

Designing High-Fidelity Gates for Dissipative Cat Qubits

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²Institut quantique et Département de physique, Université de Sherbrooke, Sherbrooke, Canada

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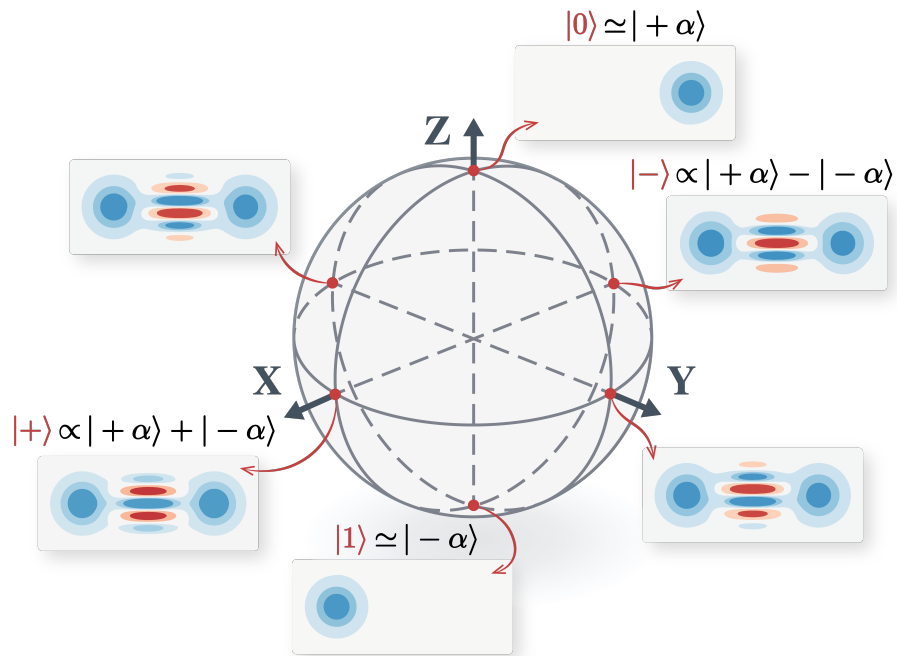
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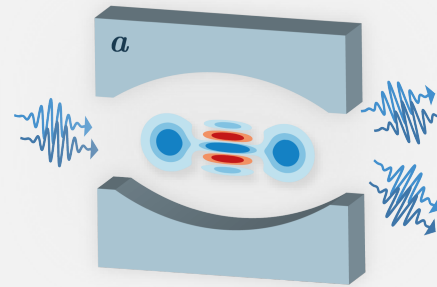
Dissipative Stabilization of Cat Qubits

Cat qubits can be stabilized with two-photon dissipation, often mediated by an **ancillary buffer mode**



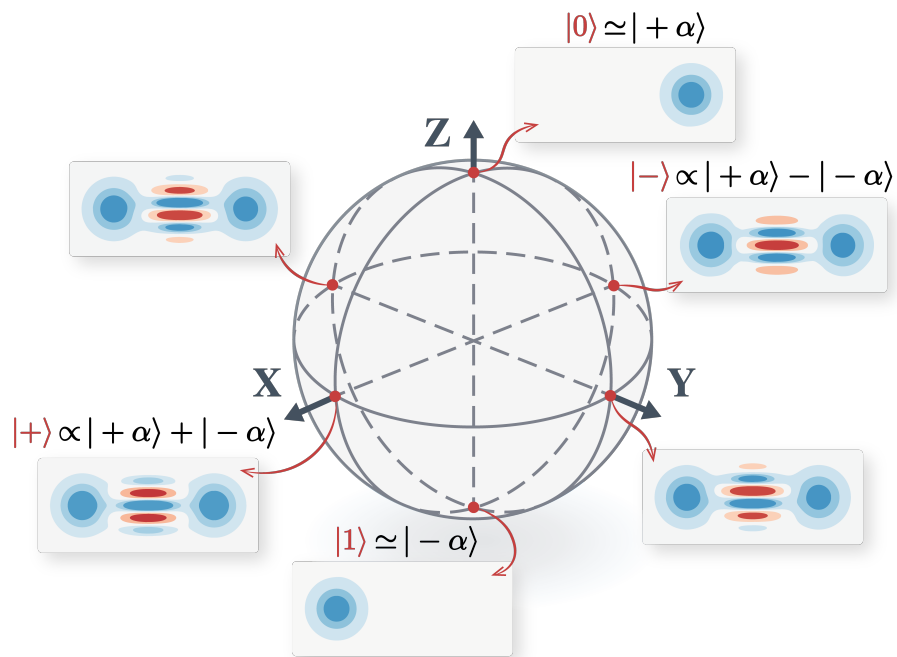
Bloch Sphere Representation of a Cat Qubit

$$\frac{d\rho}{dt} = \kappa_2 \mathcal{D}[a^2 - \alpha^2] \rho$$



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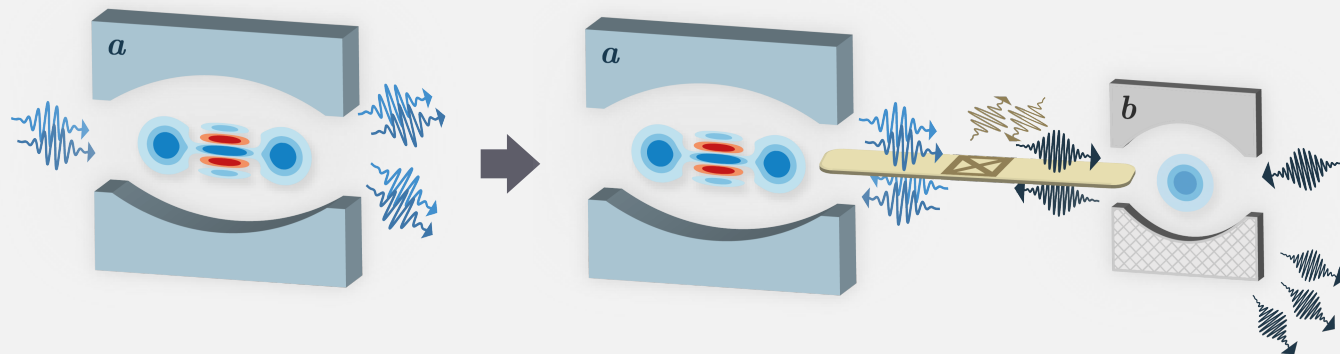
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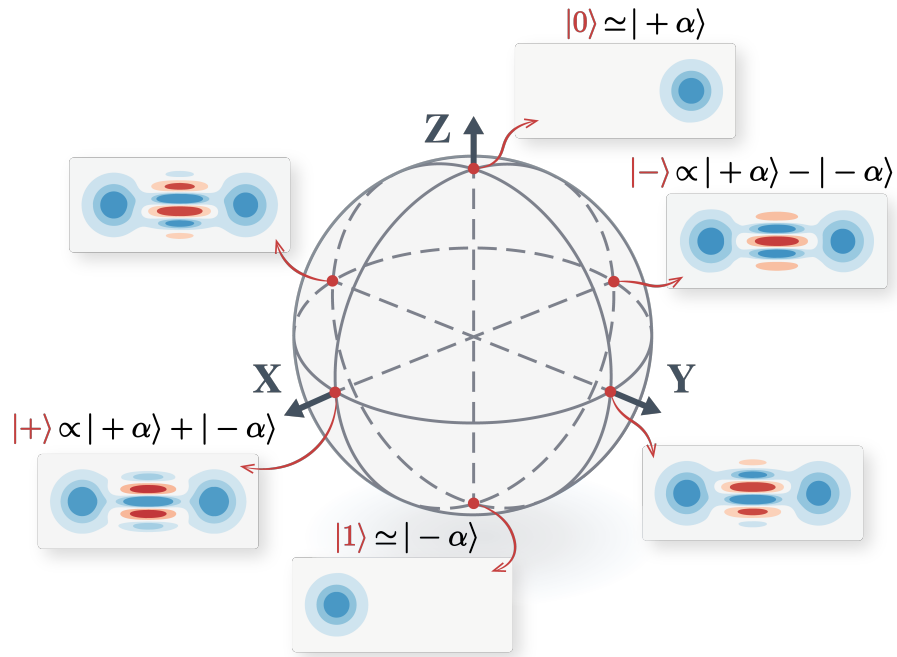
$$\frac{d\rho}{dt} = \kappa_2 \mathcal{D}[a^2 - \alpha^2] \rho \quad \rightarrow \quad \frac{d\rho}{dt} = -i[\mathbf{H}_{AB}, \rho] + \kappa_b \mathcal{D}[b] \rho$$

with $\mathbf{H}_{AB} = g_2(a^2 - \alpha^2)b^\dagger + \text{h.c.}$



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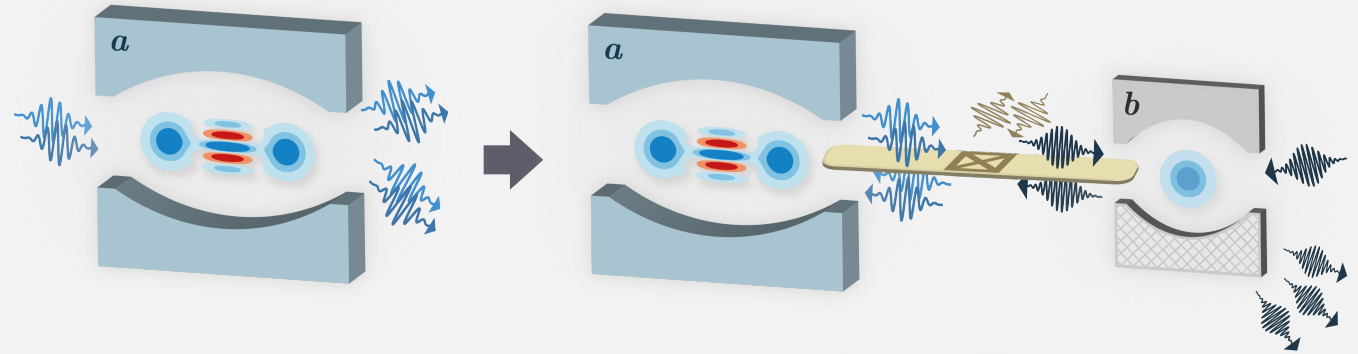
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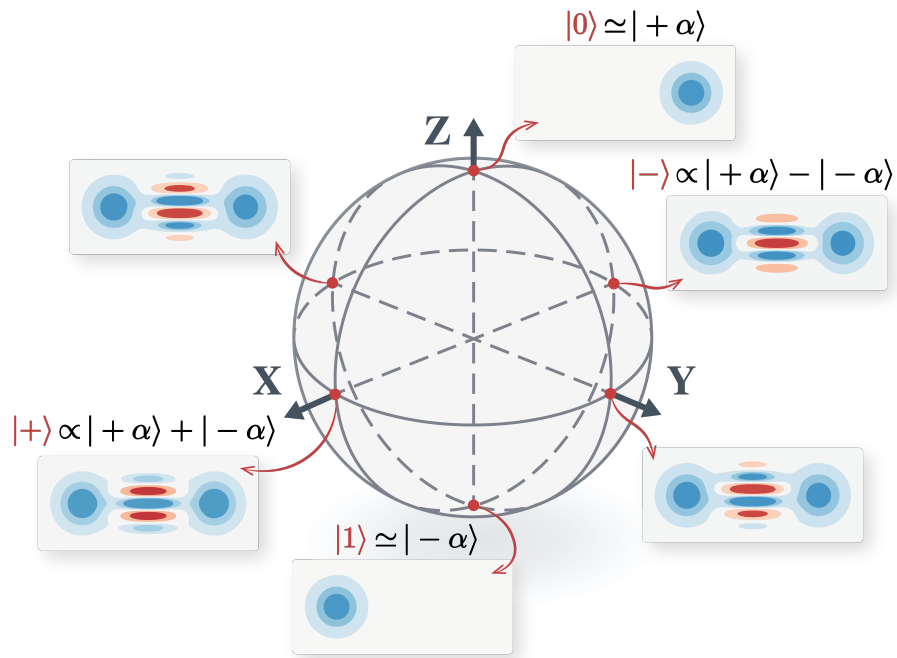
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- Adiabatic elimination $\kappa_2 \equiv \frac{4g_2^2}{\kappa_b}$

