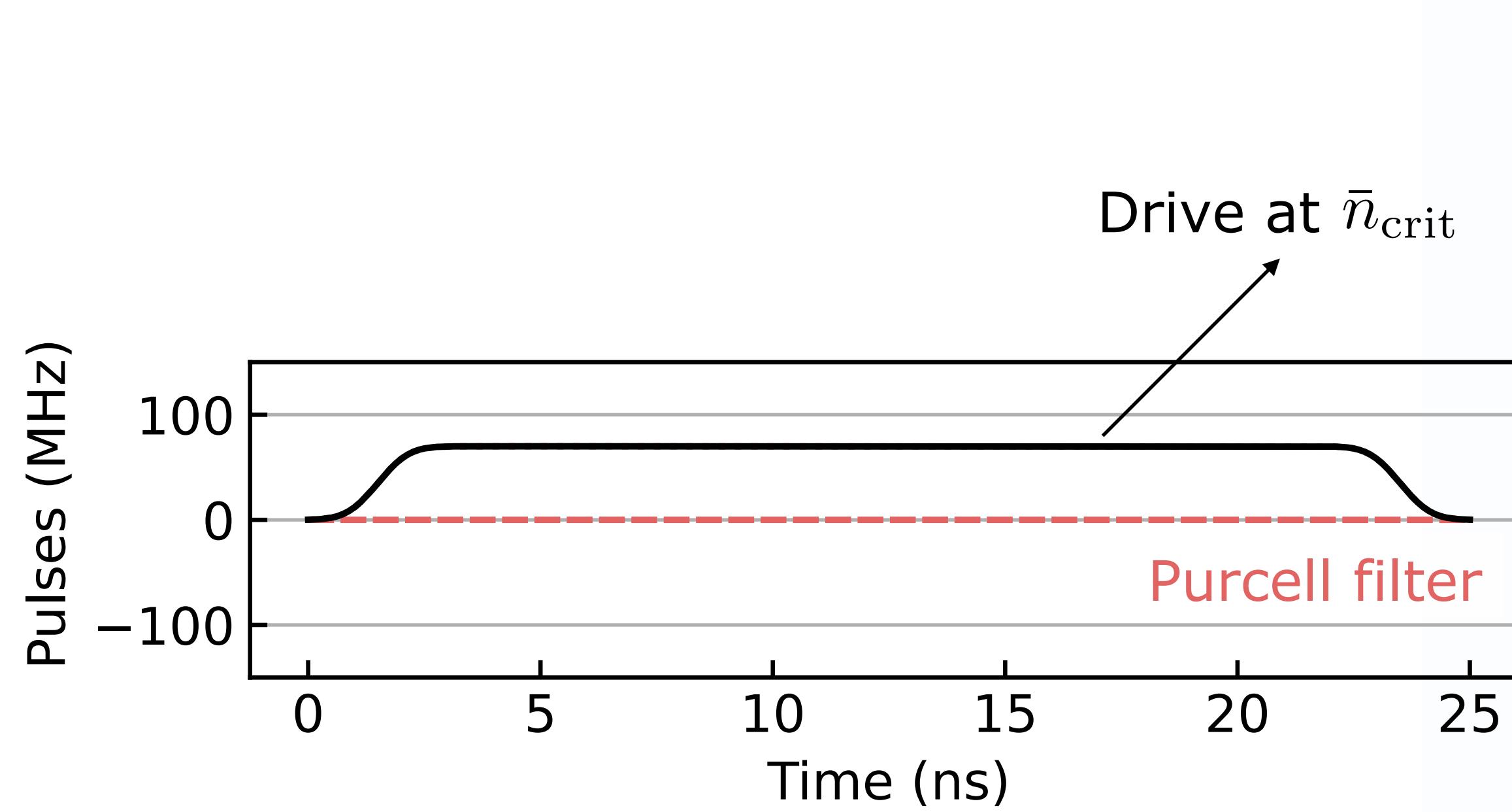
$$\text{SNR} = \sqrt{2\kappa\eta \int_0^T |\beta_e - \beta_g|^2 dt}$$