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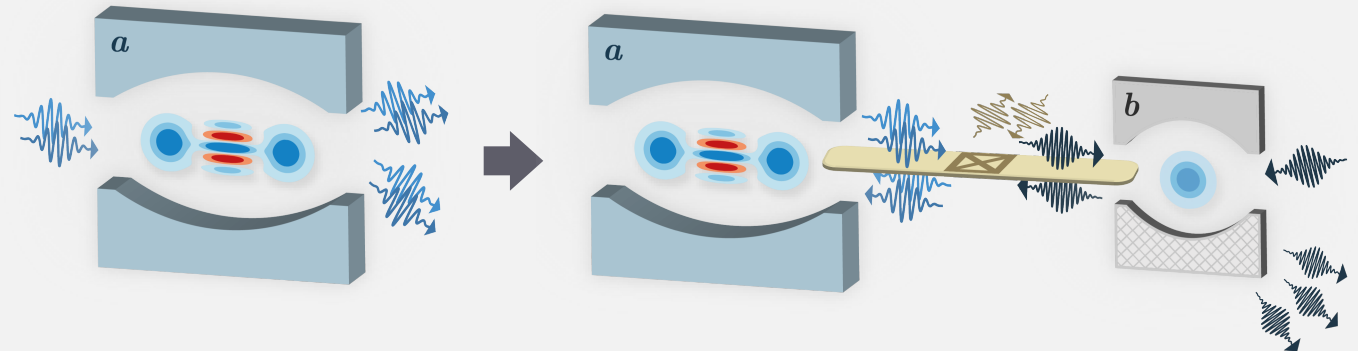
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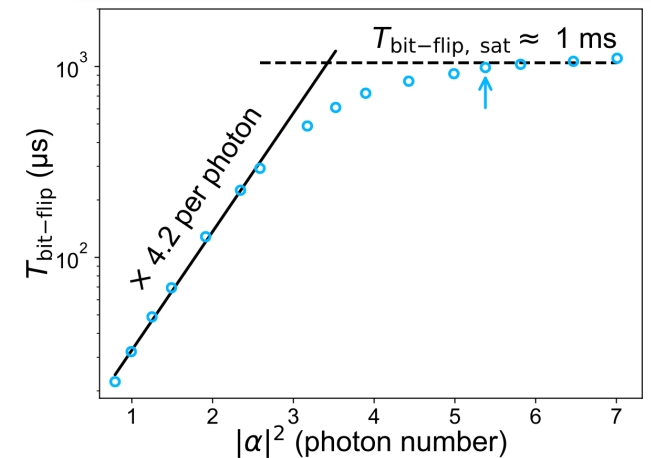
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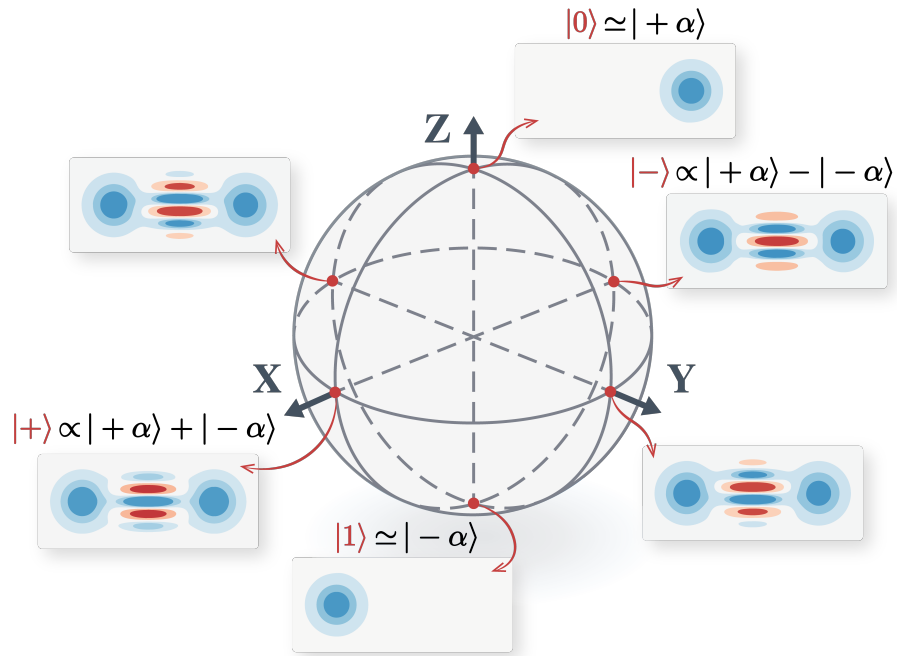


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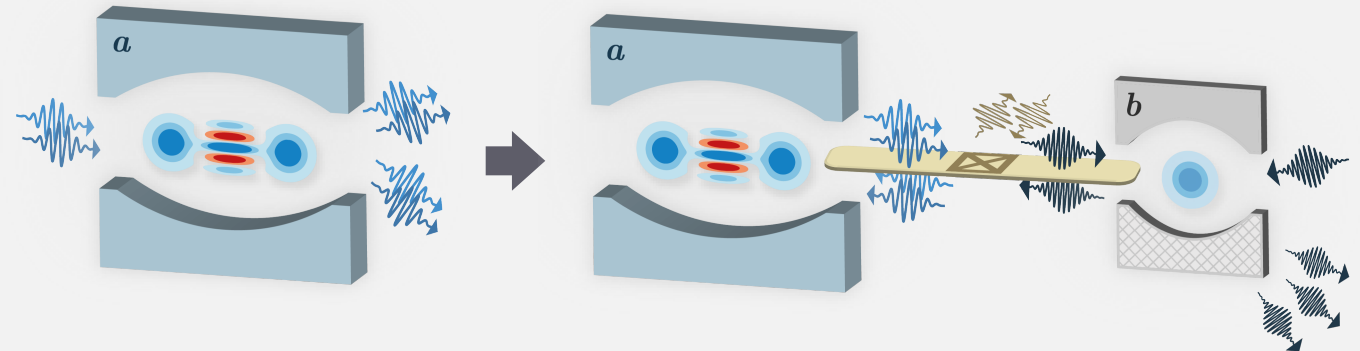
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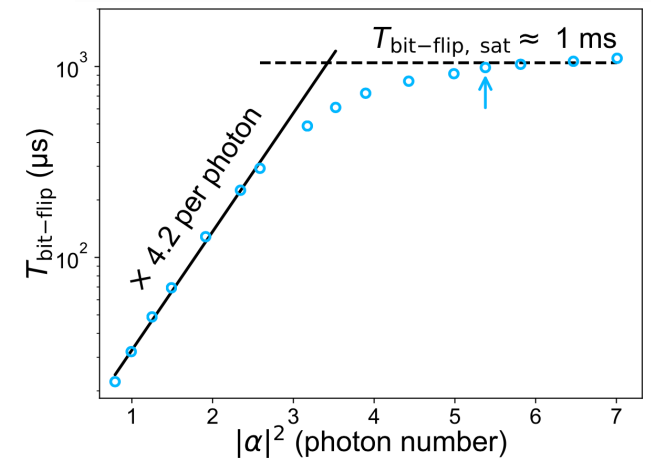
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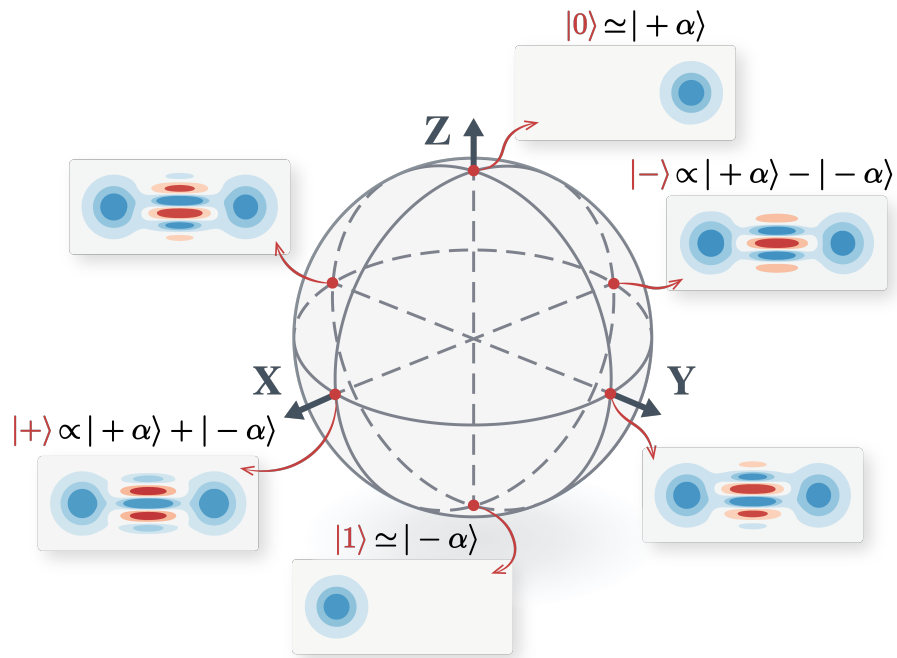


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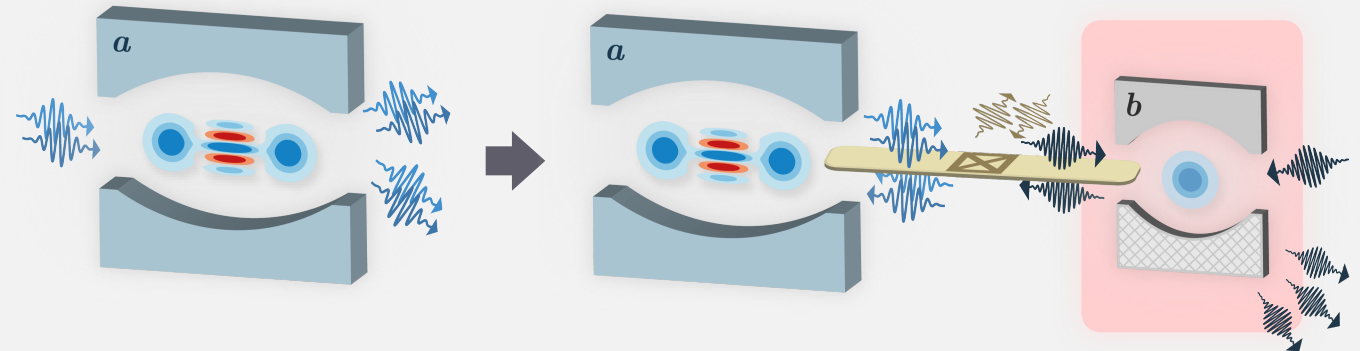
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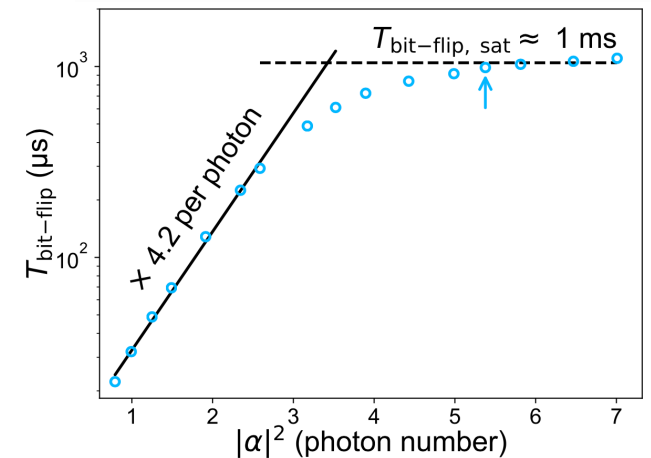
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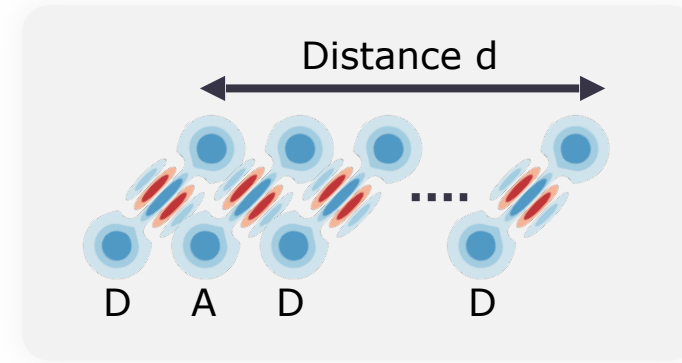


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- ➔ **Gate engineering**



Repetition Cat Qubits

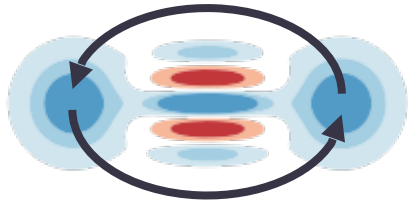
For fault-tolerant universal QC with repetition cat qubits, only four gates are required on top of preparation and measurement in $|\pm_L\rangle$



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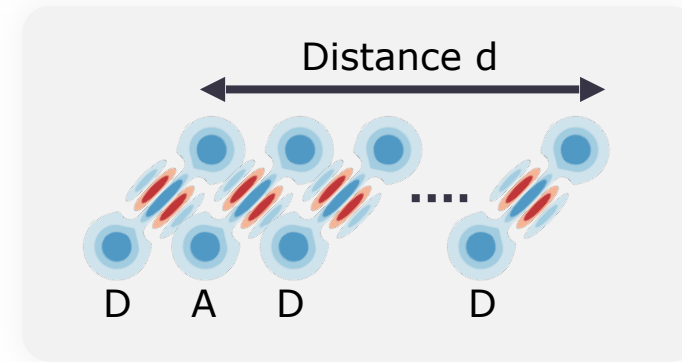
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Pauli X



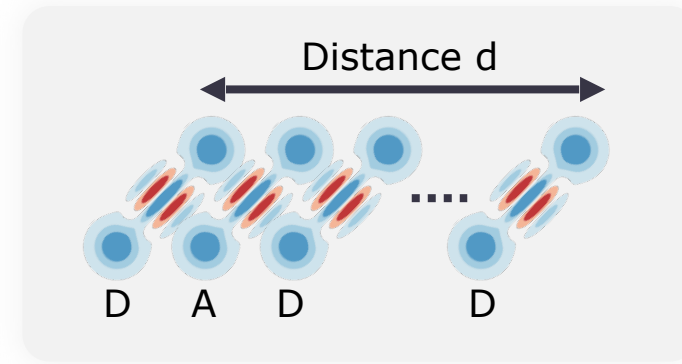
$$H = \Delta_X a^\dagger a$$

(\sim error-less)

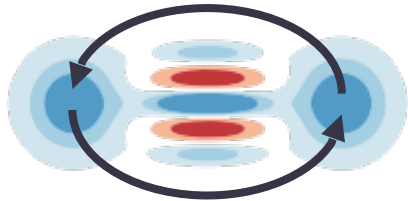


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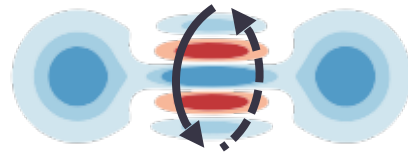
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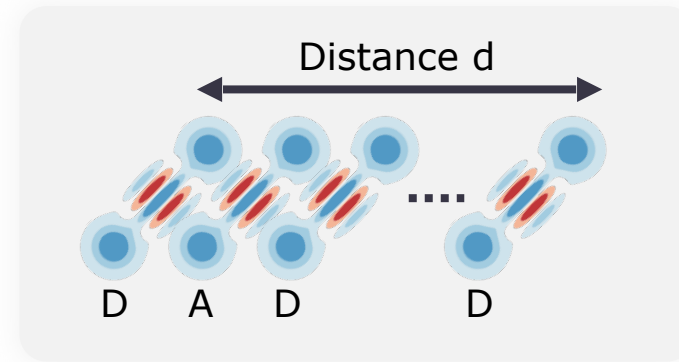
$Z(\theta)$ gate



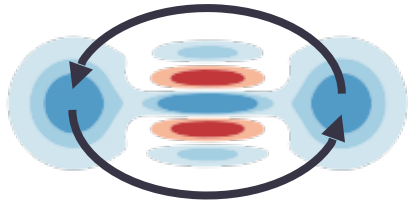
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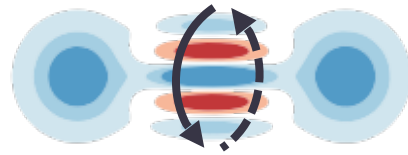
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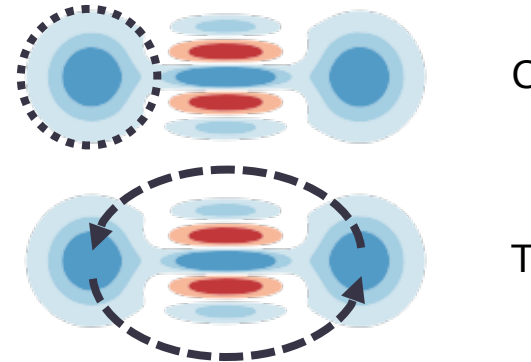
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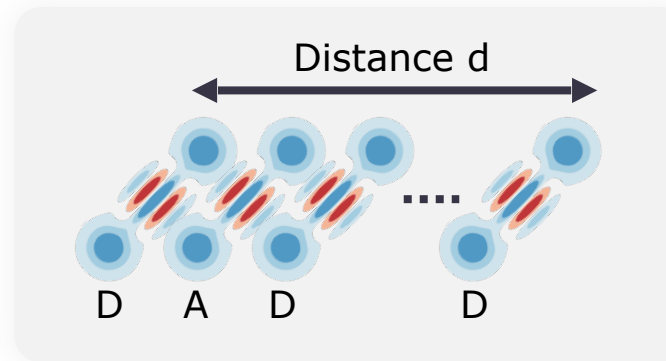
CNOT



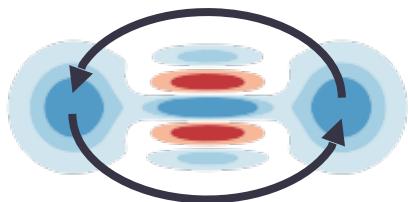
$$H = \varepsilon_{CX} (a_C^\dagger + a_C - 2\alpha) \times (a_T^\dagger a_T - |\alpha|^2)$$

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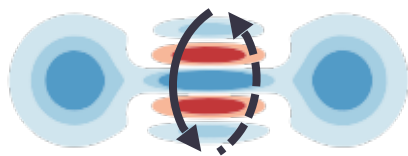
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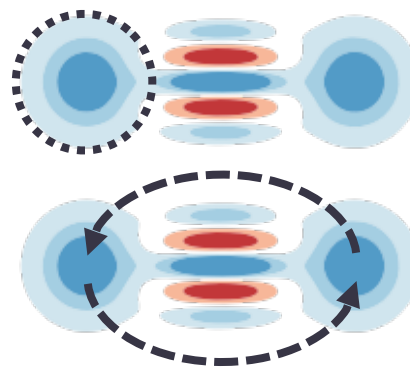
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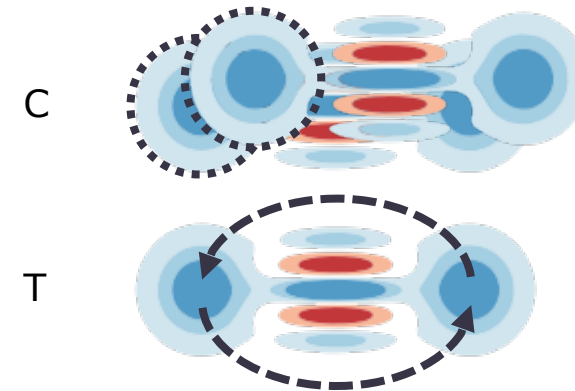
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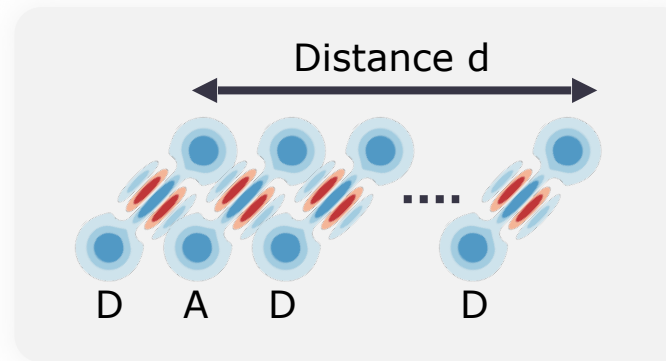
Toffoli