Assignment error

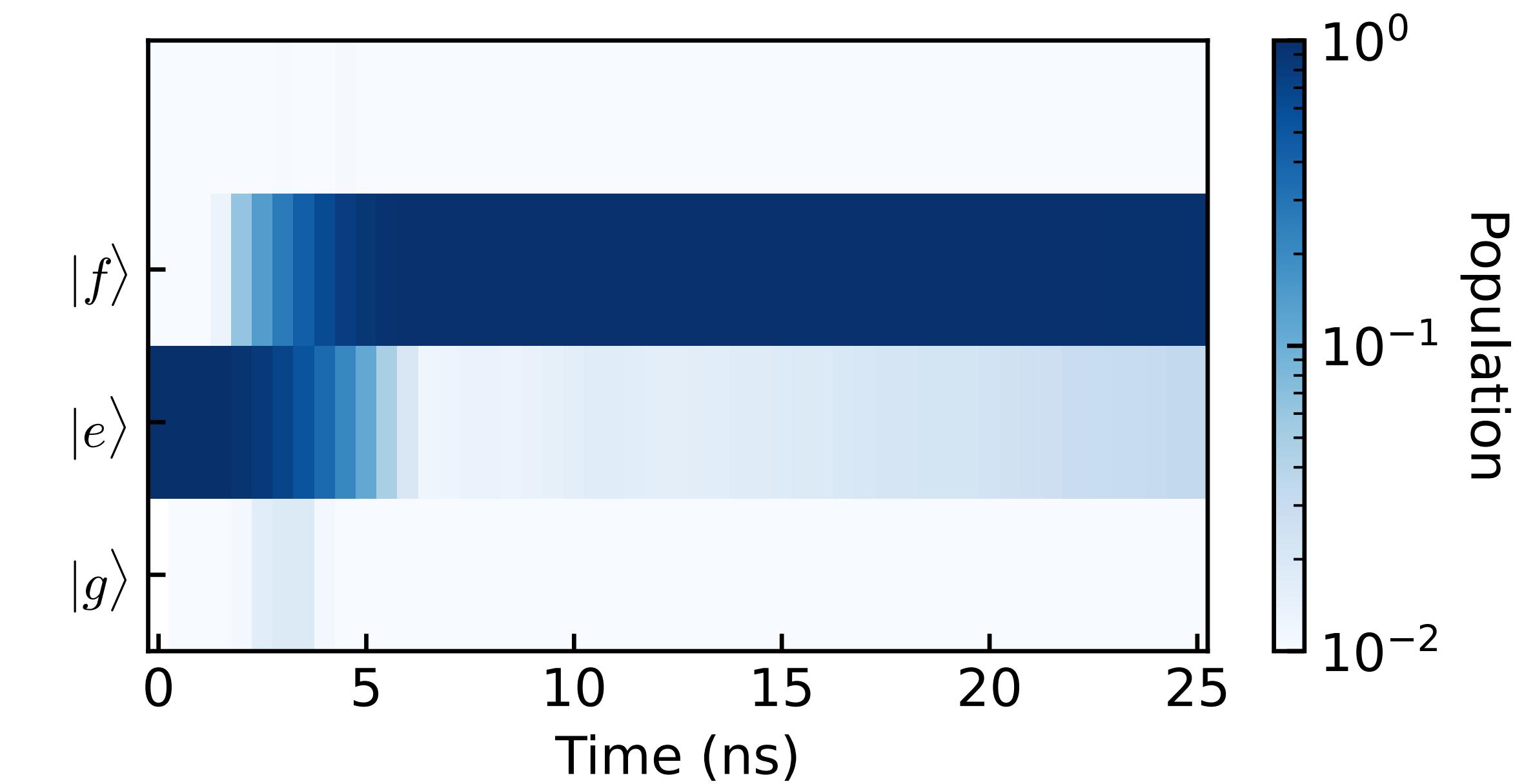
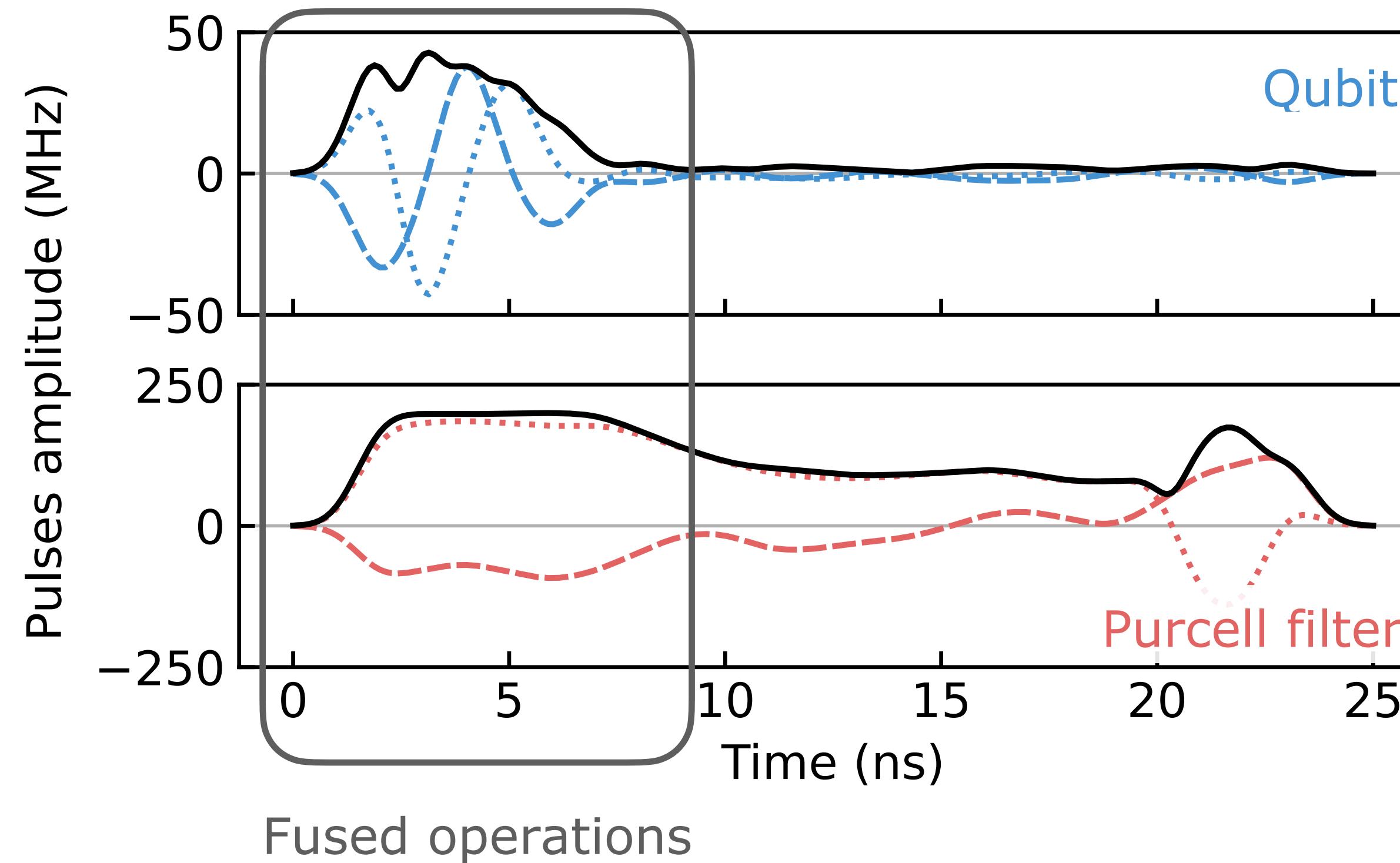
$$\begin{aligned}\varepsilon_a &= \frac{1}{2} (P(e|g) + P(g|e)) \\ &\sim \frac{1}{2} (1 - \text{erf}(\text{SNR}/2) + \tau/T_1)\end{aligned}$$