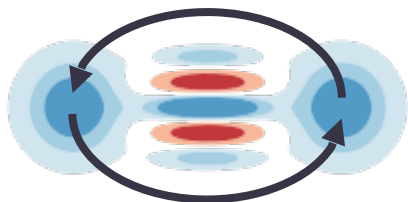
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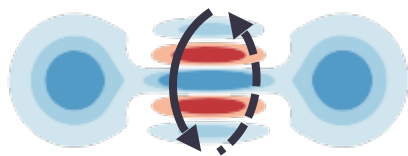
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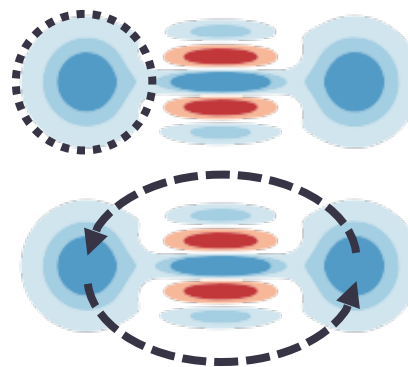
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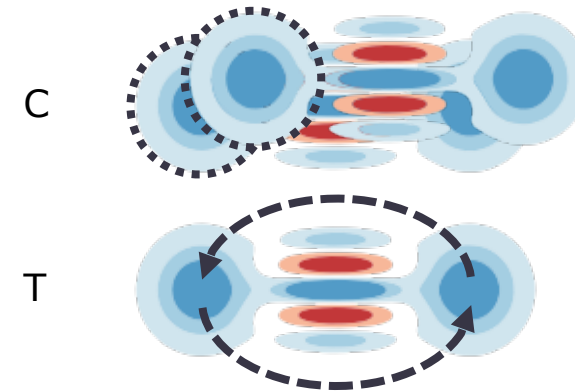
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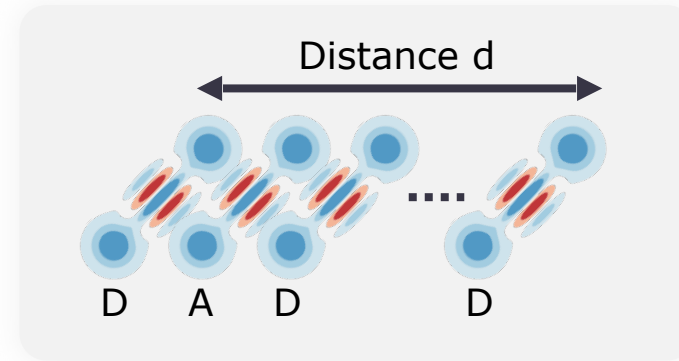
Gate-induced errors are only on **control** qubits

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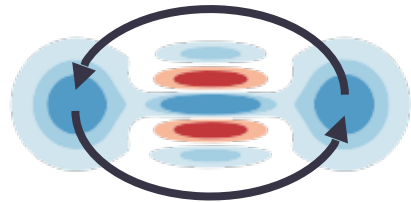
Gate errors Cavity lifetime

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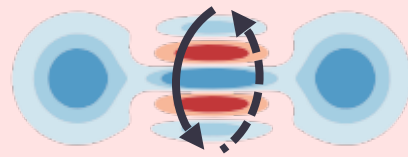
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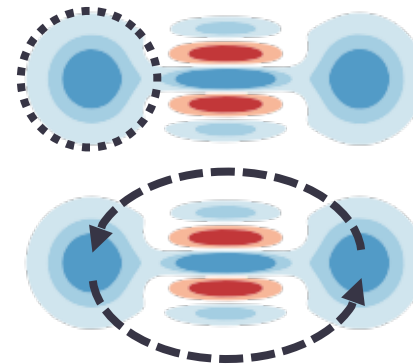
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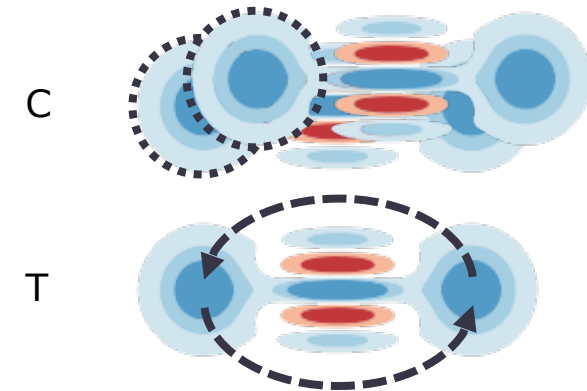
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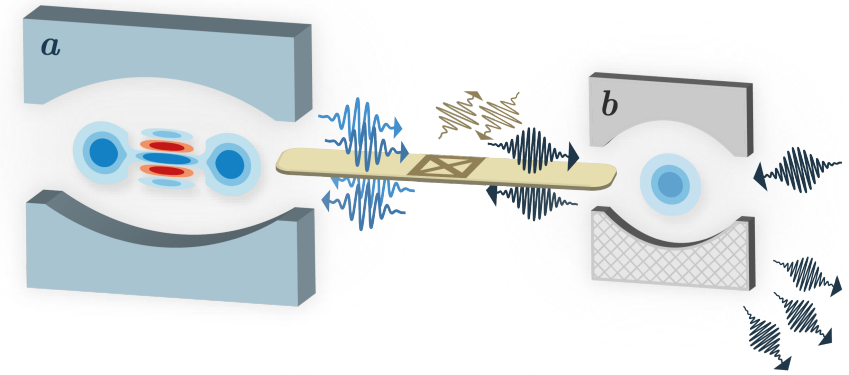
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Gate errors Cavity lifetime

On the Origin of Gate Errors

Consider the Z gate Hamiltonian with a buffer mode,

$$H = g_2(\mathbf{a}^2 - \alpha^2)\mathbf{b}^\dagger + \varepsilon_Z\mathbf{a} + \text{h.c.}$$



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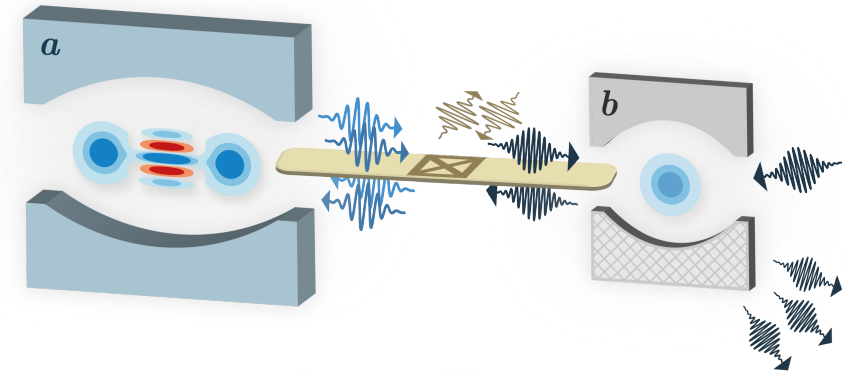
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Introduce the Shifted Fock Basis

$$\mathbf{a} \rightarrow \underbrace{\sigma_z}_{\text{qubit}} \otimes \underbrace{(\tilde{\mathbf{a}} + \alpha)}_{\text{gauge}} \rightarrow \text{Bloch sphere} \otimes \text{Fock states}$$

$$H = g_2(\tilde{\mathbf{a}}^2 + 2\alpha\tilde{\mathbf{a}})\mathbf{b}^\dagger + \varepsilon_Z \sigma_z (\tilde{\mathbf{a}} + \alpha) + \text{h.c.}$$

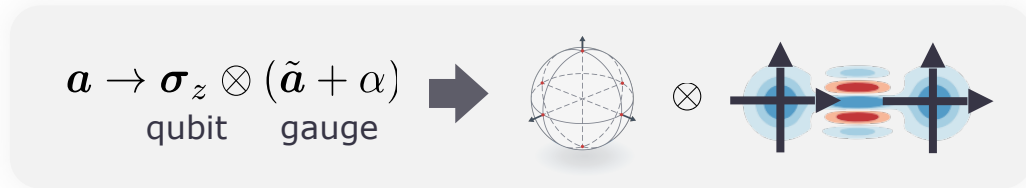


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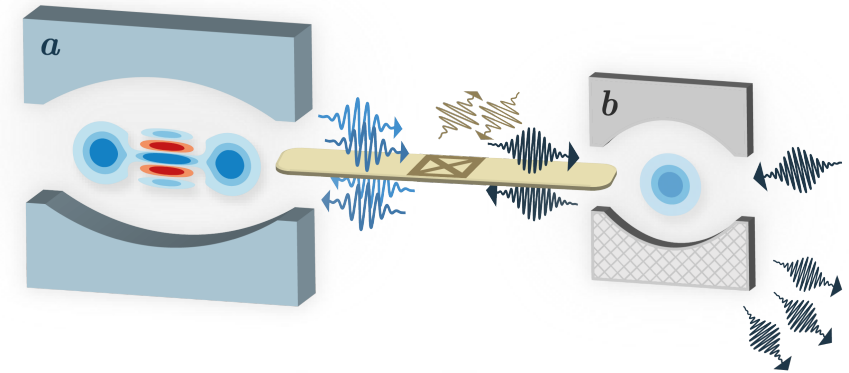


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Write master equation in Heisenberg picture

$$\dot{\tilde{\mathbf{a}}} = -2i\alpha g_2 \mathbf{b} - i\varepsilon_Z \sigma_z$$

$$\dot{\mathbf{b}} = -2i\alpha g_2 \tilde{\mathbf{a}} - \kappa_b \mathbf{b}/2$$



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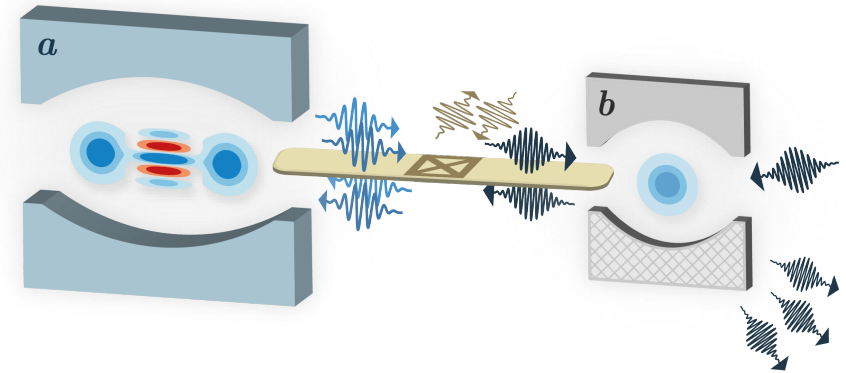
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Decouple these equations

$$\ddot{\mathbf{b}} + \kappa \dot{\mathbf{b}} + \omega_0^2 \mathbf{b} = -\omega_0 \varepsilon_Z \sigma_z \Rightarrow \mathbf{b}_{eq} = -\varepsilon_Z / \omega_0 \sigma_z$$

$$\text{with } \kappa = \kappa_b/2 \text{ and } \omega_0 = 2\alpha g_2$$



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qubit gauge



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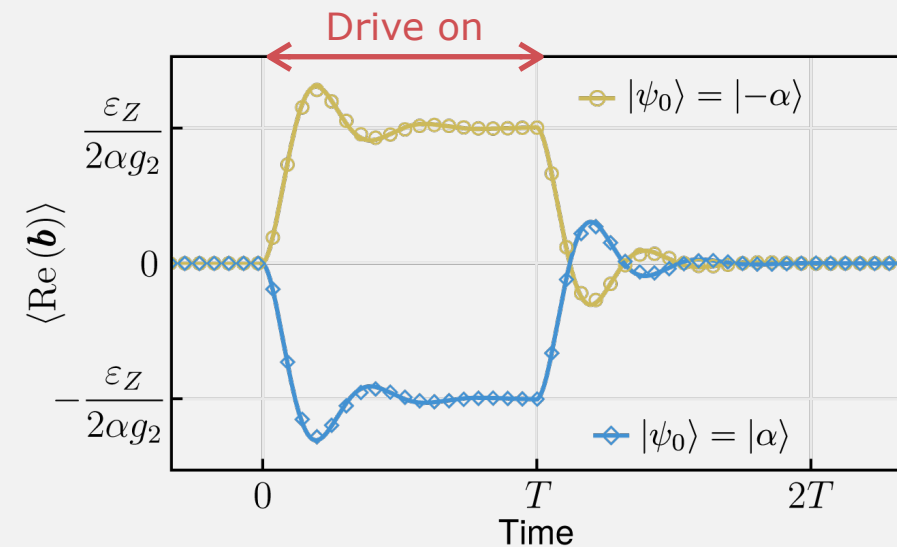
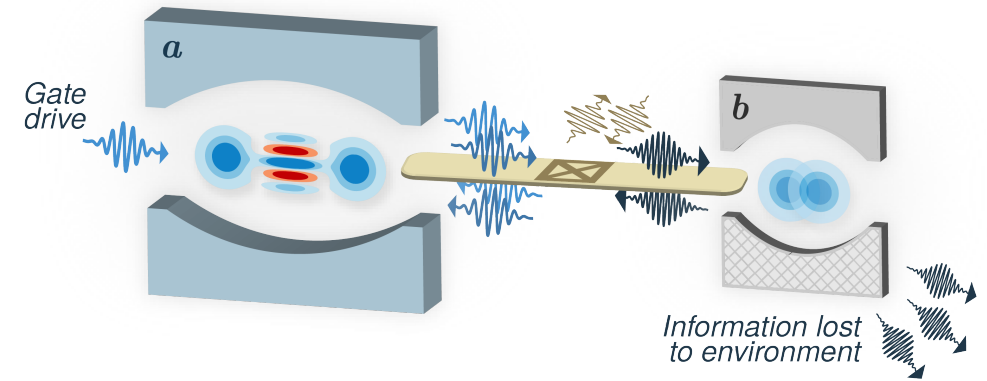
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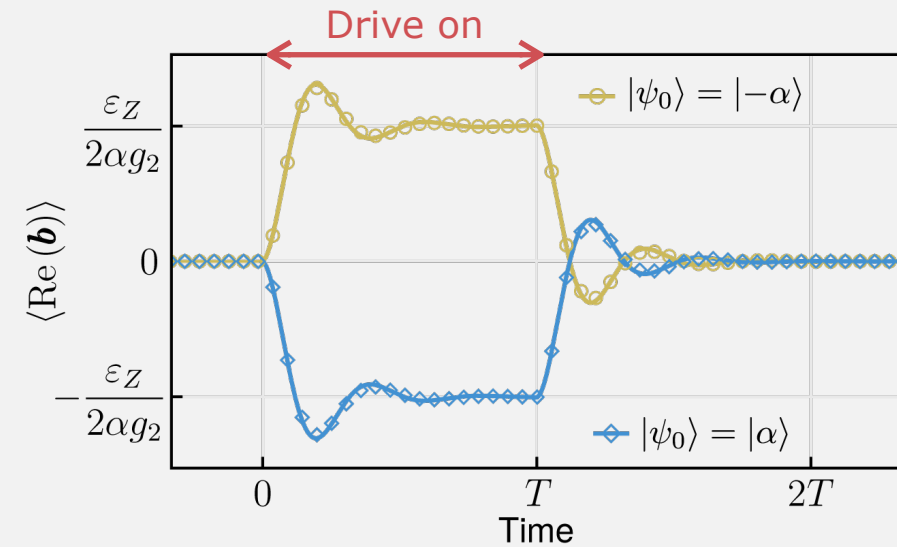
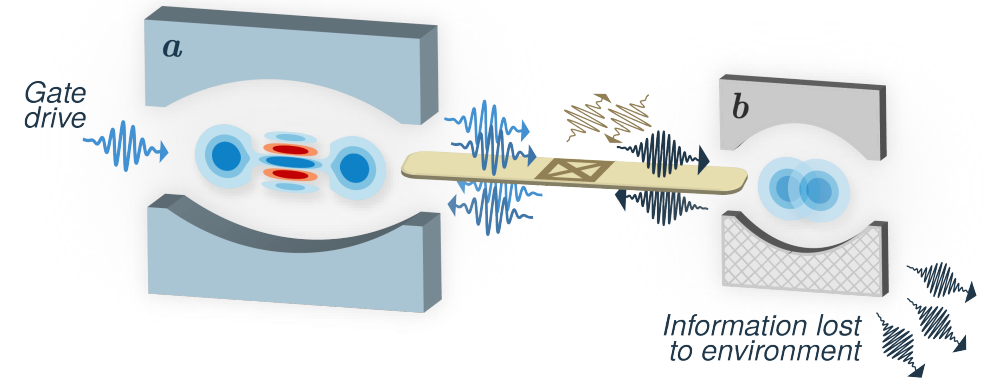
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Can re-derive Zeno gate errors

$$D[b] \approx D[b_{\text{eq}}] = \frac{\varepsilon_Z^2}{\omega_0^2} D[\sigma_z] \Rightarrow p_Z = \frac{\pi^2}{16|\alpha|^2 T} \frac{\kappa_b}{4g_2^2}$$

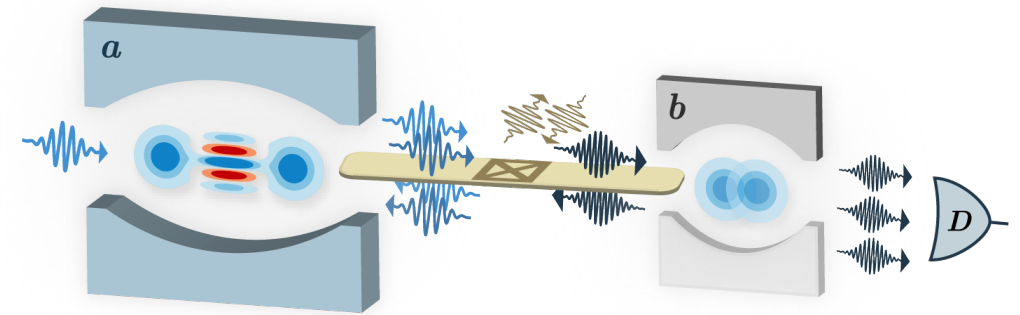
Buffer Photodetection – Rabi oscillations

Information is lost through the buffer mode

→ **measure the buffer output** to retrieve it

$$d\rho = -i[H, \rho]dt + \underbrace{\kappa_b (\mathcal{D}_\eta[b]\rho dt)}_{\text{no-jump}} + \underbrace{\mathcal{J}[b]\rho dN_\eta}_{\text{jump}}$$

with $0 \leq \eta \leq 1$ (detection efficiency)



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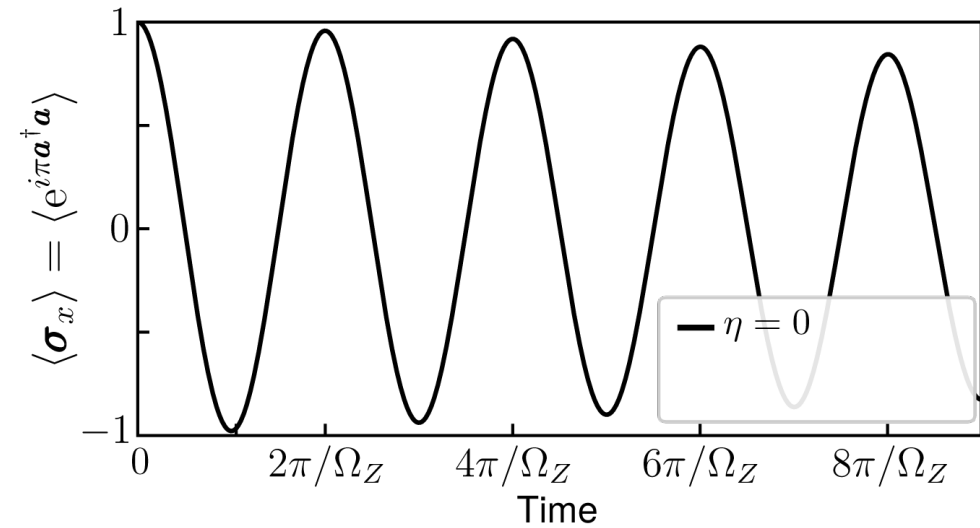
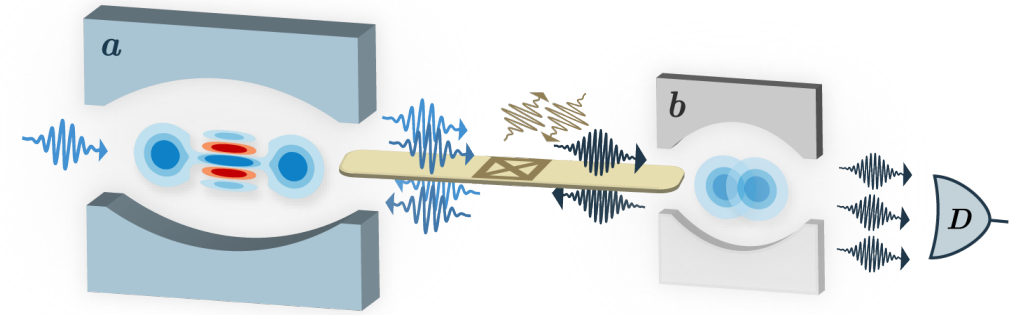
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no-jump

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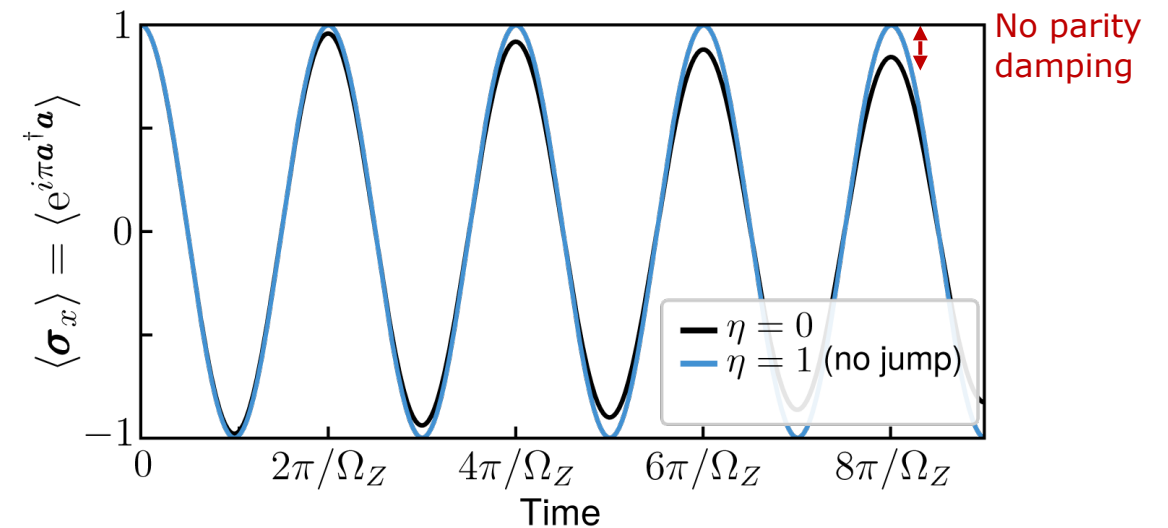
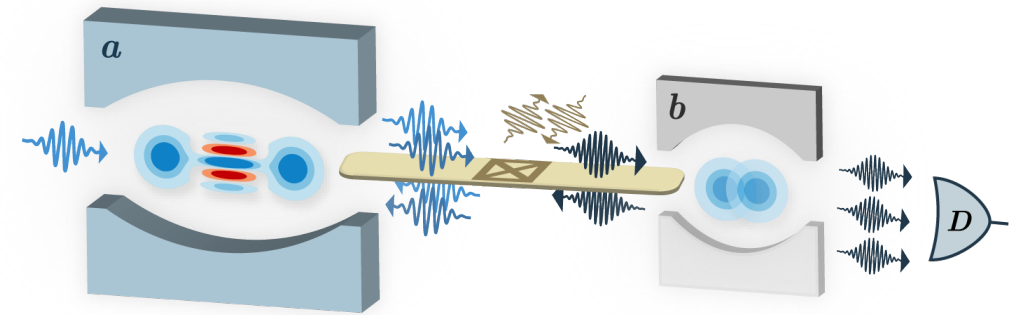
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Assuming $\eta = 1$