~55ns to ~30ns !

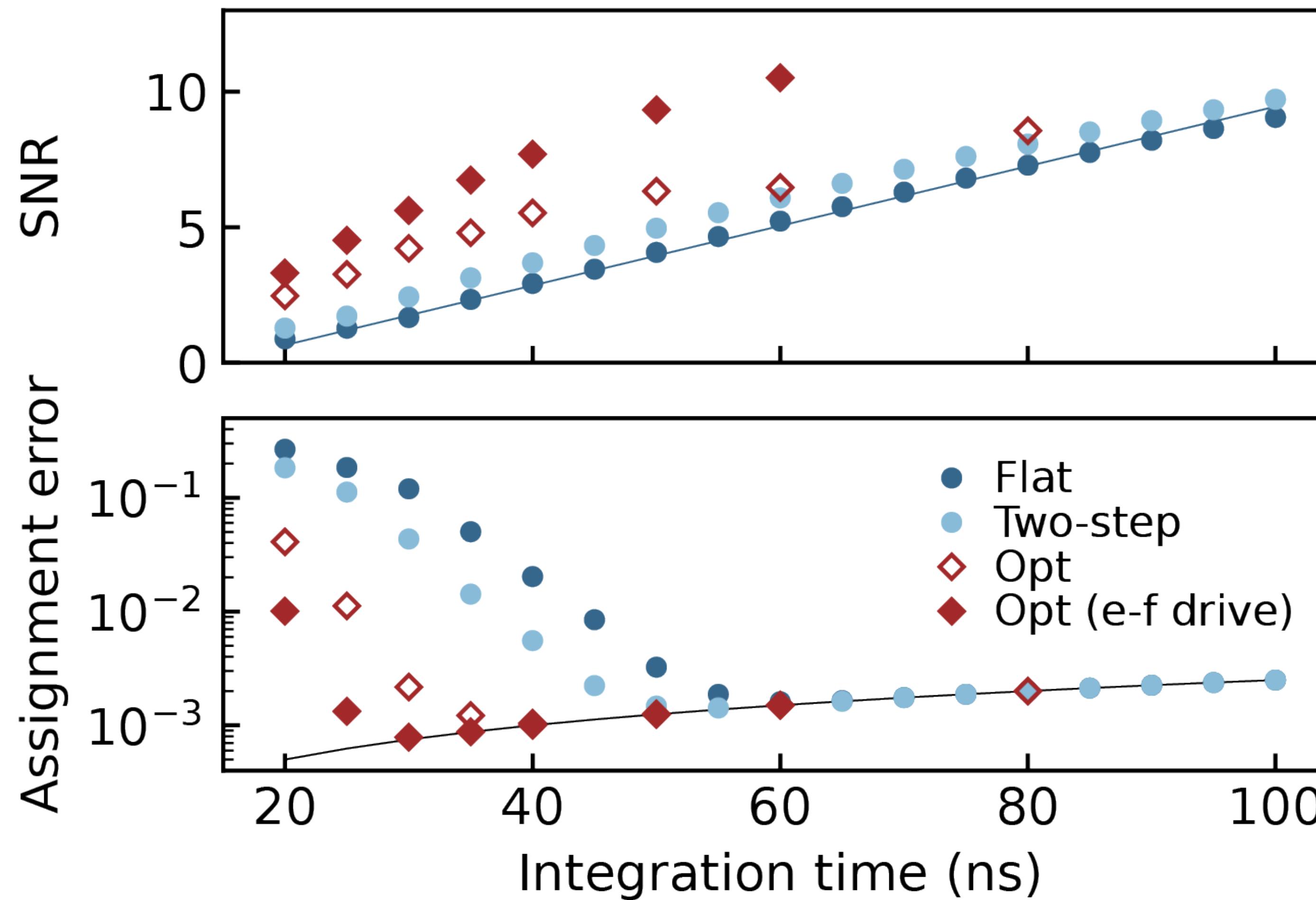
Optimizing readout - Flat readout



Optimizing readout - Preliminary results



Optimizing readout - Preliminary results



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Assignment error

$$\begin{aligned}\varepsilon_a &= \frac{1}{2} (P(e|g) + P(g|e)) \\ &\sim \frac{1}{2} (1 - \text{erf}(\text{SNR}/2) + \tau/T_1)\end{aligned}$$

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A library for efficient differentiable solvers



Differentiable solvers with PyTorch

- Solve SE, ME and SME
- Gradients with automatic or adjoint differentiation



Features

- GPU & CPU support
- Supports batching for density matrices & Hamiltonians
- Useful for optimal control, parameter estimation, state tomography...
- Many solvers (Dormand-Prince 5, Rouchon 1&2, Explicit exponentiation,...)



Open-source @ *github.com/dynamiqs/dynamiqs*

Co-developed with Pierre Guilmin, Adrien Bocquet & Élie Genois



dynamiqs

The dynamiqs library

(1) Works as a drop-in replacement to QuTiP

```
1 import qutip as qt
2 import dynamiqs as dq
3
4 # parameters
5 N = 32
6 delta = 4.0
7 kappa = 0.1
8 alpha = 2.0j
9
10 # operators
11 a = qt.destroy(N)
12 H = delta * a.dag() * a
13 L = sqrt(kappa) * a
14
15 # mesolve arguments
16 rho0 = qt.coherent_dm(N, alpha)
17 tsave = np.linspace(0.0, 1.0, 11)
18
19 # solve ME
20 rhos, _ = dq.mesolve(H, L, rho0, tsave)
```

(2) Compute gradients in a few lines

```
1 import torch
2 import dynamiqs as dq
3
4 # parameters
5 N = 32
6 delta = torch.tensor(4.0, requires_grad=True)
7 kappa = torch.tensor(0.1)
8 alpha = torch.tensor(2.0j, requires_grad=True)
9
10 # operators
11 a, adag = dq.destroy(N), dq.create(N)
12 H = delta * adag @ a
13 L = sqrt(kappa) * a
14
15 # mesolve arguments
16 rho0 = dq.coherent_dm(N, alpha)
17 tsave = torch.linspace(0.0, 1.0, 11)
18
19 # solve ME
20 rhos, _ = dq.mesolve(H, L, rho0, tsave)
21
22 # compute gradient of some loss
23 loss = dq.expect(adag @ a, rhos[-1])
24 loss.backward()
```