- Preserves purity (no information lost)



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no-jump

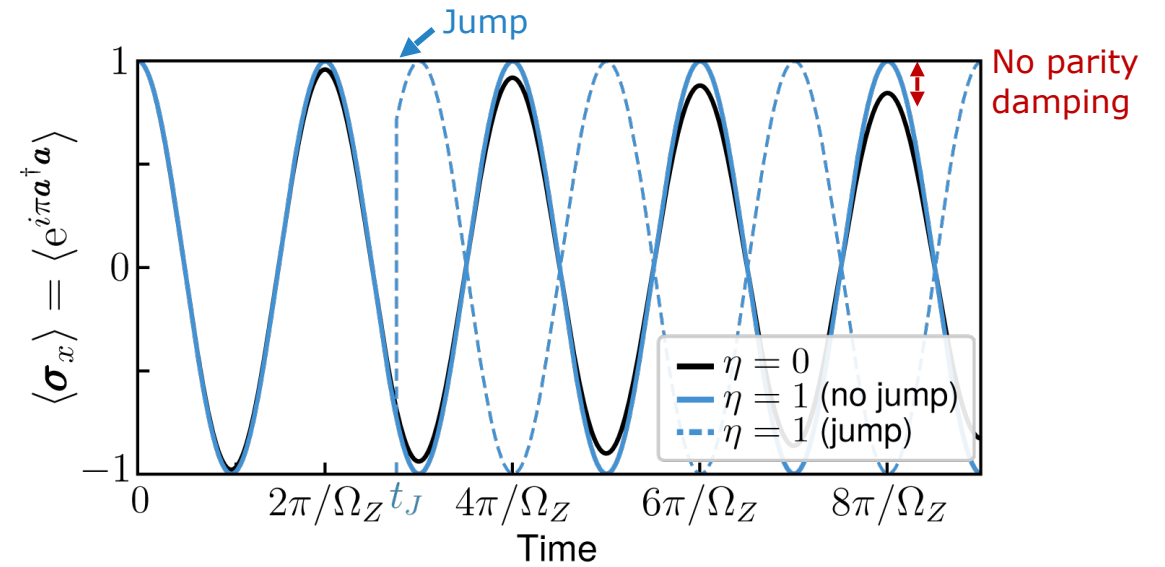
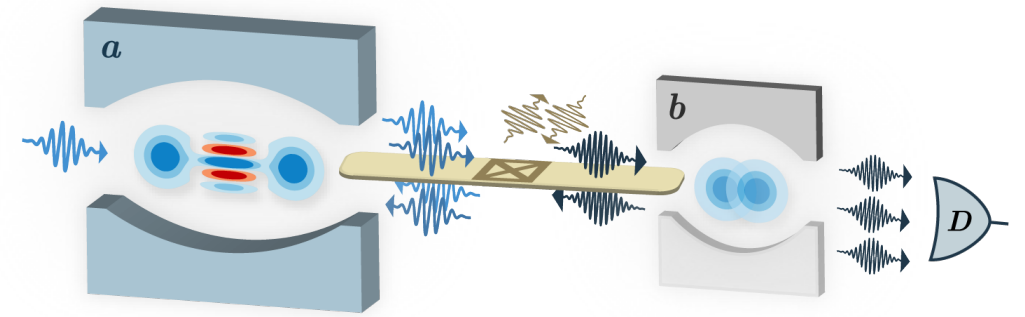
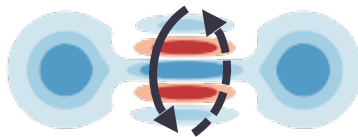
jump

with $0 \leq \eta \leq 1$ (detection efficiency)

Assuming $\eta = 1$

- Preserves purity (no information lost)
- Jump detected = parity swap

→ $b \propto \sigma_z$



Buffer Photodetection – Gates and buffer relaxation

Information is lost through the buffer mode

→ measure the buffer output to retrieve it + **classical feedback**

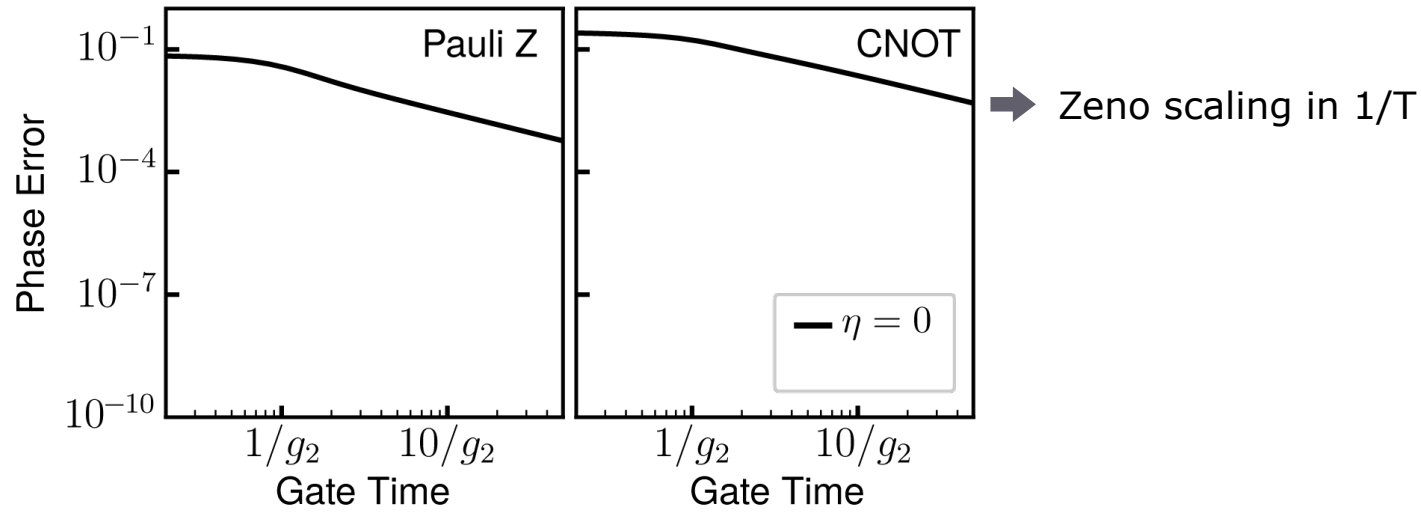
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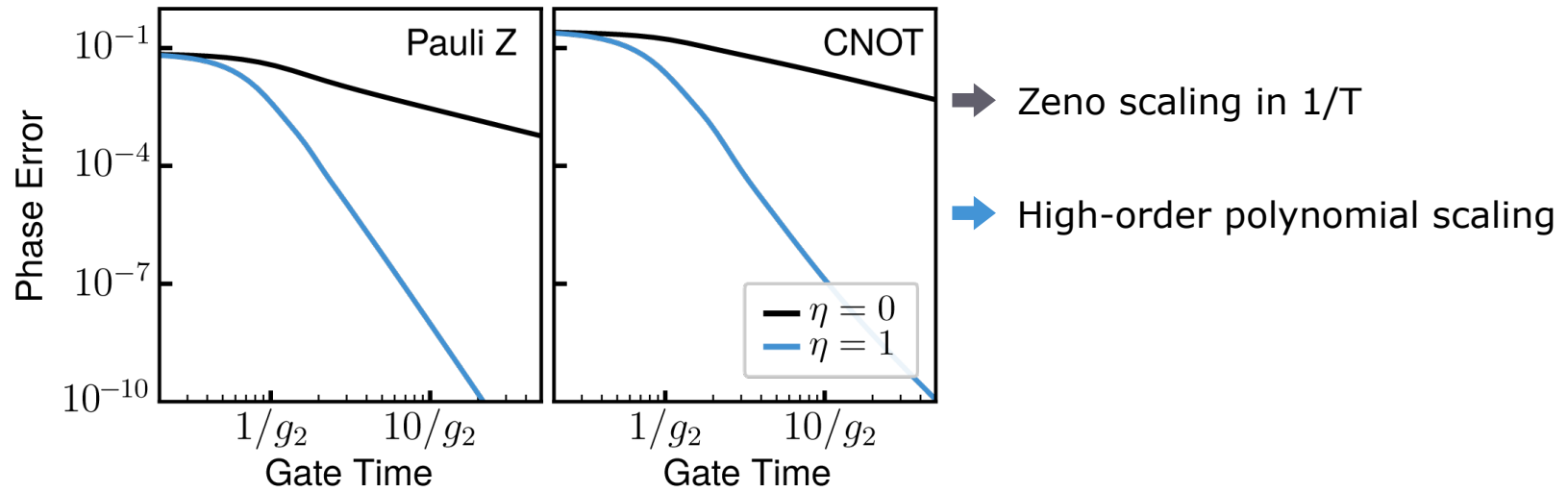


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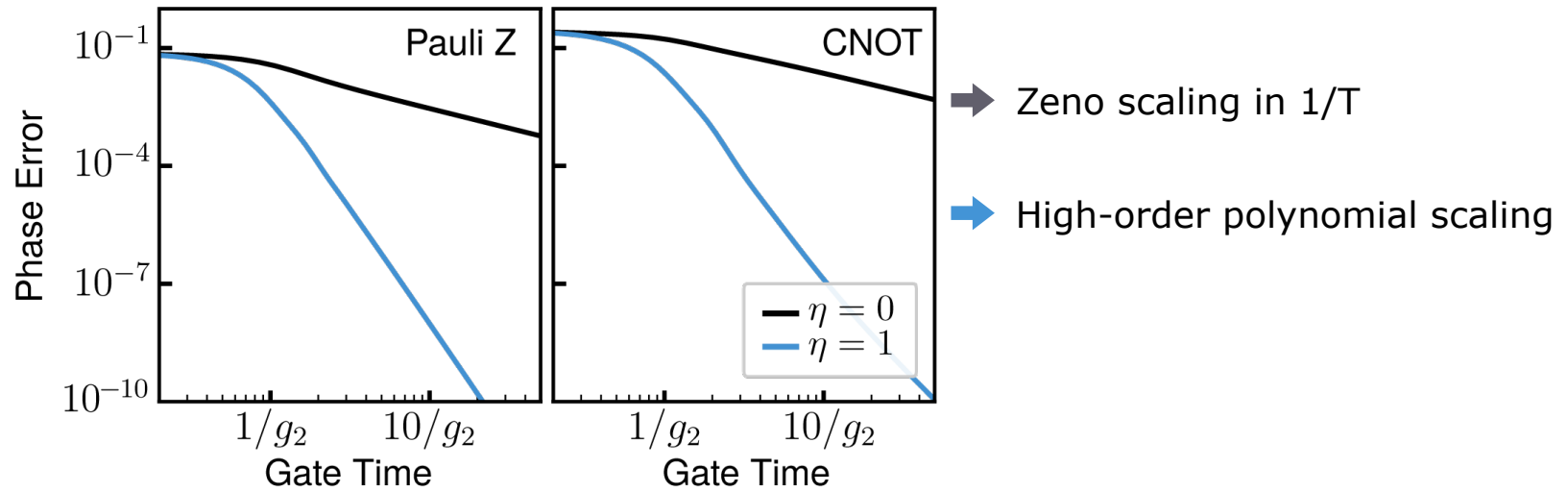


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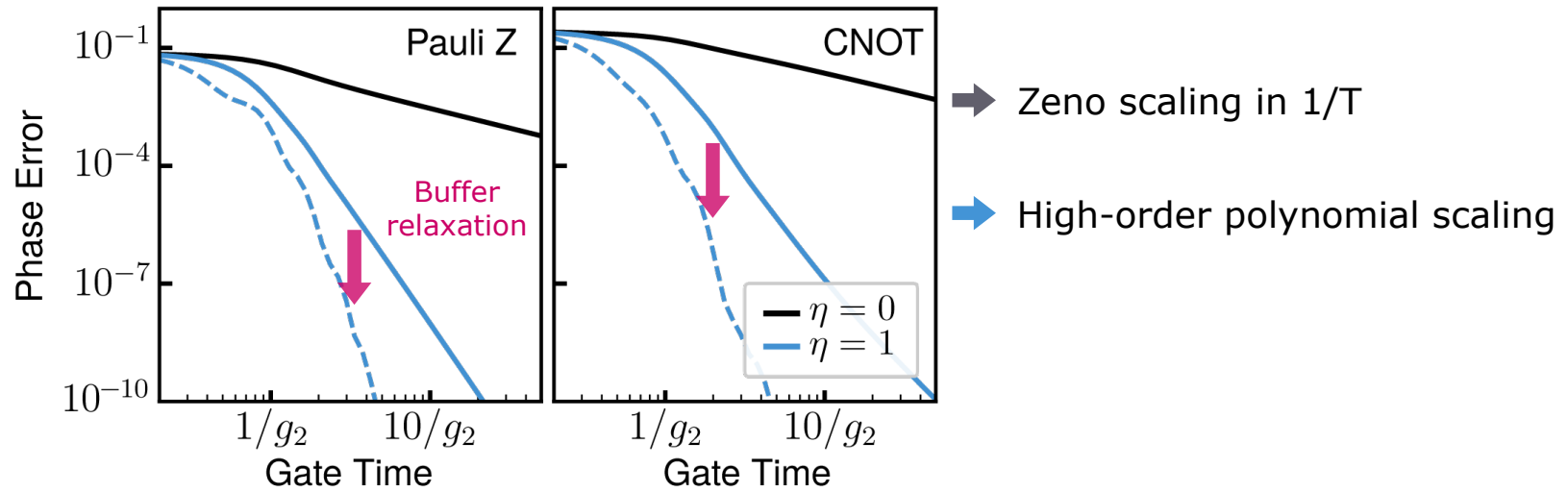
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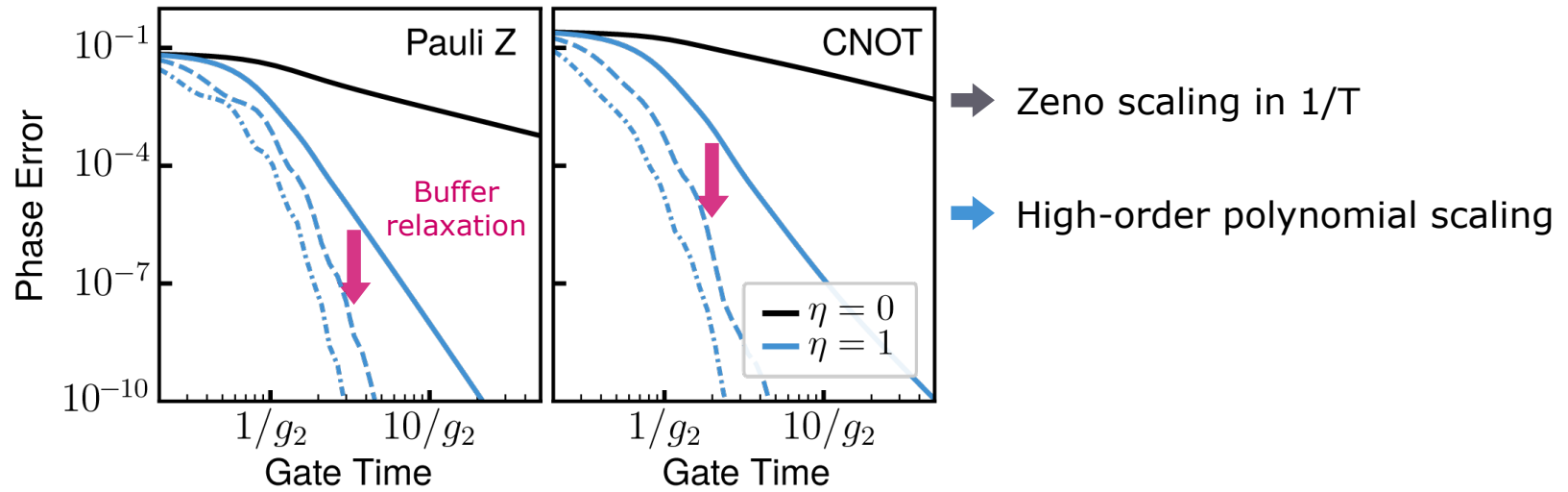
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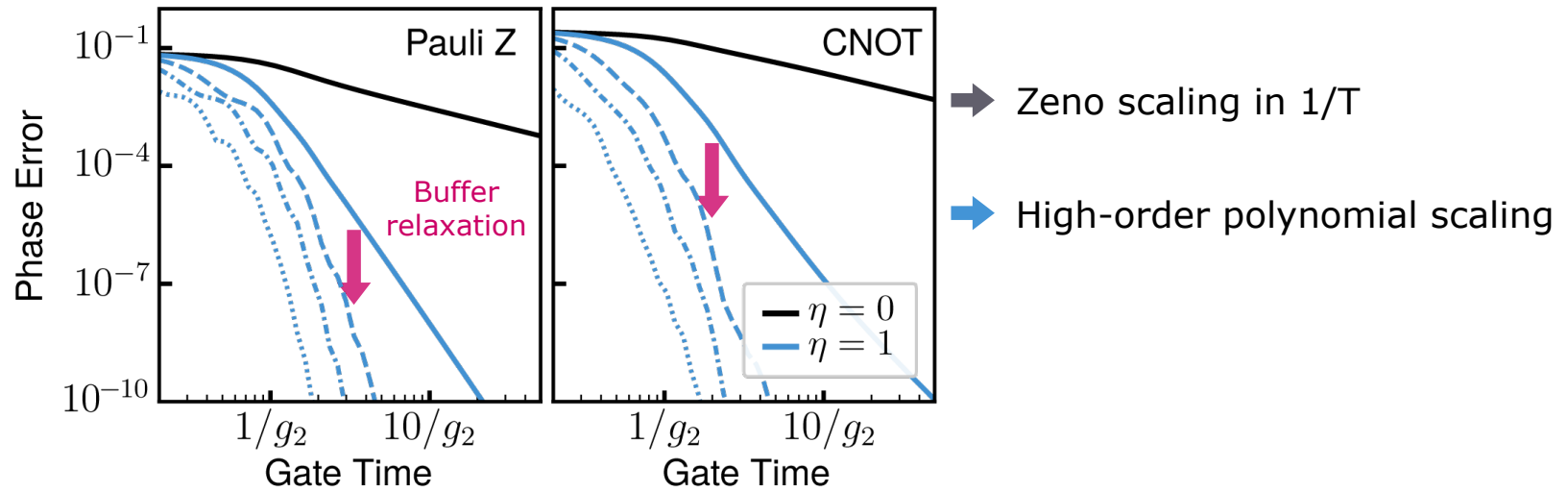
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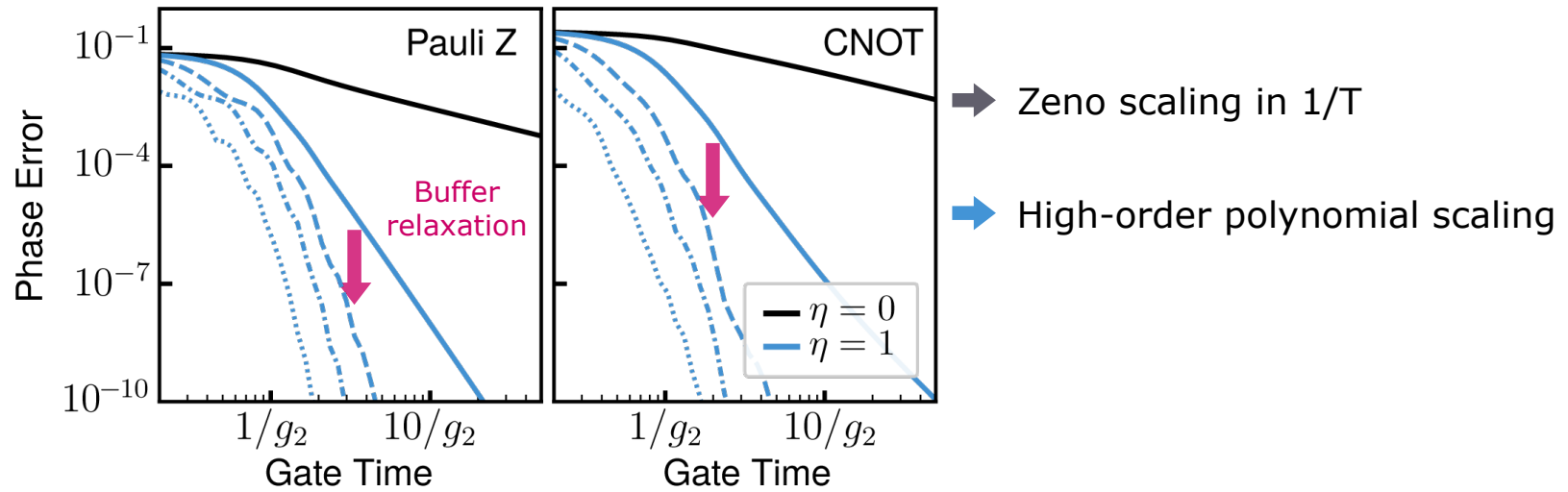
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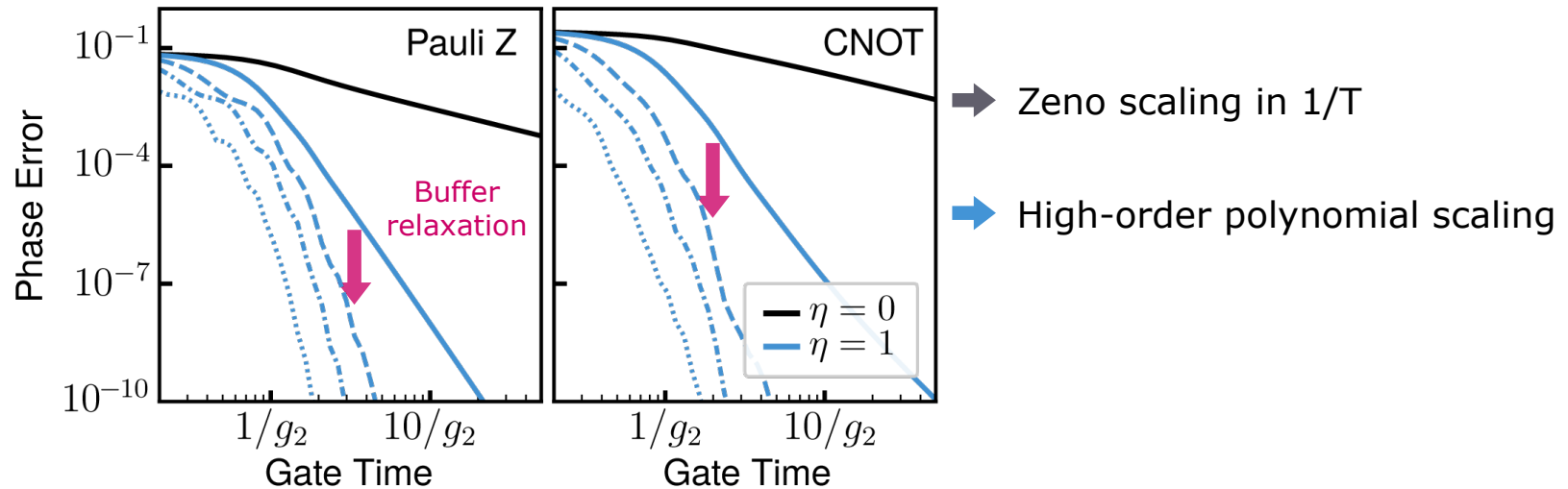
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- Detection of buffer photons ➔ Erasure errors or classical feedback
- Can apply N gates with a single relaxation per QEC cycle
- Gate fidelity limited by detection efficiency ➔ **Autonomous feedback**

Autonomous Feedback

Information is lost through the buffer mode

➔ **Feedback buffer detection autonomously**

$$\frac{d\rho}{dt} = -i[\mathbf{H}_{AB} + \mathbf{H}_Z, \rho] + \kappa_{ab}\mathcal{D}[ab]\rho$$

➔ Correlate buffer photon losses with cat parity-swaps

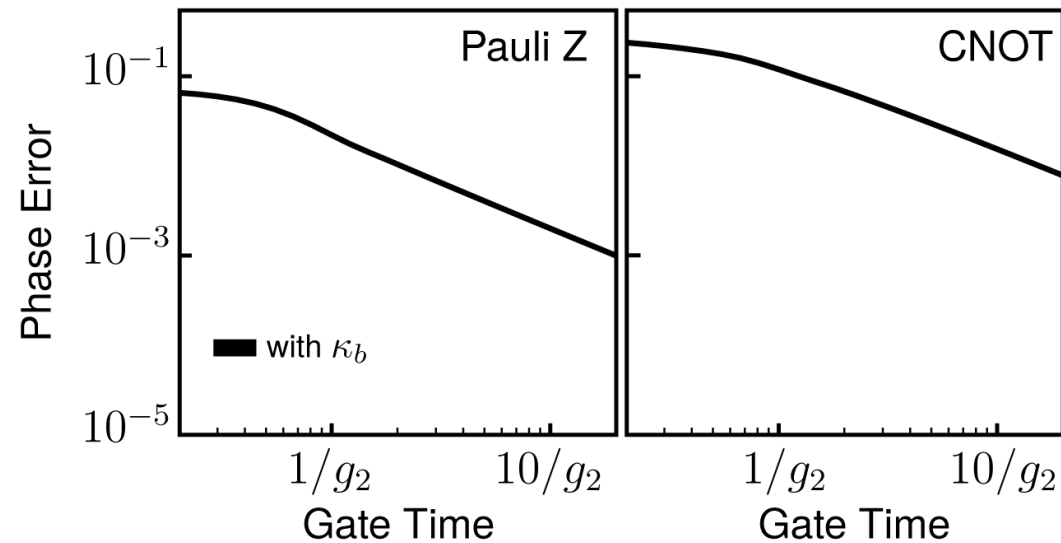
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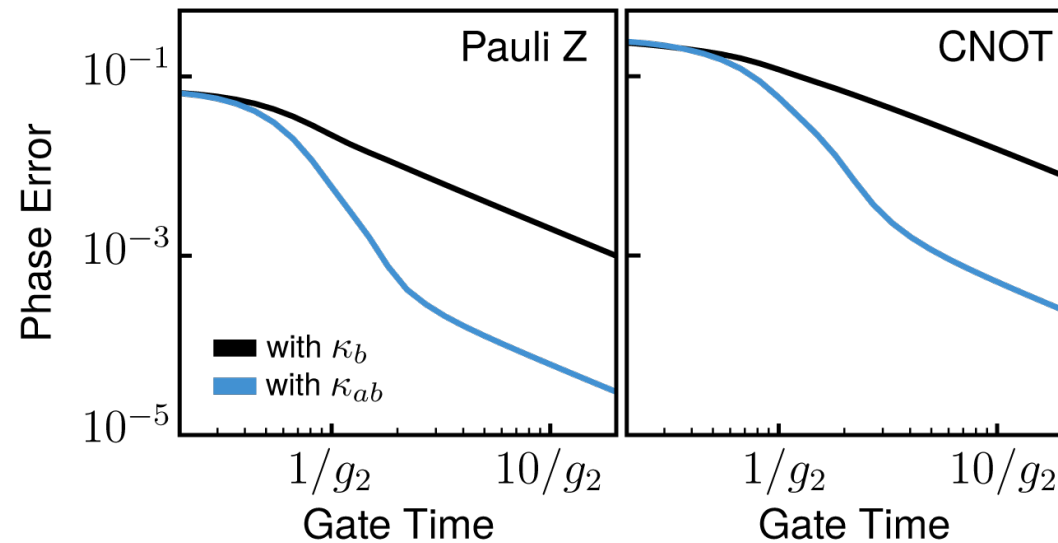
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- **x50** gate fidelity improvement (limited by second-order effects), independent of $|\alpha|^2$

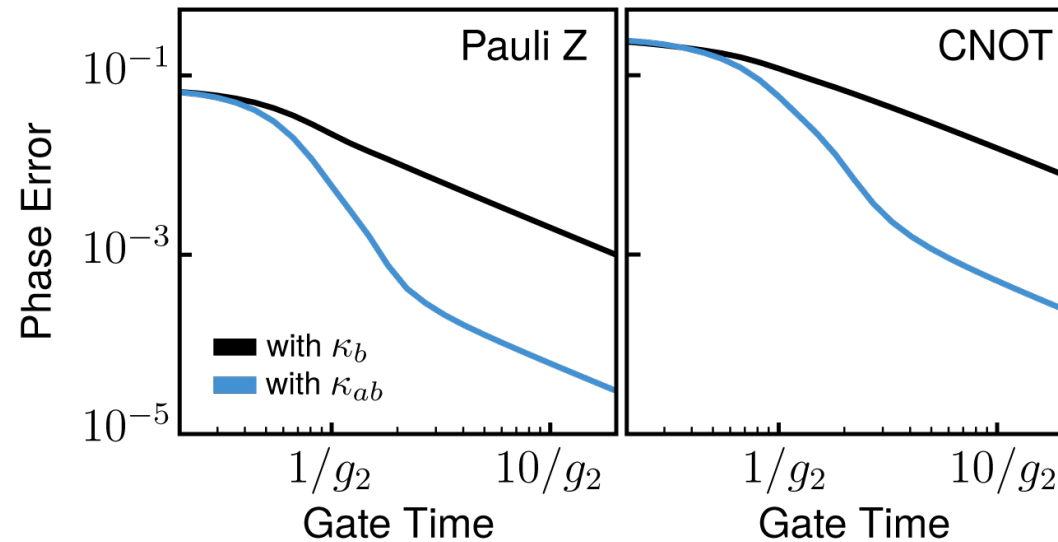
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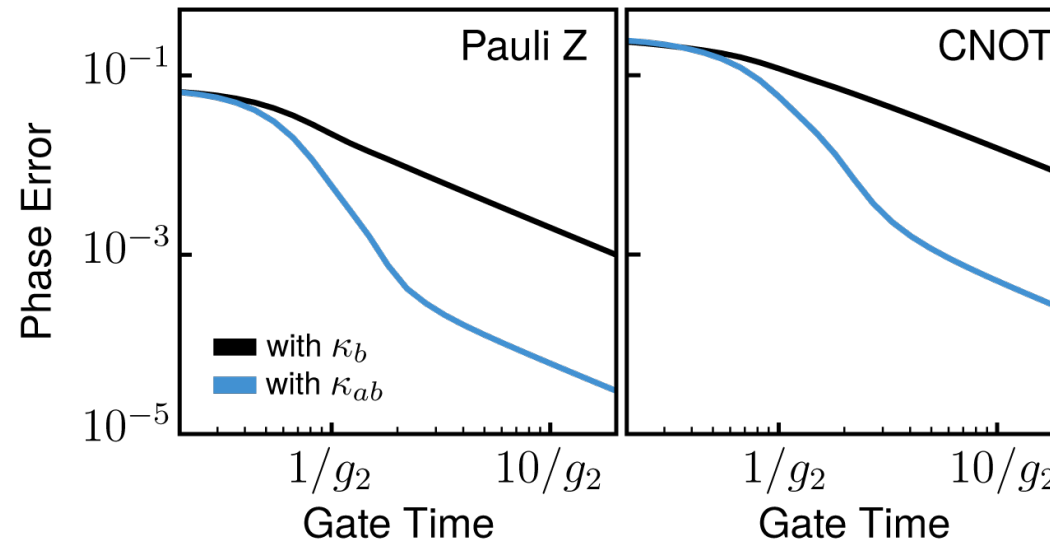
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- Autonomous correction of cavity losses with **squeezed** cats

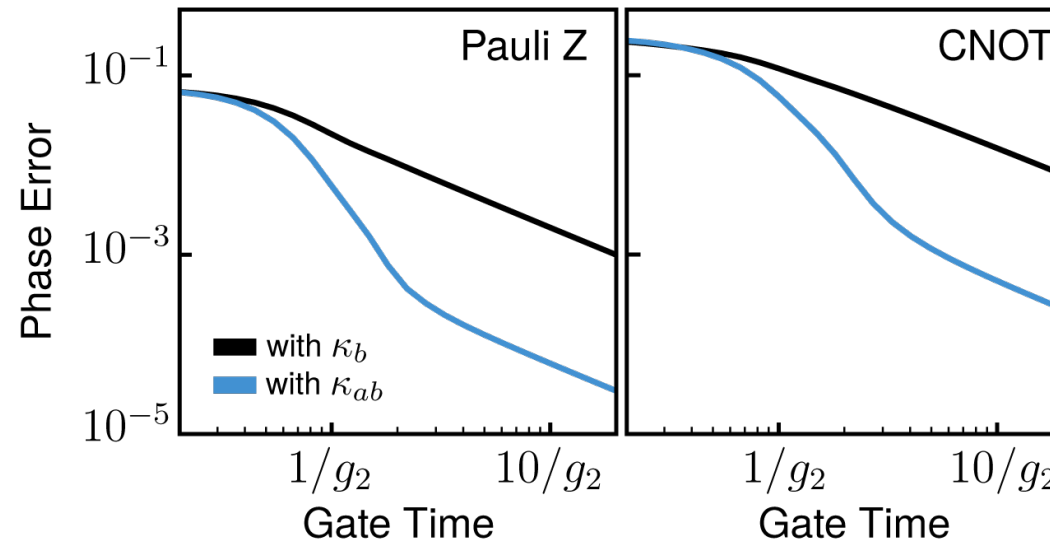
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- Autonomous correction of cavity losses with **squeezed** cats
- Can tune dissipation in situ

Designing High-Fidelity Gates for Dissipative Cat Qubits



- Buffer mode useful for more than dissipation engineering
- Information retrieval through photodetection: interesting conceptually but limited in practice
- Devised a scheme for autonomous correction of gate errors
- Check out [arXiv:2303.00760](https://arxiv.org/abs/2303.00760) for more details

